# Quadratic Relations of the Deformed $\boldsymbol{W}$-Algebra for the Twisted Affine Lie Algebra of Type $A_{2 N}^{(2)}$ 

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#### Abstract

We revisit the free field construction of the deformed $W$-algebra by Frenkel and Reshetikhin [Comm. Math. Phys. 197 (1998), 1-32], where the basic $W$-current has been identified. Herein, we establish a free field construction of higher $W$-currents of the deformed $W$-algebra associated with the twisted affine Lie algebra $A_{2 N}^{(2)}$. We obtain a closed set of quadratic relations and duality, which allows us to define deformed $W$-algebra $\mathcal{W}_{x, r}\left(A_{2 N}^{(2)}\right)$ using generators and relations.


Key words: deformed $W$-algebra; twisted affine algebra; quadratic relation; free field construction; exactly solvable model

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This paper is dedicated to Professor Michio Jimbo on the occasion of his 70th anniversary

## 1 Introduction

The deformed $W$-algebra $\mathcal{W}_{x, r}(\mathfrak{g})$ is a two-parameter deformation of the classical $W$-algebra $\mathcal{W}(\mathfrak{g})$. The deformation theory of the $W$-algebra has been studied in papers [2, 3, 4, 5, $6,8,10,12,13,14,16,17]$. For instance, free field constructions of the basic $W$-current $T_{1}(z)$ of $\mathcal{W}_{x, r}(\mathfrak{g})$ were suggested in the case when the underlying Lie algebra is of classical type. However, in comparison with the conformal case, the deformation theory of $W$-algebras is still not fully developed and understood. Moreover, finding quadratic relations of the deformed $W$-algebra $\mathcal{W}_{x, r}(\mathfrak{g})$ is still an unresolved problem.

In this paper, we generalize the study for $\mathcal{W}_{x, r}\left(A_{2}^{(2)}\right)^{1}$ by Brazhnikov and Lukyanov [3]. They obtained a quadratic relation for the $W$-current $T_{1}(z)$ of the deformed $W$-algebra $\mathcal{W}_{x, r}\left(A_{2}^{(2)}\right)$

$$
\begin{aligned}
f\left(\frac{z_{2}}{z_{1}}\right) & T_{1}\left(z_{1}\right) T_{1}\left(z_{2}\right)-f\left(\frac{z_{1}}{z_{2}}\right) T_{1}\left(z_{2}\right) T_{1}\left(z_{1}\right) \\
& =\delta\left(\frac{x^{-2} z_{2}}{z_{1}}\right) T_{1}\left(x^{-1} z_{2}\right)-\delta\left(\frac{x^{2} z_{2}}{z_{1}}\right) T_{1}\left(x z_{2}\right)+c\left(\delta\left(\frac{x^{-3} z_{2}}{z_{1}}\right)-\delta\left(\frac{x^{3} z_{2}}{z_{1}}\right)\right)
\end{aligned}
$$

with an appropriate constant $c$ and a function $f(z)$. This study aims to generalize the result for the cases $A_{2}^{(2)}$ to $A_{2 N}^{(2)}$. We introduce higher $W$-currents $T_{i}(z), 1 \leq i \leq 2 N$, by fusion of the free field construction of the basic $W$-current $T_{1}(z)$ of $\mathcal{W}_{x, r}\left(A_{2 N}^{(2)}\right)$ [8] (see formula (3.2)). We obtain a closed set of quadratic relations for the $W$-currents $T_{i}(z)$, which is completely different from

[^0]those in the case of deformed $W$-algebras associated with affine Lie algebras of types $A_{N}^{(1)}$ and $A(M, N)^{(1)}$ (see formula (3.4)). We refer the reader to references $[18,19]$ for the affine Lie superalgebra notation. We obtain the duality $T_{2 N+1-i}(z)=c_{i} T_{i}(z)$ with $1 \leq i \leq N$, which is a new phenomenon that does not occur in the case of deformed $W$-algebras associated with affine Lie algebras of types $A_{2}^{(2)}, A_{N}^{(1)}$, and $A(M, N)^{(1)}$ (see formula (3.3)). This allows us to define $\mathcal{W}_{x, r}\left(A_{2 N}^{(2)}\right)$ using generators and relations. We believe that this paper presents a key step toward extending our construction for general affine Lie algebras $\mathfrak{g}$, because the structures of the free field construction of the basic $W$-current $T_{1}(z)$ for the affine algebras other than that of type $A_{N}^{(1)}$ are quite similar to those of type $A_{2 N}^{(2)}$, not $A_{N}^{(1)}$. We have checked that there are similar quadratic relations as those for type $A_{2 N}^{(2)}$ in the case of type $B_{N}^{(1)}$ with small rank $N$.

The remainder of this paper is organized as follows. In Section 2, we review the free field construction of the basic $W$-current $T_{1}(z)$ of the deformed $W$-algebra $\mathcal{W}_{x, r}\left(A_{2 N}^{(2)}\right)$ [8]. In Section 3 , we introduce higher $W$-currents $T_{i}(z)$ and present a closed set of quadratic relations and duality. We also obtain the $q$-Poisson algebra in the classical limit. In Section 4 , we establish proofs of Proposition 3.1 and Theorem 3.2. Section 5 is devoted to discussion. In Appendices A and $B$, we summarize normal ordering rules.

## 2 Free field construction

In this section, we define notation and review the free field construction of the basic $W$-current $T_{1}(z)$ of $\mathcal{W}_{x, r}\left(A_{2 N}^{(2)}\right)$. Throughout this paper, we fix a natural number $N=1,2,3, \ldots$, a real number $r>1$, and a complex number $x$ with $0<|x|<1$.

### 2.1 Notation

In this section, we use complex numbers $a, w, q$, and $p$ with $w \neq 0, q \neq 0, \pm 1$, and $|p|<1$. For any integer $n$, we define $q$-integers

$$
[n]_{q}=\frac{q^{n}-q^{-n}}{q-q^{-1}}
$$

We use symbols for infinite products,

$$
(a ; p)_{\infty}=\prod_{k=0}^{\infty}\left(1-a p^{k}\right), \quad\left(a_{1}, a_{2}, \ldots, a_{N} ; p\right)_{\infty}=\prod_{i=1}^{N}\left(a_{i} ; p\right)_{\infty}
$$

for complex numbers $a_{1}, a_{2}, \ldots, a_{N}$. The following standard formulas are used,

$$
\exp \left(-\sum_{m=1}^{\infty} \frac{1}{m} a^{m}\right)=1-a, \quad \exp \left(-\sum_{m=1}^{\infty} \frac{1}{m} \frac{a^{m}}{1-p^{m}}\right)=(a ; p)_{\infty}
$$

We use the elliptic theta function $\Theta_{p}(w)$ and the compact notation $\Theta_{p}\left(w_{1}, w_{2}, \ldots, w_{N}\right)$,

$$
\Theta_{p}(w)=\left(p, w, p w^{-1} ; p\right)_{\infty}, \quad \Theta_{p}\left(w_{1}, w_{2}, \ldots, w_{N}\right)=\prod_{i=1}^{N} \Theta_{p}\left(w_{i}\right)
$$

for complex numbers $w_{1}, w_{2}, \ldots, w_{N} \neq 0$. Define $\delta(z)$ by the formal series

$$
\delta(z)=\sum_{m \in \mathbb{Z}} z^{m}
$$

### 2.2 Twisted affine Lie algebra of type $A_{2 N}^{(2)}$

In this section we recall the definition of the twisted affine Lie algebra of type $A_{2 N}^{(2)}, N=$ $1,2,3, \ldots$, in [11]. The Dynkin diagram of type $A_{2 N}^{(2)}$ is given by

$$
A_{2 N}^{(2)} \text { with } N \geq 2
$$



The corresponding Cartan matrix $A=\left(A_{i, j}\right)_{i, j=0}^{N}$ of type $A_{2 N}^{(2)}$ is given by

$$
A=\left(\begin{array}{rrrrrrc}
2 & -2 & 0 & \cdots & \cdots & \cdots & 0 \\
-1 & 2 & -1 & \ddots & & & \vdots \\
0 & -1 & 2 & \ddots & \ddots & & \vdots \\
\vdots & \ddots & \ddots & \ddots & \ddots & \ddots & \vdots \\
\vdots & & \ddots & \ddots & 2 & -1 & 0 \\
\vdots & & & \ddots & -1 & 2 & -2 \\
0 & \cdots & \cdots & \cdots & 0 & -1 & 2
\end{array}\right)
$$

with $N \geq 2$, and

$$
A=\left(\begin{array}{rr}
2 & -4 \\
-1 & 2
\end{array}\right)
$$

with $N=1$. We set the labels $a_{i}=2,0 \leq i \leq N-1, a_{N}=1$, and the co-labels $a_{0}^{\vee}=1, a_{i}^{\vee}=2$, $1 \leq i \leq N$. We set $D=\operatorname{diag}\left(a_{0} a_{0}^{\vee-1}, a_{1} a_{1}^{\vee-1}, \ldots, a_{N} a_{N}^{\vee-1}\right)$. We obtain $A=D B$, where $B$ is a symmetric matrix. Thus, the Cartan matrix $A$ is symmetrizable. Let $\mathfrak{h}$ be an $(N+2)$ dimensional vector space over $\mathbf{C}$. Let $\left\{h_{0}, h_{1}, \ldots, h_{N}, d\right\}$ be a basis of $\mathfrak{h}$, and $\left\{\alpha_{0}, \alpha_{1}, \ldots, \alpha_{N}, \Lambda_{0}\right\}$ a basis of $\mathfrak{h}^{*}=\operatorname{Hom}_{\mathbf{C}}(\mathfrak{h}, \mathbf{C})$ such that we have with respect to pairing $\langle\cdot, \cdot\rangle: \mathfrak{h} \times \mathfrak{h}^{*} \rightarrow \mathbf{C}$

$$
\begin{array}{llll}
\left\langle h_{i}, \alpha_{j}\right\rangle=A_{i, j}, & & 0 \leq i, j \leq N, &
\end{array}\left\langle d, \alpha_{i}\right\rangle=\delta_{0, i}, ~ 子 \begin{array}{ll}
\left\langle h_{i}, \Lambda_{0}\right\rangle=\delta_{i, 0}, & \\
0 \leq i \leq N, & \\
\left\langle d, \Lambda_{0}\right\rangle=0 .
\end{array}
$$

Let $\mathfrak{g}(A)$ be the affine Lie algebra associated with the Cartan matrix $A$. Since $A$ is symmetrizable, it is defined as the Lie algebra generated by $e_{i}, f_{i}, 0 \leq i \leq N$, and $\mathfrak{h}$ with the following relations:

$$
\begin{aligned}
& {\left[e_{i}, f_{j}\right]=\delta_{i, j} h_{i}, \quad 0 \leq i, j \leq N, \quad\left[h, h^{\prime}\right]=0, \quad h, h^{\prime} \in \mathfrak{h}, \quad\left[h, e_{i}\right]=\left\langle h, \alpha_{i}\right\rangle e_{i},} \\
& {\left[h, f_{i}\right]=-\left\langle h, \alpha_{i}\right\rangle f_{i}, \quad h \in \mathfrak{h}, \quad 0 \leq i \leq N,} \\
& \left(\operatorname{ad} e_{i}\right)^{-A_{i, j}+1} e_{j}=0, \quad\left(\operatorname{ad} f_{i}\right)^{-A_{i, j}+1} f_{j}=0, \quad 0 \leq i, j \leq N, \quad i \neq j .
\end{aligned}
$$

Here we used the adjoint action $(\operatorname{ad} x) y=[x, y]$.

### 2.3 Free field construction

In this section, we recall the free field construction of the basic $W$-current $T_{1}(z)$ and of the screening operators $S_{i}$ of the deformed $W$-algebra $\mathcal{W}_{x, r}\left(A_{2 N}^{(2)}\right)$ introduced by Frenkel and Reshetikhin [8].

First, we define the $N \times N$ symmetric matrix $B(m)=\left(B_{i, j}(m)\right)_{i, j=1}^{N}, m \in \mathbb{Z}$, associated with $A_{2 N}^{(2)}, N=1,2,3, \ldots$, as follows:

$$
\begin{aligned}
& B_{i, j}(m)= \begin{cases}\frac{[2 m]_{x}}{[m]_{x}}, & 1 \leq i, j \leq N-1, \quad i=j, \\
\frac{[2 m]_{x}-[m]_{x}}{[m]_{x}}, & i=j=N, \\
-1, & |i-j|=1, \\
0, & |i-j| \geq 2,\end{cases} \\
& B_{i, j}(0)=\left\{\begin{aligned}
2, & 1 \leq i, j \leq N-1, \quad i=j, \\
1, & i=j=N, \\
-1, & |i-j|=1, \\
0, & |i-j| \geq 2 .
\end{aligned}\right.
\end{aligned}
$$

We introduce the Heisenberg algebra $\mathcal{H}_{x, r}$ with generators $a_{i}(m), Q_{i}, m \in \mathbb{Z}, 1 \leq i \leq N$, satisfying

$$
\begin{aligned}
& {\left[a_{i}(m), a_{j}(n)\right]=\frac{1}{m}[r m]_{x}[(r-1) m]_{x} B_{i, j}(m)\left(x-x^{-1}\right)^{2} \delta_{m+n, 0}, \quad m, n \neq 0, \quad 1 \leq i, j \leq N,} \\
& {\left[a_{i}(0), Q_{j}\right]=B_{i, j}(0), \quad 1 \leq i, j \leq N}
\end{aligned}
$$

The remaining commutators vanish. The generators $a_{i}(m), Q_{i}$ are "root" type generators of $\mathcal{H}_{x, r}$. There is a unique set of "fundamental weight" type generators $y_{i}(m), Q_{i}^{y}, m \in \mathbb{Z}, 1 \leq i \leq N$, which satisfy the following relations

$$
\begin{aligned}
& {\left[y_{i}(m), a_{j}(n)\right]=\frac{1}{m}[r m]_{x}[(r-1) m]_{x}\left(x-x^{-1}\right)^{2} \delta_{i, j} \delta_{m+n, 0}, \quad m, n \neq 0, \quad 1 \leq i, j \leq N} \\
& {\left[y_{i}(0), Q_{j}\right]=\delta_{i, j}, \quad\left[a_{i}(0), Q_{j}^{y}\right]=\delta_{i, j}, \quad\left[y_{i}(0), a_{j}(m)\right]=0, \quad m \in \mathbb{Z}, \quad 1 \leq i, j \leq N}
\end{aligned}
$$

The explicit formulas for $y_{i}(m)$ and $Q_{j}^{y}$ are given in (A.7). We use the normal ordering : : on $\mathcal{H}_{x, r}$ that satisfies

$$
: a_{i}(m) a_{j}(n):=\left\{\begin{array}{ll}
a_{i}(m) a_{j}(n), & m<0, \\
a_{j}(n) a_{i}(m), & m \geq 0,
\end{array} \quad m, n \in \mathbb{Z}, \quad 1 \leq i, j \leq N .\right.
$$

Let $|0\rangle \neq 0$ be the Fock vacuum of the Fock space of $\mathcal{H}_{x, r}$ such that $a_{i}(m)|0\rangle=0, m \geq 0$, $1 \leq i \leq N$. Let $\pi_{\lambda}$ be the Fock space of $\mathcal{H}_{x, r}$ generated by $|\lambda\rangle=\mathrm{e}^{\lambda}|0\rangle, \lambda=\sum_{j=1}^{N} \lambda_{j} Q_{j}^{y}$. We obtain

$$
\begin{equation*}
a_{i}(0)|\lambda\rangle=\lambda_{i}|\lambda\rangle, \quad a_{i}(m)|\lambda\rangle=0, \quad m>0, \quad 1 \leq i \leq N . \tag{2.1}
\end{equation*}
$$

We work in the Fock space $\pi_{\lambda}$ of the Heisenberg algebra $\mathcal{H}_{x, r}$. Let the vertex operators $A_{i}(z)$, $Y_{i}(z)$, and $S_{i}(z), 1 \leq i \leq N$, be

$$
\begin{equation*}
A_{i}(z)=x^{r a_{i}(0)}: \exp \left(\sum_{m \neq 0} a_{i}(m) z^{-m}\right):, \tag{2.2}
\end{equation*}
$$

$$
\begin{align*}
& Y_{i}(z)=x^{r y_{i}(0)}: \exp \left(\sum_{m \neq 0} y_{i}(m) z^{-m}\right):,  \tag{2.3}\\
& S_{i}(z)=z^{\frac{r-1}{2 r} B_{i, i}(0)} \mathrm{e}^{-\sqrt{\frac{r-1}{r}} Q_{i}} z^{-\sqrt{\frac{r-1}{r}} a_{i}(0)}: \exp \left(\sum_{m \neq 0} \frac{a_{i}(m)}{x^{r m}-x^{-r m}} z^{-m}\right): . \tag{2.4}
\end{align*}
$$

The main parts of (2.2), (2.3), and (2.4) are the same as those of [8]. We corrected the misprints in the formulas for $A_{i}(z), Y_{i}(z)$, and $S_{i}(z)$ in [8] by multiplying (2.2) and (2.3) by constants and multiplying (2.4) by $z^{\frac{r-1}{2 r} B_{i, i}(0)}$. With our fine-tuning, both (3.3) and (3.5) hold.

Let $J_{N}=\{1,2, \ldots, N, 0, \bar{N}, \ldots, \overline{2}, \overline{1}\}$. Here, the indices are ordered as

$$
1 \prec 2 \prec \cdots \prec N \prec 0 \prec \bar{N} \prec \cdots \prec \overline{2} \prec \overline{1} .
$$

Let $\overline{\bar{k}}=k, k=1,2, \ldots, N$, and $\overline{0}=0$. The indices $i, j \in J_{N}$ satisfy $i \prec j$ if and only if $\bar{j} \prec \bar{i}$. We define $\bar{I}=\left\{\overline{i_{1}}, \overline{i_{2}}, \ldots, \overline{i_{k}}\right\}$ for a subset $I \subset J_{N}, I=\left\{i_{1}, i_{2}, \ldots, i_{k}\right\}$. Let $T_{1}(z)$ be the generating series with operator valued coefficients acting on the Fock space $\pi_{\lambda}$,

$$
T_{1}(z)=\sum_{i \in J_{N}} \Lambda_{i}(z),
$$

where

$$
\begin{align*}
& \Lambda_{1}(z)=Y_{1}(z), \quad \Lambda_{k}(z)=: \Lambda_{k-1}(z) A_{k-1}\left(x^{-k+1} z\right)^{-1}:, \quad 2 \leq k \leq N, \\
& \Lambda_{0}(z)=\frac{\left[r-\frac{1}{2}\right]_{x}}{\left[\frac{1}{2}\right]_{x}}: \Lambda_{N}(z) A_{N}\left(x^{-N} z\right)^{-1}: \\
& \Lambda_{\bar{N}}(z)=\frac{\left[\frac{1}{2}\right]_{x}}{\left[r-\frac{1}{2}\right]_{x}}: \Lambda_{0}(z) A_{N}\left(x^{-N-1} z\right)^{-1}:, \\
& \Lambda_{\bar{k}}(z)=: \Lambda_{\overline{k+1}}(z) A_{k}\left(x^{-2 N+k-1} z\right)^{-1}:, \quad 1 \leq k \leq N-1 . \tag{2.5}
\end{align*}
$$

We call $T_{1}(z)$ the basic $W$-current of the deformed $W$-algebra $\mathcal{W}_{x, r}\left(A_{2 N}^{(2)}\right)$.
Let $\pi_{\mu}$ be the Fock space of $\mathcal{H}_{x, r}$ generated by $|\mu\rangle=\mathrm{e}^{\mu}|0\rangle$ with $\mu=\sum_{i=1}^{N} \mu_{i} Q_{i}^{y}$, where we choose $\mu_{i} \in \frac{1}{2} \sqrt{\frac{r-1}{r}} B_{i, i}(0)+\sqrt{\frac{r}{r-1}} \mathbb{Z}, 1 \leq i \leq N$. From (2.1) and (2.4) the power of $w$ in $S_{i}(w)$, $w^{\frac{r-1}{2 r} B_{i, i}(0)} w^{-\sqrt{\frac{r-1}{r}} a_{i}(0)}$, takes values in integers on $\pi_{\mu}$. Hence, $S_{i}$ is well-defined on $\pi_{\mu}$. We define the screening operators $S_{i}, 1 \leq i \leq N$, acting on the Fock space $\pi_{\mu}$ as

$$
\begin{equation*}
S_{i}=\oint \frac{\mathrm{d} w}{2 \pi \sqrt{-1} w} S_{i}(w) \tag{2.6}
\end{equation*}
$$

The integral in formula (2.6) means the residue at zero.

## 3 Quadratic relations

In this section, we introduce the higher $W$-currents $T_{i}(z)$ and present a set of quadratic relations between $T_{i}(z)$ for the deformed $W$-algebra $\mathcal{W}_{x, r}\left(A_{2 N}^{(2)}\right)$.

### 3.1 Quadratic relations

We define the formal series $\Delta(z) \in \mathbf{C}[[z]]$ and the constant $c(x, r)$ as

$$
\Delta(z)=\frac{\left(1-x^{2 r-1} z\right)\left(1-x^{-2 r+1} z\right)}{(1-x z)\left(1-x^{-1} z\right)}, \quad c(x, r)=[r]_{x}[r-1]_{x}\left(x-x^{-1}\right)
$$

The formal series $\Delta(z)$ satisfies

$$
\begin{aligned}
& \Delta(z)-\Delta\left(z^{-1}\right)=c(x, r)\left(\delta\left(x^{-1} z\right)-\delta(x z)\right) \\
& \Delta(z) \Delta\left(x^{s} z\right)-\Delta\left(z^{-1}\right) \Delta\left(x^{-s} z^{-1}\right) \\
& \quad=c(x, r)\left\{\Delta\left(x^{s+1}\right)\left(\delta\left(x^{-1} z\right)-\delta\left(x^{s+1} z\right)\right)+\Delta\left(x^{s-1}\right)\left(\delta\left(x^{s-1} z\right)-\delta(x z)\right)\right\}, \quad s \neq 0, \pm 2
\end{aligned}
$$

We define the structure functions $f_{i, j}(z), i, j=0,1,2, \ldots$, as

$$
\begin{align*}
f_{i, j}(z)= & \exp \left(-\sum_{m=1}^{\infty} \frac{1}{m}[(r-1) m]_{x}[r m]_{x}\left(x-x^{-1}\right)^{2}\right. \\
& \left.\times \frac{[\operatorname{Min}(i, j) m]_{x}\left([(N+1-\operatorname{Max}(i, j)) m]_{x}-[(N-\operatorname{Max}(i, j)) m]_{x}\right)}{[m]_{x}\left([(N+1) m]_{x}-[N m]_{x}\right)} z^{m}\right) . \tag{3.1}
\end{align*}
$$

The ratio of the structure functions $f_{1,1}(z)$ is

$$
\frac{f_{1,1}\left(z^{-1}\right)}{f_{1,1}(z)}=-z \frac{\Theta_{x^{4 N+2}}\left(x^{2} z, x^{2 N-1} z, x^{4 N+2-2 r} z, x^{4 N+2 r} z, x^{2 N+1+2 r} z, x^{2 N-2 r+3} z\right)}{\Theta_{x^{4 N+2}}\left(x^{2} / z, x^{2 N-1} / z, x^{4 N+2-2 r} / z, x^{4 N+2 r} / z, x^{2 N+1+2 r} / z, x^{2 N-2 r+3} / z\right)} .
$$

We introduce higher $W$-currents $T_{i}(z)$ as follows:

$$
\begin{align*}
& T_{0}(z)=1, \quad T_{1}(z)=\sum_{i \in J_{N}} \Lambda_{i}(z), \\
& T_{i}(z)=\sum_{\substack{\Omega_{i} \subset J_{N} \\
\left|\Omega_{i}\right|=i}} d_{\Omega_{i}}(x, r) \vec{\Lambda}_{\Omega_{i}}(z), \quad 2 \leq i \leq 2 N+1 . \tag{3.2}
\end{align*}
$$

Here, for a subset $\Omega_{i}=\left\{s_{1}, s_{2}, \ldots, s_{i}\right\} \subset J_{N}$ with $s_{1} \prec s_{2} \prec \cdots \prec s_{i}$, we set

$$
\begin{aligned}
& d_{\Omega_{i}}(x, r)=\prod_{\substack{1 \leq p<q \leq i \\
s_{q}=\bar{s}_{p}}} \Delta\left(x^{2\left(q-p+s_{p}-N-1\right)}\right), \quad d_{\varnothing}(x, r)=1, \\
& \vec{\Lambda}_{\Omega_{i}}(z)=: \Lambda_{s_{1}}\left(x^{-i+1} z\right) \Lambda_{s_{2}}\left(x^{-i+3} z\right) \cdots \Lambda_{s_{i}}\left(x^{i-1} z\right):, \quad \vec{\Lambda}_{\varnothing}(z)=1 .
\end{aligned}
$$

Proposition 3.1. The $W$-currents $T_{i}(z)$ satisfy the duality

$$
\begin{equation*}
T_{2 N+1-i}(z)=\frac{\left[r-\frac{1}{2}\right]_{x}}{\left[\frac{1}{2}\right]_{x}} \prod_{k=1}^{N-i} \Delta\left(x^{2 k}\right) T_{i}(z), \quad 0 \leq i \leq N . \tag{3.3}
\end{equation*}
$$

Theorem 3.2. The $W$-currents $T_{i}(z)$ satisfy the set of quadratic relations

$$
\begin{align*}
& f_{i, j}\left(\frac{z_{2}}{z_{1}}\right) T_{i}\left(z_{1}\right) T_{j}\left(z_{2}\right)-f_{j, i}\left(\frac{z_{1}}{z_{2}}\right) T_{j}\left(z_{2}\right) T_{i}\left(z_{1}\right) \\
& \quad=c(x, r) \sum_{k=1}^{i} \prod_{l=1}^{k-1} \Delta\left(x^{2 l+1}\right)\left(\delta\left(\frac{x^{-j+i-2 k} z_{2}}{z_{1}}\right) f_{i-k, j+k}\left(x^{j-i}\right) T_{i-k}\left(x^{k} z_{1}\right) T_{j+k}\left(x^{-k} z_{2}\right)\right. \\
& \left.\quad-\delta\left(\frac{x^{j-i+2 k} z_{2}}{z_{1}}\right) f_{i-k, j+k}\left(x^{-j+i}\right) T_{i-k}\left(x^{-k} z_{1}\right) T_{j+k}\left(x^{k} z_{2}\right)\right) \\
& \quad+c(x, r) \prod_{l=1}^{i-1} \Delta\left(x^{2 l+1}\right) \prod_{l=N+1-j}^{N+i-j} \Delta\left(x^{2 l}\right)\left(\delta\left(\frac{x^{-2 N+j-i-1} z_{2}}{z_{1}}\right) T_{j-i}\left(x^{-i} z_{2}\right)\right. \\
& \left.\quad-\delta\left(\frac{x^{2 N-j+i+1} z_{2}}{z_{1}}\right) T_{j-i}\left(x^{i} z_{2}\right)\right), \quad 1 \leq i \leq j \leq N . \tag{3.4}
\end{align*}
$$

Here, we use $f_{i, j}(z)$ introduced in (3.1).

In view of Proposition 3.1 and Theorem 3.2, we obtain the following definition.
Definition 3.3. Let $W$ be the free complex associative algebra generated by elements $\bar{T}_{i}[m]$, $m \in \mathbb{Z}, 1 \leq i \leq 2 N, I_{K}$ the left ideal generated by elements $\bar{T}_{i}[m], m \geq K \in \mathbb{N}, 1 \leq i \leq 2 N$, and

$$
\widehat{W}=\lim _{\leftarrow} W / I_{K}
$$

The deformed $W$-algebra $\mathcal{W}_{x, r}\left(A_{2 N}^{(2)}\right)$ is the quotient of $\widehat{W}$ by the two-sided ideal generated by the coefficients of the generating series which are the differences of the right hand sides and of the left hand sides of the relations (3.3) and (3.4), where the generating series $T_{i}(z)$ are replaced with $\bar{T}_{i}(z)=\sum_{m \in \mathbb{Z}} \bar{T}_{i}[m] z^{-m}, 1 \leq i \leq 2 N$, and $\bar{T}_{0}(z)=1$.

The justification of this definition is presented later. We compare this definition of the deformed $W$-algebra with other definitions in Section 5.

Lemma 3.4. The $W$-currents $T_{i}(z)$ commute with the screening operators $S_{j}$,

$$
\begin{equation*}
\left[T_{i}(z), S_{j}\right]=0, \quad 1 \leq i \leq 2 N, \quad 1 \leq j \leq N . \tag{3.5}
\end{equation*}
$$

We present the proofs of Proposition 3.1, Theorem 3.2, and Lemma 3.4 in Section 4.

### 3.2 Classical limit

The deformed $W$-algebra $\mathcal{W}_{x, r}(\mathfrak{g})$ yields a $q$-Poisson $W$-algebra [7, 8, 9, 15] in the classical limit. As an application of the quadratic relations (3.4), we obtain a $q$-Poisson $W$-algebra of type $A_{2 N}^{(2)}$. We set parameters $q=x^{2 r}$ and $\beta=(r-1) / r$. We define the $q$-Poisson bracket $\{\cdot, \cdot\}$ by taking the classical limit $\beta \rightarrow 0$ with $q$ fixed as

$$
\left\{T_{i}^{\mathrm{PB}}[m], T_{j}^{\mathrm{PB}}[n]\right\}=\lim _{\beta \rightarrow 0} \frac{1}{2 \beta \log q}\left[T_{i}[m], T_{j}[n]\right] .
$$

Here, we introduce $T_{i}^{\mathrm{PB}}[m]$ by

$$
T_{i}(z)=\sum_{m \in \mathbb{Z}} T_{i}[m] z^{-m} \longrightarrow T_{i}^{\mathrm{PB}}(z)=\sum_{m \in \mathbb{Z}} T_{i}^{\mathrm{PB}}[m] z^{-m}, \quad \beta \rightarrow 0, \quad q \text { fixed } .
$$

The $\beta$-expansions of the structure functions are given as

$$
\begin{aligned}
f_{i, j}(z)= & 1-2 \beta \log q\left(q-q^{-1}\right) \sum_{m=1}^{\infty}[\operatorname{Min}(i, j) m]_{q} \\
& \times \frac{[(N+1-\operatorname{Max}(i, j)) m]_{q}-[(N-\operatorname{Max}(i, j)) m]_{q}}{[(N+1) m]_{q}-[N m]_{q}} z^{m}+O\left(\beta^{2}\right), \quad i, j \geq 1, \\
c(x, r)= & 2 \beta \log q+O\left(\beta^{2}\right) .
\end{aligned}
$$

As corollaries of Proposition 3.1 and Theorem 3.2 we obtain the following.
Corollary 3.5. For the $q$-Poisson $W$-algebra associated with affine Lie algebra of type $A_{2 N}^{(2)}$, the currents $T_{i}^{\mathrm{PB}}(z)$ satisfy

$$
\begin{aligned}
& \left\{T_{i}^{\mathrm{PB}}\left(z_{1}\right), T_{j}^{\mathrm{PB}}\left(z_{2}\right)\right\} \\
& \quad=\left(q-q^{-1}\right) C_{i, j}\left(\frac{z_{2}}{z_{1}}\right) T_{i}^{\mathrm{PB}}\left(z_{1}\right) T_{j}^{\mathrm{PB}}\left(z_{2}\right)+\sum_{k=1}^{i}\left(\delta\left(\frac{q^{-j+i-2 k} z_{2}}{z_{1}}\right) T_{i-k}^{\mathrm{PB}}\left(q^{k} z_{1}\right) T_{j+k}^{\mathrm{PB}}\left(q^{-k} z_{2}\right)\right.
\end{aligned}
$$

$$
\begin{align*}
& \left.-\delta\left(\frac{q^{j-i+2 k} z_{2}}{z_{1}}\right) T_{i-k}^{\mathrm{PB}}\left(q^{-k} z_{1}\right) T_{j+k}^{\mathrm{PB}}\left(q^{k} z_{2}\right)\right)+\delta\left(\frac{q^{-2 N+j-i-1} z_{2}}{z_{1}}\right) T_{j-i}^{\mathrm{PB}}\left(q^{-i} z_{2}\right) \\
& -\delta\left(\frac{q^{2 N-j+i+1} z_{2}}{z_{1}}\right) T_{j-i}^{\mathrm{PB}}\left(q^{i} z_{2}\right), \quad 1 \leq i \leq j \leq N . \tag{3.6}
\end{align*}
$$

Here, the structure functions $C_{i, j}(z)$ are given by

$$
\begin{aligned}
& C_{i, j}(z)=\sum_{m \in \mathbb{Z}} \frac{[\operatorname{Min}(\mathrm{i}, \mathrm{j}) \mathrm{m}]_{\mathrm{q}}\left([(\mathrm{~N}+1-\operatorname{Max}(\mathrm{i}, \mathrm{j})) \mathrm{m}]_{\mathrm{q}}-[(\mathrm{N}-\operatorname{Max}(\mathrm{i}, \mathrm{j})) \mathrm{m}]_{\mathrm{q}}\right)}{[(N+1) m]_{q}-[N m]_{q}} z^{m} \\
& 1 \leq i, j \leq N .
\end{aligned}
$$

Corollary 3.6. The currents $T_{i}^{\mathrm{PB}}(z)$ satisfy the duality relations

$$
\begin{equation*}
T_{2 N+1-i}^{\mathrm{PB}}(z)=T_{i}^{\mathrm{PB}}(z), \quad 0 \leq i \leq N . \tag{3.7}
\end{equation*}
$$

## 4 Proof of Theorem 3.2

In this section, we prove Proposition 3.1, Theorem 3.2, and Lemma 3.4.

### 4.1 Proof of Proposition 3.1

Lemma 4.1. The $\Lambda_{i}(z), i \in J_{N}$, satisfy

$$
\begin{align*}
& f_{1,1}\left(\frac{z_{2}}{z_{1}}\right) \Lambda_{i}\left(z_{1}\right) \Lambda_{j}\left(z_{2}\right)=\Delta\left(\frac{x^{-1} z_{2}}{z_{1}}\right): \Lambda_{i}\left(z_{1}\right) \Lambda_{j}\left(z_{2}\right):, \quad i, j \in J_{N}, \quad i \prec j, \quad j \neq \bar{i}, \\
& f_{1,1}\left(\frac{z_{2}}{z_{1}}\right) \Lambda_{j}\left(z_{1}\right) \Lambda_{i}\left(z_{2}\right)=\Delta\left(\frac{x z_{2}}{z_{1}}\right): \Lambda_{j}\left(z_{1}\right) \Lambda_{i}\left(z_{2}\right):, \quad i, j \in J_{N}, \quad i \prec j, \quad j \neq \bar{i}, \\
& f_{1,1}\left(\frac{z_{2}}{z_{1}}\right) \Lambda_{0}\left(z_{1}\right) \Lambda_{0}\left(z_{2}\right)=\Delta\left(\frac{z_{2}}{z_{1}}\right): \Lambda_{0}\left(z_{1}\right) \Lambda_{0}\left(z_{2}\right):, \\
& f_{1,1}\left(\frac{z_{2}}{z_{1}}\right) \Lambda_{i}\left(z_{1}\right) \Lambda_{i}\left(z_{2}\right)=: \Lambda_{i}\left(z_{1}\right) \Lambda_{i}\left(z_{2}\right):, \quad i \in J_{N} \backslash\{0\},  \tag{4.1}\\
& f_{1,1}\left(\frac{z_{2}}{z_{1}}\right) \Lambda_{k}\left(z_{1}\right) \Lambda_{\bar{k}}\left(z_{2}\right)=\Delta\left(\frac{x^{-1} z_{2}}{z_{1}}\right) \Delta\left(\frac{x^{-2 N-2+2 k} z_{2}}{z_{1}}\right): \Lambda_{k}\left(z_{1}\right) \Lambda_{\bar{k}}\left(z_{2}\right):, \quad 1 \leq k \leq N, \\
& f_{1,1}\left(\frac{z_{2}}{z_{1}}\right) \Lambda_{\bar{k}}\left(z_{1}\right) \Lambda_{k}\left(z_{2}\right)=\Delta\left(\frac{x z_{2}}{z_{1}}\right) \Delta\left(\frac{x^{2 N+2-2 k} z_{2}}{z_{1}}\right): \Lambda_{\bar{k}}\left(z_{1}\right) \Lambda_{k}\left(z_{2}\right):, \quad 1 \leq k \leq N .
\end{align*}
$$

Proof. Using (A.2) and (A.8), we obtain the normal ordering rules (4.1).
Lemma 4.2. The $\Lambda_{i}(z), i \in J_{N}$, satisfy

$$
\begin{align*}
& : \Lambda_{0}(z) \Lambda_{0}(x z):=\Delta(1): \Lambda_{N}(z) \Lambda_{\bar{N}}(x z):  \tag{4.2}\\
& : \Lambda_{1}(z) \Lambda_{\overline{1}}\left(x^{2 N+1} z\right):=1  \tag{4.3}\\
& : \Lambda_{k}(z) \Lambda_{\bar{k}}\left(x^{2 N-2 k+3} z\right):=: \Lambda_{k-1}(z) \Lambda_{\overline{k-1}}\left(x^{2 N-2 k+3} z\right): \quad 2 \leq k \leq N \tag{4.4}
\end{align*}
$$

Proof. From (2.5), we obtain (4.2) and (4.4). From (2.2), (2.3) and (2.5), we obtain (4.3).
Lemma 4.3. The $\Delta(z)$ and $f_{i, j}(z)$ satisfy the following fusion relations:

$$
\begin{equation*}
f_{i, j}(z)=f_{j, i}(z)=\prod_{k=1}^{i} f_{1, j}\left(z^{-i-1+2 k} z\right), \quad 1 \leq i \leq j \tag{4.5}
\end{equation*}
$$

$$
\begin{align*}
& f_{1, i}(z)=\left(\prod_{k=1}^{i-1} \Delta\left(x^{-i+2 k} z\right)\right)^{-1} \prod_{k=1}^{i} f_{1,1}\left(x^{-i-1+2 k} z\right), \quad i \geq 2,  \tag{4.6}\\
& f_{i, 2 N+1}(z)=\prod_{k=1}^{i} \Delta\left(x^{-i-1+2 k} z\right), \quad i \geq 1,  \tag{4.7}\\
& f_{i, j}(z)=f_{i, 2 N+1-j}(z)=f_{2 N+1-j, i}(z)=f_{j, i}(z), \quad i \geq 1,1 \leq j \leq N,  \tag{4.8}\\
& f_{1, j}(z) f_{1, j}\left(x^{2 N+1} z\right)=\Delta\left(x^{j} z\right) \Delta\left(x^{2 N+1-j} z\right), \quad j \geq 1,  \tag{4.9}\\
& f_{1, i}(z) f_{j, i}\left(x^{ \pm(j+1)} z\right)= \begin{cases}f_{j+1, i}\left(x^{ \pm j} z\right) \Delta\left(x^{ \pm i} z\right), & 1 \leq i \leq j, \\
f_{j+1, i}\left(x^{ \pm j} z\right), & 1 \leq j<i,\end{cases}  \tag{4.10}\\
& f_{1, i}(z) f_{1, j}\left(x^{ \pm(i+j)} z\right)=f_{1, i+j}\left(x^{ \pm j} z\right) \Delta\left(x^{ \pm i} z\right), \quad i, j \geq 1,  \tag{4.11}\\
& f_{1, i}(z) f_{1, j}\left(x^{ \pm(i-j-2 k)} z\right)=f_{1, i-k}\left(x^{\mp k} z\right) f_{1, j+k}\left(x^{ \pm(i-j-k)} z\right), \quad i, j, i-k, j+k \geq 1 . \tag{4.12}
\end{align*}
$$

Proof. We show (4.6) here. From the definitions, we have

$$
\begin{aligned}
& \left(\prod_{k=1}^{i-1} \Delta_{1}\left(x^{-i+2 k} z\right)\right)^{-1} \prod_{k=1}^{i} f_{1,1}\left(x^{-i-1+2 k} z\right) \\
& \quad=\exp \left(-\sum_{m=1}^{\infty} \frac{1}{m} \frac{[r m]_{x}[(r-1) m]_{x}}{[(N+1) m]_{x}-[N m]_{x}}\left(x-x^{-1}\right)^{2}\left\{\left([N m]_{x}-[(N-1) m]_{x}\right)\right.\right. \\
& \left.\left.\quad \times \sum_{k=1}^{i} x^{(-i+2 k-1) m}-\left([(N+1) m]_{x}-[N m]_{x}\right) \sum_{k=1}^{i-1} x^{(-i+2 k) m}\right\} z^{m}\right) .
\end{aligned}
$$

Using the relation

$$
[(a-1) m]_{x} \sum_{k=1}^{i} x^{(-i+2 k-1) m}-[a m]_{x} \sum_{k=1}^{i-1} x^{(-i+2 k) m}=[(a-i) m]_{x}, \quad a=N, N+1,
$$

we obtain $f_{1, i}(z)$ in the right hand side of the previous formula. We obtain (4.5), (4.7), (4.8), and (4.9) by straightforward calculation from the definitions. Using (4.5) and (4.6), we obtain the relations (4.10), (4.11), and (4.12).

Lemma 4.4. The following relation holds for $A \subset J_{N}$ :

$$
\vec{\Lambda}_{J_{N} \backslash A}(z)=\vec{\Lambda}_{A}(z) \times \begin{cases}\frac{\left[r-\frac{1}{2}\right]_{x}}{\left[\frac{1}{2}\right]_{x}}, & 0 \notin A,  \tag{4.13}\\ \frac{\left[\frac{1}{2}\right]_{x}}{\left[r-\frac{1}{2}\right]_{x}}, & 0 \in A .\end{cases}
$$

Proof. First, we consider the case $A=\varnothing$ and $\overline{J_{N} \backslash A}=J_{N}$. In this case, (4.13) can be rewritten as

$$
\begin{equation*}
: \Lambda_{1}\left(x^{-2 N} z\right) \cdots \Lambda_{N}\left(x^{-2} z\right) \Lambda_{0}(z) \Lambda_{\bar{N}}\left(x^{2} z\right) \cdots \Lambda_{\overline{1}}\left(x^{2 N} z\right):=\frac{\left[r-\frac{1}{2}\right]_{x}}{\left[\frac{1}{2}\right]_{x}} . \tag{4.14}
\end{equation*}
$$

Using (2.2), (2.3), and (2.5), the left side of (4.14) can be written as

$$
\frac{\left[r-\frac{1}{2}\right]_{x}}{\left[\frac{1}{2}\right]_{x}}: \exp \left(\sum_{m \neq 0}\left(\frac{[(2 N+1) m]_{x}}{[m]_{x}} y_{1}(m)-\sum_{j=1}^{N} \frac{[(2 N+1-j) m]_{x}+[j m]_{x}}{[m]_{x}} a_{j}(m)\right) z^{-m}\right): .
$$

Using the relation $\frac{[(2 N+1-j) m]_{x}+[j m]_{x}}{[(2 N+1) m]_{x}}=\frac{[(N+1-j) m]_{x}-[(N-j) m]_{x}}{[(N+1) m]_{x}-[N m]_{x}}$, the generators $y_{1}(m)$ in (A.7) are rewritten as $y_{1}(m)=\sum_{j=1}^{N} \frac{[(2 N+1-j) m]_{x}+[j m]_{x}}{[(2 N+1) m]_{x}} a_{j}(m)$. Hence, we obtain (4.14).

Next, we show (4.13) for $A \subset J_{N}$. Cases $(i), 0 \in A$ and (ii), $0 \notin A$ are proved separately. First, we study case $(i), 0 \in A$. Let

$$
A=\left\{k_{1}, \ldots, k_{K}, 0, \overline{l_{L}}, \ldots, \overline{l_{1}} \mid k_{1} \prec \cdots \prec k_{K} \prec 0 \prec \overline{l_{L}} \prec \cdots \prec \overline{l_{1}}, 1 \leq K, L \leq N\right\}
$$

Multiplying (4.14) by $\overrightarrow{\Lambda_{A}}\left(x^{L-K+1} z\right)$ on the left, and using (4.1) and (4.7) yields

$$
\begin{equation*}
: \vec{\Lambda}_{J_{N}}(z) \vec{\Lambda}_{A}\left(x^{L-K+1} z\right):=\frac{\left[r-\frac{1}{2}\right]_{x}}{\left[\frac{1}{2}\right]_{x}} \vec{\Lambda}_{A}\left(x^{L-K+1} z\right) \tag{4.15}
\end{equation*}
$$

Using (4.2), (4.3) and (4.4) yields

$$
\begin{aligned}
& : \vec{\Lambda}_{J_{N}}(z) \Lambda_{0}(x z):=\Delta(1) \vec{\Lambda}_{J_{N} \backslash\{0\}}(x z) \\
& : \vec{\Lambda}_{J_{N} \backslash\{0\}}(z) \Lambda_{\bar{l}_{L}}\left(x^{2} z\right):=\vec{\Lambda}_{J_{N} \backslash\left\{l_{L}, 0\right\}}(x z) \\
& : \vec{\Lambda}_{J_{N} \backslash\left\{l_{L-s+1}, \ldots, l_{L}, 0\right\}}(z) \Lambda_{\bar{l}_{L-s}}\left(x^{2+s} z\right):=\vec{\Lambda}_{J_{N} \backslash\left\{l_{L-s}, \ldots, l_{L}, 0\right\}}(x z), \quad 1 \leq s \leq L-1 \\
& : \vec{\Lambda}_{J_{N} \backslash\left\{l_{1}, \ldots, l_{L}, 0\right\}}(z) \Lambda_{k_{K}}\left(x^{-L-2} z\right):=\vec{\Lambda}_{J_{N} \backslash\left\{l_{1}, \ldots, l_{L}, 0, \bar{k}_{K}\right\}}\left(x^{-1} z\right) \\
& : \vec{\Lambda}_{J_{N} \backslash\left\{l_{1}, \ldots, l_{L}, 0, \bar{k}_{K}, \ldots, \bar{k}_{K-s+1}\right\}}(z) \Lambda_{k_{K-s}}\left(x^{-L-2-s} z\right):=\vec{\Lambda}_{J_{N} \backslash\left\{l_{1}, \ldots, l_{L}, 0, \bar{k}_{K}, \ldots, \bar{k}_{K-s}\right\}}\left(x^{-1} z\right), \\
& \quad 1 \leq s \leq K-1
\end{aligned}
$$

Using the above five relations yields

$$
: \vec{\Lambda}_{J_{N}}(z) \vec{\Lambda}_{A}\left(x^{L-K+1} z\right):=\Delta(1) \vec{\Lambda}_{\overline{J_{N} \backslash A}}\left(x^{L-K+1} z\right)
$$

From (4.15) we obtain (4.13) for $0 \in A$.
Next, we study case $(i i), 0 \notin A$. The proof for this case is similar to that of case $(i)$. Let

$$
A=\left\{k_{1}, \ldots, k_{K}, \overline{l_{L}}, \ldots, \overline{l_{1}} \mid k_{1} \prec \cdots \prec k_{K} \prec \overline{l_{L}} \prec \cdots \prec \overline{l_{1}}, 1 \leq K, L \leq N\right\}
$$

Multiplying (4.14) by $\overrightarrow{\Lambda_{A}}\left(x^{L-K} z\right)$ on the left, and using (4.1) and (4.7) yields

$$
\begin{equation*}
: \vec{\Lambda}_{J_{N}}(z) \vec{\Lambda}_{A}\left(x^{L-K} z\right):=\frac{\left[r-\frac{1}{2}\right]_{x}}{\left[\frac{1}{2}\right]_{x}} \vec{\Lambda}_{A}\left(x^{L-K} z\right) \tag{4.16}
\end{equation*}
$$

Using (4.2), (4.3), and (4.4) yields

$$
\begin{aligned}
& : \vec{\Lambda}_{J_{N}}(z) \Lambda_{\bar{l}_{L}}(x z):=\vec{\Lambda}_{J_{N} \backslash\left\{l_{L}\right\}}(x z) \\
& : \vec{\Lambda}_{J_{N} \backslash\left\{l_{L-s+1}, \ldots, l_{L}\right\}}(z) \Lambda_{\bar{l}_{L-s}}\left(x^{1+s} z\right):=\vec{\Lambda}_{J_{N} \backslash\left\{l_{L-s}, \ldots, l_{L}\right\}}(x z), \quad 1 \leq s \leq L-1 \\
& : \vec{\Lambda}_{J_{N} \backslash\left\{l_{1}, \ldots, l_{L}\right\}}(z) \Lambda_{k_{K}}\left(x^{-L-1} z\right):=\vec{\Lambda}_{J_{N} \backslash\left\{l_{1}, \ldots, l_{L}, \bar{k}_{K}\right\}}\left(x^{-1} z\right) \\
& : \vec{\Lambda}_{J_{N} \backslash\left\{l_{1}, \ldots, l_{L}, \bar{k}_{K}, \ldots, \bar{k}_{K-s+1}\right\}}(z) \Lambda_{k_{K-s}}\left(x^{-L-1-s} z\right):=\vec{\Lambda}_{J_{N} \backslash\left\{l_{1}, \ldots, l_{L}, \bar{k}_{K}, \ldots, \bar{k}_{K-s}\right\}}\left(x^{-1} z\right), \\
& \quad 1 \leq s \leq K-1
\end{aligned}
$$

Using the above five relations yields

$$
: \vec{\Lambda}_{J_{N}}(z) \vec{\Lambda}_{A}\left(x^{L-K} z\right):=\vec{\Lambda}_{\overline{J_{N} \backslash A}}\left(x^{L-K} z\right)
$$

From (4.16) we obtain (4.13) for $0 \notin A$.

Lemma 4.5. The following relation holds for $A \subset J_{N}$ with $|A| \leq N$ :

$$
\frac{d_{J_{N} \backslash A}(x, r)}{d_{A}(x, r)}=\prod_{k=1}^{N-|A|} \Delta\left(x^{2 k}\right) \times \begin{cases}\Delta(1), & 0 \in A,  \tag{4.17}\\ 1, & 0 \notin A .\end{cases}
$$

Proof. We define the map $\sigma: J_{N} \rightarrow J_{N+1}$ by

$$
\sigma(j)= \begin{cases}k+1, & j=k, \quad 1 \leq k \leq N \\ 0, & j=0, \\ \overline{k+1}, & j=\bar{k}, \quad 1 \leq k \leq N\end{cases}
$$

For $T \subset J_{N}$ with $|T| \leq N$, relation (4.17) is rewritten as

$$
\frac{d_{\sigma\left(J_{N} \backslash T\right)}(x, r)}{d_{\sigma(T)}(x, r)}=\prod_{k=1}^{N-|T|} \Delta\left(x^{2 k}\right) \times \begin{cases}\Delta(1), & 0 \in T \\ 1, & 0 \notin T\end{cases}
$$

Hence, the relation

$$
\frac{d_{\left(J_{N} \backslash B\right) \cap \sigma\left(J_{N}\right)}(x, r)}{d_{B \cap \sigma\left(J_{N}\right)}(x, r)}=\prod_{k=1}^{N-\left|B \cap \sigma\left(J_{N}\right)\right|} \Delta\left(x^{2 k}\right) \times \begin{cases}\Delta(1), & 0 \in B,  \tag{4.18}\\ 1, & 0 \notin B,\end{cases}
$$

for $B \subset J_{N+1}$ with $\left|B \cap \sigma\left(J_{N}\right)\right| \leq N$ holds if relation (4.17) for $A \subset J_{N}$ with $|A| \leq N$ is assumed. Here, we used $\left|B \cap \sigma\left(J_{N}\right)\right|=\left|\sigma^{-1}\left(B \cap \sigma\left(J_{N}\right)\right)\right|$.

We prove (4.17) by induction on $N$. First, we establish the base $N=1$ using case-by-case analysis. For $A=\varnothing$, we obtain $d_{A}(x, r)=1$ and $d_{J_{N} \backslash A}(x, r)=\Delta\left(x^{2}\right)$. For $A=\{1\}$, we obtain $d_{A}(x, r)=1$ and $d_{J_{N} \backslash A}(x, r)=1$. For $A=\{0\}$, we obtain $d_{A}(x, r)=1$ and $d_{J_{N} \backslash A}(x, r)=\Delta(1)$. For $A=\{\overline{1}\}$, we obtain $d_{A}(x, r)=1$ and $d_{J_{N} \backslash A}(x, r)=1$. This implies that (4.17) holds for $N=1$.

Next, we assume that relation (4.17) holds for some $N$, and show (4.17) for $N$ replaced by $N+1$. Let $A \subset J_{N+1}$. From the definition of $d_{A}(x, r)$, we obtain

$$
\begin{align*}
\frac{d_{J_{N} \backslash A}(x, r)}{d_{A}(x, r)}= & \frac{d_{\left(J_{N} \backslash A\right) \cap \sigma\left(J_{N}\right)}(x, r)}{d_{A \cap \sigma\left(J_{N}\right)}(x, r)} \\
& \times \begin{cases}1, & 1 \in A, \quad \overline{1} \notin A \quad \text { or } \quad 1 \notin A, \quad \overline{1} \in A, \\
\Delta\left(x^{2\left(N-\left|J_{N} \backslash A\right|+1\right)}\right)^{-1}, & 1, \\
\Delta\left(x^{2(N-|A|+1)}\right), & 1, \\
\overline{1} \notin A,\end{cases} \tag{4.19}
\end{align*}
$$

Cases $(i), 1 \in A, \overline{1} \notin A$ (or $1 \notin A, \overline{1} \in A$ ), (ii), $1, \overline{1} \in A$, and (iii), $1, \overline{1} \notin A$ are proved separately.

First, we study case $(i), 1 \in A, \overline{1} \notin A$ (or $1 \notin A, \overline{1} \in A$ ). In this case, we obtain $\left|A \cap \sigma\left(J_{N}\right)\right|=$ $|A|-1 \leq N$. Hence, (4.18) holds with $B=A$. Using (4.18), (4.19) and $\left|A \cap \sigma\left(J_{N}\right)\right|=|A|-1$ yields (4.17) with $N$ replaced by $N+1$.

Next, we study case (ii), $1, \overline{1} \in A$. In this case, we obtain $\left|A \cap \sigma\left(J_{N}\right)\right|=|A|-2 \leq N-1$. Hence, (4.18) holds with $B=A$. Using (4.18) and (4.19), $\left|A \cap \sigma\left(J_{N}\right)\right|=|A|-2$ and $\left|J_{N} \backslash A\right|=2 N+3-|A|$ yields (4.17) with $N$ replaced by $N+1$.

Finally, we examine case (iii), $1, \overline{1} \notin A$. Case ( $i$ iii) is further subdivided into (iii.1), $|A| \leq N$, $1, \overline{1} \notin A$ and (iii.2), $|A|=N+1,1, \overline{1} \notin A$.

For the condition (iii.1), we obtain $\left|A \cap \sigma\left(J_{N}\right)\right|=|A| \leq N$. Hence, (4.18) holds with $B=A$. Using (4.18), (4.19), and $\left|A \cap \sigma\left(J_{N}\right)\right|=|A|$ yields (4.17) with $N$ replaced by $N+1$.

For condition (iii.2), we obtain $\left|\left(J_{N} \backslash A\right) \cap \sigma\left(J_{N}\right)\right|=N$. Hence, (4.18) holds with $B=J_{N} \backslash A$. Using (4.18) and (4.19), $|A|=N+1$ and $\left|\left(J_{N} \backslash A\right) \cap \sigma\left(J_{N}\right)\right|=N$ yields (4.17) with $N$ replaced by $N+1$.

Proof. Here we will show Proposition 3.1. Using (4.13), (4.17), and $d \overline{\overline{J_{N} \backslash \Omega_{i}}}(x, r)=d_{J_{N} \backslash \Omega_{i}}(x, r)$ yields

$$
\begin{equation*}
d \overline{J_{N} \backslash \Omega_{i}}(x, r) \vec{\Lambda} \frac{\left[r-\frac{1}{2}\right]_{x}}{\left[\frac{1}{2}\right]_{x}} \prod_{k=1}^{N-\left|\Omega_{i}\right|} \Delta\left(x^{2 k}\right) d_{\Omega_{i}}(x, r) \vec{\Lambda}_{\Omega_{i}}(z) \tag{4.20}
\end{equation*}
$$

Adding relations (4.20) over all $\Omega_{i} \subset J_{N}$ for each fixed $i, 0 \leq i \leq N$, yields (3.3).

### 4.2 Proof of Theorem 3.2

Lemma 4.6. The $W$-currents $T_{j}(z), 1 \leq j \leq N$, satisfy the set of quadratic relations

$$
\begin{align*}
& f_{1, j}\left(\frac{z_{2}}{z_{1}}\right) T_{1}\left(z_{1}\right) T_{j}\left(z_{2}\right)-f_{j, 1}\left(\frac{z_{1}}{z_{2}}\right) T_{j}\left(z_{2}\right) T_{1}\left(z_{1}\right) \\
&= c(x, r)\left(\delta\left(\frac{x^{-j-1} z_{2}}{z_{1}}\right) T_{j+1}\left(x^{-1} z_{2}\right)-\delta\left(\frac{x^{j+1} z_{2}}{z_{1}}\right) T_{j+1}\left(x z_{2}\right)\right)+c(x, r) \Delta\left(x^{2 N+2-2 j}\right) \\
& \times\left(\delta\left(\frac{x^{-2 N+j-2} z_{2}}{z_{1}}\right) T_{j-1}\left(x^{-1} z_{2}\right)-\delta\left(\frac{x^{2 N-j+2} z_{2}}{z_{1}}\right) T_{j-1}\left(x z_{2}\right)\right), \quad 1 \leq j \leq N . \tag{4.21}
\end{align*}
$$

Here, we use $f_{i, j}(z)$ introduced in (3.1).
Proof. In this proof, we frequently use exchange relations (B.1)-(B.7) in Appendix B. We start from

$$
\mathrm{LHS}_{1, j}=f_{1, j}\left(z_{2} / z_{1}\right) T_{1}\left(z_{1}\right) T_{j}\left(z_{2}\right)-f_{j, 1}\left(z_{1} / z_{2}\right) T_{j}\left(z_{2}\right) T_{1}\left(z_{1}\right), \quad 1 \leq j \leq N
$$

From the definition of $T_{j}(z)$ introduced in $(3.2), \mathrm{LHS}_{1, j}$ can be written as the sum of

$$
f_{1, j}\left(z_{2} / z_{1}\right) \Lambda_{s}\left(z_{1}\right) \vec{\Lambda}_{\Omega_{j}}\left(z_{2}\right)-f_{j, 1}\left(z_{1} / z_{2}\right) \vec{\Lambda}_{\Omega_{j}}\left(z_{2}\right) \Lambda_{s}\left(z_{1}\right) \quad \text { over } \quad s \in J_{N}, \Omega_{j} \subset J_{N},\left|\Omega_{j}\right|=j
$$

summarized in Appendix B. Adding exchange relations (B.1)-(B.7) over $s \in J_{N}, \Omega_{j} \subset J_{N}$, $\left|\Omega_{j}\right|=j$ yields

$$
\begin{align*}
\mathrm{LHS}_{1, j}=c(x, r)\{ & \sum_{m=0}^{\left[\frac{j}{2}\right]}\left(\delta\left(x^{-j-1+2 m} \frac{z_{2}}{z_{1}}\right) \bar{G}_{j+1-2 m}\left(z_{2}\right)-\delta\left(x^{j+1-2 m} \frac{z_{2}}{z_{1}}\right) G_{j+1-2 m}\left(z_{2}\right)\right) \\
& +\sum_{m=0}^{N-\left[\frac{j-1}{2}\right]}\left(\delta\left(x^{-2 N+j-2+2 m} \frac{z_{2}}{z_{1}}\right) \bar{H}_{2 N-j+2-2 m}\left(z_{2}\right)\right. \\
& \left.\left.-\delta\left(x^{2 N-j+2-2 m} \frac{z_{2}}{z_{1}}\right) H_{2 N-j+2-2 m}\left(z_{2}\right)\right)\right\} . \tag{4.22}
\end{align*}
$$

Formulas for $\quad \bar{G}_{j+1}(z), \quad G_{j+1}(z), \quad \bar{H}_{2 N-j+2}(z), \quad H_{2 N-j+2}(z), \quad \bar{G}_{j+1-2 m}(z), \quad G_{j+1-2 m}(z)$, $\bar{H}_{2 N-j+2-2 m}(z)$, and $H_{2 N-j+2-2 m}(z)$ will be given below. In (4.38) we define $H_{0}(z)=0$ to avoid ambiguity of $\bar{H}_{0}(z)$ and $H_{0}(z)$. In the case when $j$ is even, we have $\mathrm{LHS}_{1, j}=$ $c(x, r) \bar{H}_{0}\left(z_{2}\right)-H_{0}\left(z_{2}\right) \delta\left(z_{2} / z_{1}\right)+\bar{H}_{2}\left(z_{2}\right) \delta\left(x^{-2} z_{2} / z_{1}\right)-H_{2}\left(z_{2}\right) \delta\left(x^{2} z_{2} / z_{1}\right)+\cdots$.

First, we define $\bar{G}_{j+1}(z), 1 \leq j \leq N$, as the coefficient of $\delta\left(x^{-j-1} z_{2} / z_{1}\right)$ in (4.22). In what follows, for a subset $\Omega_{j} \subset J_{N}$ with $\left|\Omega_{j}\right|=j$, we write its elements as $s_{1}, s_{2}, \ldots, s_{j}, s_{1} \prec s_{2} \prec$ $\cdots \prec s_{j}$. Adding the first term in (B.1) and the first term in (B.6) yields

$$
\begin{aligned}
\bar{G}_{j+1}(z)= & \sum_{\Omega_{j} \subset J_{N}} \sum_{\substack{s \in J_{N} \\
s \prec s_{1}, \bar{s} \notin \Omega_{j}}} d_{\Omega_{j}}(x, r): \Lambda_{s}\left(x^{-j-1} z\right) \vec{\Lambda}_{\Omega_{j}}(z): \\
& +\sum_{\Omega_{j} \subset J_{N}} \sum_{\substack{n=1 \\
n \prec s_{1}}}^{N} \sum_{\substack{l=1 \\
s_{l}=\bar{n}}}^{j} \Delta\left(x^{2(N+1-l-n)}\right) d_{\Omega_{j}}(x, r): \Lambda_{n}\left(x^{-j-1} z\right) \vec{\Lambda}_{\Omega_{j}}(z): .
\end{aligned}
$$

Using : $\Lambda_{n}\left(x^{-j-1} z\right) \vec{\Lambda}_{\Omega_{j}}(z):=\vec{\Lambda}_{\Omega_{j} \cup\{n\}}\left(x^{-1} z\right)$ and

$$
d_{\Omega_{j} \cup\{n\}}(x, r)=d_{\Omega_{j}}(x, r) \times\left\{\begin{array}{ll}
1, & \bar{n} \notin \Omega_{j}, \\
\Delta\left(x^{2(N+1-l-n)}\right), & \bar{n}=s_{l}
\end{array} \quad \text { with } n \prec s_{1}, 1 \leq n \leq N,\right.
$$

yields

$$
\bar{G}_{j+1}(z)=\sum_{\Omega_{j} \subset J_{N}} \sum_{\substack{s \in J_{N} \\ s<s_{1}}} d_{\Omega_{j} \cup\{s\}}(x, r) \vec{\Lambda}_{\Omega_{j} \cup\{s\}}\left(x^{-1} z\right) .
$$

Hence, we obtain $\bar{G}_{j+1}(z)=T_{j+1}\left(x^{-1} z\right), 1 \leq j \leq N$.
Next, we define $G_{j+1}(z), 1 \leq j \leq N$, as the coefficient of $\delta\left(x^{j+1} z_{2} / z_{1}\right)$ in (4.22). Adding the second term in (B.1) and the third term in (B.7) yields

$$
\begin{aligned}
G_{j+1}(z)= & \sum_{\Omega_{j} \subset J_{N}} \sum_{\substack{s \in J_{N} \\
s_{j}<s, \bar{s} \notin \Omega_{j}}} d_{\Omega_{j}}(x, r): \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{s}\left(x^{j+1} z\right): \\
& +\sum_{\Omega_{j} \subset J_{N}} \sum_{\substack{n=1 \\
s_{j}<\bar{n} \\
s_{k}=n}}^{N} \sum_{\substack{k=1 \\
s_{k}=n}}^{j} \Delta\left(x^{2(N+k-j-n)}\right) d_{\Omega_{j}}(x, r): \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{\bar{n}}\left(x^{j+1} z\right): .
\end{aligned}
$$

Using $: \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{\bar{n}}\left(x^{j+1} z\right):=\vec{\Lambda}_{\Omega_{j} \cup\{\bar{n}\}}(x z)$ and

$$
d_{\Omega_{j} \cup\{\bar{n}\}}(x, r)=d_{\Omega_{j}}(x, r) \times\left\{\begin{array}{ll}
1, & n \notin \Omega_{j}, \\
\Delta\left(x^{2(N+k-j-n)}\right), & n=s_{k}
\end{array} \quad \text { with } s_{j} \prec \bar{n}, 1 \leq n \leq N,\right.
$$

yields

$$
G_{j+1}(z)=\sum_{\Omega_{j} \subset J_{N}} \sum_{\substack{s \in J_{N} \\ s_{j} \not s s}} d_{\Omega_{j} \cup\{s\}}(x, r) \vec{\Lambda}_{\Omega_{j} \cup\{s\}}(x z) .
$$

Hence, we obtain $G_{j+1}(z)=T_{j+1}(z), 1 \leq j \leq N$.
We define $\bar{H}_{2 N-j+2}(z), 1 \leq j \leq N$, as the coefficient of $\delta\left(x^{-2 N+j-2} z_{2} / z_{1}\right)$ in (4.22). Adding the first term in (B.3), the second term in (B.3), the second term in (B.6), and the fourth term in (B.6) yields

$$
\begin{align*}
& \bar{H}_{2 N-j+2}(z)=\sum_{n=1}^{j-1} \sum_{k=1}^{j-n} \sum_{\substack{\Omega_{j} C_{N} \\
s_{k}=n \\
s_{l}=\bar{n}, l=j-n+1}} d_{\Omega_{j}}(x, r): \Lambda_{n}\left(x^{-2 N+j-2} z\right) \vec{\Lambda}_{\Omega_{j}}(z): \\
& -\sum_{n=1}^{j-2} \sum_{k=1}^{j-n-1} \sum_{\substack{\Omega_{j} \subset J_{N} \\
s_{k}=n \\
s_{l}=\bar{n}, l=j-n}} d_{\Omega_{j}}(x, r): \Lambda_{n}\left(x^{-2 N+j-2} z\right) \vec{\Lambda}_{\Omega_{j}}(z): \\
& +\sum_{n=1}^{j} \sum_{k=1}^{j-n+1} \sum_{\substack{\Omega_{j} \subset J_{N} \\
s_{k-1}<n<s_{k} \\
s_{l}=\bar{n}, l=j+1-n}} \Delta\left(x^{2(N-j+k)}\right) d_{\Omega_{j}}(x, r): \Lambda_{n}\left(x^{-2 N+j-2} z\right) \vec{\Lambda}_{\Omega_{j}}(z): \\
& -\sum_{n=1}^{j-1} \sum_{k=1}^{j-n} \sum_{\substack{\Omega_{j} \subset J_{N} \\
s_{k}=1 \\
s_{l}=\bar{n}, l=j-s_{k}}} \Delta\left(x^{2(N-j+k)}\right) d_{\Omega_{j}}(x, r): \Lambda_{n}\left(x^{-2 N+j-2} z\right) \vec{\Lambda}_{\Omega_{j}}(z): . \tag{4.23}
\end{align*}
$$

The second term in (4.23) vanishes, because there doesn't exist $s_{j} \in J_{N}$, if $\Omega_{j} \subset J_{N}, s_{k}=n$, $1 \leq k \leq j+1-n, s_{l}=\bar{n}, l=j-n$, and $1 \leq n \leq j-2$ are satisfied. The fourth term in (4.23) vanishes, because there doesn't exist $s_{j} \in J_{N}$, if $\Omega_{j} \subset J_{N}, s_{k-1} \prec n \prec s_{k}, 1 \leq k \leq j-n, s_{l}=\bar{n}$, $l=j-n$, and $1 \leq n \leq j-1$ are satisfied. Rewriting the sum of the first and the third terms yields

$$
\begin{aligned}
\bar{H}_{2 N-j+2}(z)= & \sum_{n=1}^{j-1} \sum_{k=1}^{\operatorname{Min}(j-n, n)} \sum_{\substack{\Omega_{j} \subset J_{N}, s_{k}=n \\
\left(s_{j}-n+1, \ldots, s_{j-1}, s_{j}\right) \\
=(\bar{n}, \ldots, 2, \overline{1})}} d_{\Omega_{j}}(x, r): \Lambda_{n}\left(x^{-2 N+j-2} z\right) \vec{\Lambda}_{\Omega_{j}}(z): \\
& +\sum_{n=1}^{j \operatorname{Min}(j+1-n, n)} \sum_{k=1} \sum_{\substack{\Omega_{j} \subset J_{N}, s_{k-1} \prec n \prec s_{k} \\
\left(s_{j-n}+1, \ldots, s_{j}, 1, s_{j}\right) \\
=(\bar{n}, \ldots, 2, \overline{1})}} \Delta\left(x^{2(N-j+k)}\right) d_{\Omega_{j}}(x, r): \Lambda_{n}\left(x^{-2 N+j-2} z\right) \vec{\Lambda}_{\Omega_{j}}(z): .
\end{aligned}
$$

The relation $d_{\Omega_{j}}(x, r)=\Delta\left(x^{2(N+1-j)}\right) d_{\Omega_{j} \backslash\{\bar{n}\}}(x, r)$ holds, if $s_{k}=n, s_{l}=\bar{n}, 1 \leq n \leq j-1$, and $1 \leq k<l=j+1-n$ are satisfied. The relation $d_{\Omega_{j}}(x, r) \Delta\left(x^{2(N-j+k)}\right)=\Delta\left(x^{2(N+1-j)}\right) d_{\Omega_{j} \backslash\{\bar{n}\}}(x, r)$ holds, if $s_{k-1} \prec n \prec s_{k}, s_{l}=\bar{n}, 1 \leq n \leq j$, and $1 \leq k \leq l=j+1-n$ are satisfied. Using the above two relations and

$$
\begin{aligned}
& : \Lambda_{n}\left(x^{-2 N+j-2} z\right) \vec{\Lambda}_{\Omega_{j}}(z):=\vec{\Lambda}_{\Omega_{j} \backslash\{\bar{n}\}}\left(x^{-1} z\right) \\
& \left(s_{j+1-n}, \ldots, s_{j-1}, s_{j}\right)=(\bar{n}, \ldots, \overline{2}, \overline{1}), \quad 1 \leq n \leq j
\end{aligned}
$$

obtained from (4.3) and (4.4), yields

$$
\begin{aligned}
\bar{H}_{2 N-j+2}(z)= & \Delta\left(x^{2(N-j+1)}\right)\left(\sum_{n=1}^{j-1} \sum_{k=1}^{\operatorname{Max}(j-n, n)} \sum_{\substack{\Omega_{j} \subset J_{N} \\
s_{k}=n, s_{l}=\bar{n}, l=j+1-n}} d_{\Omega_{j} \backslash\left\{s_{l}\right\}}(x, r) \vec{\Lambda}_{\Omega_{j} \backslash\left\{s_{l}\right\}}\left(x^{-1} z\right)\right. \\
& \left.+\sum_{n=1}^{j} \sum_{k=1}^{\operatorname{Max}(j+1-n, n)} \sum_{\substack{\Omega_{j} \subset J_{N} \\
s_{k}<-1<s_{k} \\
s_{l}=\bar{n}, l=j+1-n}} d_{\Omega_{j} \backslash\left\{s_{l}\right\}}(x, r) \vec{\Lambda}_{\Omega_{j} \backslash\left\{s_{l}\right\}}\left(x^{-1} z\right)\right) .
\end{aligned}
$$

Hence, we obtain $\bar{H}_{2 N-j+2}(z)=\Delta\left(x^{2(N-j+1)}\right) T_{j-1}\left(x^{-1} z\right), 1 \leq j \leq N$.
We define $H_{2 N-j+2}(z), 1 \leq j \leq N$, as the coefficient of $\delta\left(x^{2 N-j+2} z_{2} / z_{1}\right)$ in (4.22). Adding the first term in (B.4), the second term in (B.4), the second term in (B.7), and the fourth term in (B.7) yields

$$
\begin{align*}
H_{2 N-j+2}(z)= & -\sum_{n=1}^{j-2} \sum_{l=n+2}^{j} \sum_{\substack{\Omega_{j} \subset J_{N} \\
s_{k}=n, k=n+1 \\
s_{l}=\bar{n}}} d_{\Omega_{j}}(x, r): \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{\bar{n}}\left(x^{2 N-j+2}\right): \\
& +\sum_{n=1}^{j-1} \sum_{l=n+1}^{j} \sum_{\substack{\Omega_{j} \subset J_{N} \\
s_{k}=n, k=n \\
s_{l}=\bar{n}}} d_{\Omega_{j}}(x, r): \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{\bar{n}}\left(x^{2 N-j+2} z\right): \\
& +\sum_{n=1}^{j} \sum_{l=n}^{j} \sum_{\substack{\Omega_{j} \subset J_{N} \\
s_{k}=n, k=n \\
s_{l}<\bar{n}<s_{l+1} \\
j}} \Delta\left(x^{2(N+1-l)}\right) d_{\Omega_{j}}(x, r): \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{\bar{n}}\left(x^{2 N-j+2} z\right): \\
& -\sum_{n=1}^{j-1} \sum_{l=n+1}^{j} \sum_{\substack{\Omega_{j} \subset J_{N} \\
s_{k}=n=k=n+1 \\
s_{l}<n<s_{l+1}}} \Delta\left(x^{2(N+1-l)}\right) d_{\Omega_{j}}(x, r): \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{\bar{n}}\left(x^{2 N-j+2} z\right): \tag{4.24}
\end{align*}
$$

The first term in (4.24) vanishes, because there doesn't exist $s_{1} \in J_{N}$, if $\Omega_{j} \subset J_{N}, s_{k}=n$, $k=n+1, s_{l}=\bar{n}, n+2 \leq l \leq j$, and $1 \leq n \leq j-2$ are satisfied. The fourth term in (4.24) vanishes, because there doesn't exist $s_{1} \in J_{N}$, if $\Omega_{j} \subset J_{N}, s_{k}=n, k=n+1, s_{l} \prec \bar{n} \prec s_{l+1}, n+1 \leq l \leq j$, and $1 \leq n \leq j-1$ are satisfied. Rewriting the sum of the second and the third terms yields

$$
\begin{aligned}
H_{2 N-j+2}(z)= & \sum_{n=1}^{j-1} \sum_{l=\operatorname{Max}(n+1, j+1-n)}^{j} \sum_{\substack{\Omega_{j} \subset J_{N} \\
\left(s_{1}, s_{2}, \ldots, s_{n}\right) \\
=(1,2, \ldots, n) \\
s_{l}=\bar{n}}} d_{\Omega_{j}}(x, r): \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{\bar{n}}\left(x^{2 N-j+2} z\right): \\
& +\sum_{n=1}^{j} \sum_{l=\operatorname{Max}(n, j+1-n)}^{j} \sum_{\substack{\Omega_{j} \subset J_{N} \\
\left(s_{1}, s_{2}, \ldots, s_{n}\right) \\
=(1,2, \ldots, n) \\
s_{l} \prec n<s_{l+1}}} \Delta\left(x^{2(N+1-l) d_{\Omega_{j}}(x, r): \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{\bar{n}}\left(x^{2 N-j+2} z\right): .}\right.
\end{aligned}
$$

The relation $d_{\Omega_{j}}(x, r)=\Delta\left(x^{2(N+1-j)}\right) d_{\Omega_{j} \backslash\{n\}}(x, r)$ holds, if $s_{k}=n, s_{l}=\bar{n}, 1 \leq n \leq j-1$, and $k=n$, and $n+1 \leq l \leq j$ are satisfied. The relation $d_{\Omega_{j}}(x, r) \Delta\left(x^{2(N+1-l)}\right)=\Delta\left(x^{2(N+1-j)}\right) \times$ $d_{\Omega_{j} \backslash\{n\}}(x, r)$ holds, if $s_{k}=n, s_{l} \prec \bar{n} \prec s_{l+1}, 1 \leq n \leq j, k=n$, and $n+1 \leq l \leq j$ are satisfied. Using the above two relations and

$$
: \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{\bar{n}}\left(x^{2 N-j+2} z\right):=\vec{\Lambda}_{\Omega_{j} \backslash\{n\}}(x z), \quad\left(s_{1}, s_{2}, \ldots, s_{n}\right)=(1,2, \ldots, n), \quad 1 \leq n \leq j
$$

obtained from (4.3) and (4.4), yields

$$
\begin{aligned}
H_{2 N-j+2}(z)=\Delta\left(x^{2(N-j+1)}\right)( & \sum_{n=1}^{j-1} \sum_{l=\operatorname{Max}(n+1, j+1-n)}^{j} \sum_{\substack{\Omega_{j} \subset J_{N} \\
s_{k}=n, k=n \\
s_{l}=\bar{n}}} d_{\Omega_{j} \backslash\left\{s_{k}\right\}}(x, r) \vec{\Lambda}_{\Omega_{j} \backslash\left\{s_{k}\right\}}(x z) \\
& \left.+\sum_{n=1}^{j} \sum_{l=\operatorname{Max}(n, j+1-n)}^{j} \sum_{\substack{\Omega_{j} \subset J_{N} \\
s_{k}=n, k=n \\
s_{l}<n<s_{l}}}^{j} d_{\Omega_{j} \backslash\left\{s_{k}\right\}}(x, r) \vec{\Lambda}_{\Omega_{j} \backslash\left\{s_{k}\right\}}(x z)\right) .
\end{aligned}
$$

Hence, we obtain $H_{2 N-j+2}(z)=\Delta\left(x^{2(N-j+1)}\right) T_{j-1}(x z), 1 \leq j \leq N$.
We define $\bar{G}_{j+1-2 m}(z), 1 \leq m \leq\left[\frac{j}{2}\right], 1 \leq j \leq N$, as the coefficient of $\delta\left(x^{-j-1+2 m} z_{2} / z_{1}\right)$ in (4.22). Adding the first term in (B.1), the second term in (B.1), the first term in (B.6), the third term in (B.6), the first term in (B.7), and the third term in (B.7) yields

$$
\begin{aligned}
& \bar{G}_{j+1-2 m}(z)=\sum_{\Omega_{j} \subset J_{N}} \sum_{\substack { s \in J_{N} \\
s m \\
\begin{subarray}{c}{\alpha \\
s \neq s_{m+1} \\
s \notin \Omega_{j}{ s \in J _ { N } \\
s m \\
\begin{subarray} { c } { \alpha \\
s \neq s _ { m + 1 } \\
s \notin \Omega _ { j } } }\end{subarray}} d_{\Omega_{j}}(x, r): \Lambda_{s}\left(x^{-j-1+2 m} z\right) \vec{\Lambda}_{\Omega_{j}}(z): \\
& -\sum_{\Omega_{j} \subset J_{N}} \sum_{\substack{s \in J_{N} \\
s_{m}<s_{m} \\
\bar{s} \notin \Omega_{j}}} d_{\Omega_{m}}(x, r): \Lambda_{s}\left(x^{-j-1+2 m} z\right) \vec{\Lambda}_{\Omega_{j}}(z): \\
& +\sum_{\Omega_{j} \subset J_{N}} \sum_{\substack{n=1}}^{N} \sum_{\substack{l=m+1 \\
s_{m} \prec s \prec s_{m+1} \\
s_{l}=\bar{s}, s=n}}^{j} \Delta\left(x^{2(l-m+n-N-1)}\right) d_{\Omega_{j}}(x, r): \Lambda_{s}\left(x^{-j-1+2 m} z\right) \vec{\Lambda}_{\Omega_{j}}(z): \\
& -\sum_{\Omega_{j} \subset J_{N}} \sum_{n=1}^{N} \sum_{\substack{l=m \\
s_{m-1}-1 \prec \prec s s_{m} \\
s_{l}=\bar{s}, s=n}}^{j} \Delta\left(x^{2(l-m+n-N-1)}\right) d_{\Omega_{j}}(x, r): \Lambda_{s}\left(x^{-j-1+2 m} z\right) \vec{\Lambda}_{\Omega_{j}}(z):
\end{aligned}
$$

$$
\begin{aligned}
& +\sum_{\Omega_{j} \subset J_{N}} \sum_{n=1}^{N} \sum_{\substack{k=1 \\
s_{k}=\bar{s}, s=\bar{n} \\
s_{m} \prec s \prec s_{m+1}}}^{m} \Delta\left(x^{2(m-k+n-N-1)}\right) d_{\Omega_{j}}(x, r): \Lambda_{s}\left(x^{-j-1+2 m} z\right) \vec{\Lambda}_{\Omega_{j}}(z): \\
& -\sum_{\Omega_{j} \subset J_{N}} \sum_{n=1}^{N} \sum_{\substack{k=1 \\
s_{k}=\bar{s}, s=\bar{n} \\
s_{m-1} \prec s \prec s_{m}}}^{m-1} \Delta\left(x^{2(m-k+n-N-1)}\right) d_{\Omega_{j}}(x, r): \Lambda_{s}\left(x^{-j-1+2 m} z\right) \vec{\Lambda}_{\Omega_{j}}(z):
\end{aligned}
$$

For a subset $\Omega_{j}=\left\{s_{1}, s_{2}, \ldots, s_{j}\right\} \subset J_{N}$ and an element $s \notin \Omega_{j}, s \in J_{N}$, we write elements of $\Omega_{j} \cup\{s\}$ as $t_{1}, t_{2}, \ldots, t_{j+1}, t_{1} \prec t_{2} \prec \cdots \prec t_{j+1}$. In what follows, we use the abbreviation $\Omega_{j+1}=\left\{t_{1}, t_{2}, \ldots, t_{j+1}\right\}$. Rewriting the sum yields

$$
\begin{aligned}
& \bar{G}_{j+1-2 m}(z)=\sum_{\substack{\Omega_{j+1} \subset J_{N} \\
\bar{t}_{m+1} \notin \Omega_{j+1} \backslash\left\{t_{m+1}\right\}}} d_{\left\{t_{1}, \ldots, t_{m}\right\}}(x, r) d_{\left\{t_{m+2}, \ldots, t_{j+1}\right\}}(x, r) \\
& \times \prod_{p=1}^{m} \prod_{\substack{q=m+2 \\
t_{q}=\bar{t}_{p}}}^{j+1} \Delta\left(x^{2\left(q-p+t_{p}-N-2\right)}\right): \Lambda_{t_{m+1}}\left(x^{-j-1+2 m} z\right) \vec{\Lambda}_{\Omega_{j+1} \backslash\left\{t_{m+1}\right\}}(z): \\
& -\sum_{\substack{\Omega_{j+1} \subset J_{N} \\
\bar{t}_{m} \notin \Omega_{j+1} \backslash\left\{t_{m}\right\}}} d_{\left\{t_{1}, \ldots, t_{m-1}\right\}}(x, r) d_{\left\{t_{m+1}, \ldots, t_{j+1}\right\}}(x, r) \\
& \times \prod_{p=1}^{m-1} \prod_{\substack{q=m+1 \\
t_{q}=\bar{t}_{p}}}^{j+1} \Delta\left(x^{2\left(q-p+t_{p}-N-2\right)}\right): \Lambda_{t_{m}}\left(x^{-j-1+2 m} z\right) \vec{\Lambda}_{\Omega_{j+1} \backslash\left\{t_{m}\right\}}(z): \\
& +\sum_{\Omega_{j+1} \subset J_{N}} \sum_{\substack{n=1 \\
t_{m+1}=n}}^{N} \sum_{\substack{l=m+2 \\
t_{l}=\bar{t}_{m+1}}}^{j+1} d_{\left\{t_{1}, \ldots, t_{m}\right\}}(x, r) d_{\left\{t_{m+2}, \ldots, t_{j+1}\right\}}(x, r) \\
& \times \prod_{p=1}^{m} \prod_{\substack{q=m+2 \\
t_{q}=\bar{t}_{p}}}^{j+1} \Delta\left(x^{2\left(q-p+t_{p}-N-2\right)}\right) \Delta\left(x^{2\left(l-m+t_{m+1}-N-2\right)}\right) \\
& \times: \Lambda_{t_{m+1}}\left(x^{-j-1+2 m} z\right) \vec{\Lambda}_{\Omega_{j+1} \backslash\left\{t_{m+1}\right\}}(z): \\
& -\sum_{\Omega_{j+1} \subset J_{N}} \sum_{\substack{n=1 \\
t_{m}=n}}^{N} \sum_{\substack{l=m+1 \\
t_{l}=\bar{t}_{m}}}^{j+1} d_{\left\{t_{1}, \ldots, t_{m-1}\right\}}(x, r) d_{\left\{t_{m+1}, \ldots, t_{j+1}\right\}}(x, r) \\
& \times \prod_{p=1}^{m-1} \prod_{\substack{q=m+1 \\
t_{q}=\bar{t}_{p}}}^{j+1} \Delta\left(x^{2\left(q-p+t_{p}-N-2\right)}\right) \Delta\left(x^{2\left(l-m+t_{m}-N-2\right)}\right) \\
& \times: \Lambda_{t_{m}}\left(x^{-j-1+2 m} z\right) \vec{\Lambda}_{\Omega_{j+1} \backslash\left\{t_{m}\right\}}(z): \\
& +\sum_{\Omega_{j+1} \subset J_{N}} \sum_{\substack{n=1 \\
t_{m+1}=\bar{n}}}^{N} \sum_{\substack{k=1 \\
t_{m+1}=\bar{t}_{k}}}^{m} d_{\left\{t_{1}, \ldots, t_{m}\right\}}(x, r) d_{\left\{t_{m+2}, \ldots, t_{j+1}\right\}}(x, r) \\
& \times \prod_{p=1}^{m} \prod_{\substack{q=m+2 \\
t_{q}=\bar{t}_{p}}}^{j+1} \Delta\left(x^{2\left(q-p+t_{p}-N-2\right)}\right) \Delta\left(x^{2\left(m-k+t_{k}-N-1\right)}\right)
\end{aligned}
$$

$$
\begin{aligned}
& \times: \Lambda_{t_{m+1}}\left(x^{-j-1+2 m} z\right) \vec{\Lambda}_{\Omega_{j+1} \backslash\left\{t_{m+1}\right\}}(z): \\
& -\sum_{\Omega_{j+1} \subset J_{N}} \sum_{n=1}^{N} \sum_{\substack{k=1 \\
t_{m=n}}}^{m-1} d_{\left\{t_{1}, \ldots, t_{m-1}\right\}}(x, r) d_{\left\{t_{m+1}, \ldots, t_{j+1}\right\}}(x, r) \\
& \times \prod_{p=1}^{m-1} \prod_{\substack{q=m+1 \\
t_{q}=\bar{t}_{k}}}^{j+1} \Delta\left(x^{2\left(q-p+t_{p}-N-2\right)}\right) \Delta\left(x^{2\left(m-k+t_{k}-N-1\right)}\right) \\
& \times: \Lambda_{t_{m}}\left(x^{-j-1+2 m} z\right) \vec{\Lambda}_{\Omega_{j+1} \backslash\left\{t_{m}\right\}}(z):
\end{aligned}
$$

Rewriting the sum yields

$$
\begin{aligned}
\bar{G}_{j+1-2 m}(z)= & \sum_{\Omega_{j+1} \subset J_{N}} d_{\left\{t_{1}, \ldots, t_{m}\right\}}(x, r) d_{\left\{t_{m+2}, \ldots, t_{j+1}\right\}}(x, r) \prod_{p=1}^{m} \prod_{\substack{q=m+2 \\
t_{q}=\bar{t}_{p}}}^{j+1} \Delta\left(x^{2\left(q-p+t_{p}-N-2\right)}\right) \\
& \times \prod_{\substack{p=1 \\
t_{m+1}=\bar{t}_{p}}}^{m} \Delta\left(x^{2\left(m-p+t_{p}-N-1\right)}\right) \prod_{\substack{q=m+2 \\
t_{q}=\bar{t}_{m+1}}}^{j+1} \Delta\left(x^{2\left(q-m+t_{m+1}-N-2\right)}\right) \\
& \times: \Lambda_{t_{m+1}}\left(x^{-j-1+2 m} z\right) \vec{\Lambda}_{\Omega_{j+1} \backslash\left\{t_{m+1}\right\}}(z): \\
& -\sum_{\Omega_{j+1} \subset J_{N}} d_{\left\{t_{1}, \ldots, t_{m-1}\right\}}(x, r) d_{\left\{t_{m+1}, \ldots, t_{j+1}\right\}}(x, r) \prod_{p=1}^{m-1} \prod_{q=m+1}^{j+1} \Delta\left(x^{2\left(q-p+t_{p}-N-2\right)}\right) \\
& \times \prod_{p=1}^{m} \Delta\left(x^{2\left(m-p+t_{p}-N-1\right)}\right) \prod_{\substack{q=\bar{t}_{p}}}^{j+1} \Delta\left(x^{2\left(q-m+t_{m}-N-2\right)}\right) \\
& \times: \Lambda_{t_{m}}\left(x^{-j-1+2 m} z\right) \vec{\Lambda}_{\Omega_{j+1} \backslash\left\{t_{m}\right\}}(z): .
\end{aligned}
$$

Using

$$
\begin{aligned}
& d_{\left\{t_{1}, \ldots, t_{m}\right\}}(x, r) d_{\left\{t_{m+2}, \ldots, t_{j+1}\right\}}(x, r) \prod_{p=1}^{m} \prod_{\substack{q=m+2 \\
t_{q}=\bar{t}_{p}}}^{j+1} \Delta\left(x^{2\left(q-p+t_{p}-N-2\right)}\right) \prod_{\substack{p=1 \\
t_{m+1}=\bar{t}_{p}}}^{m} \Delta\left(x^{2\left(m-p+t_{p}-N-1\right)}\right) \\
& \quad \times \prod_{\substack{q=m+2 \\
t_{q}=\bar{t}_{m+1}}}^{j+1} \Delta\left(x^{2\left(q-m+t_{m+1}-N-2\right)}\right)=d_{\left\{t_{1}, \ldots, t_{m-1}\right\}}(x, r) d_{\left\{t_{m+1}, \ldots, t_{j+1}\right\}}(x, r) \\
& \quad \times \prod_{p=1}^{m-1} \prod_{\substack{q=m+1 \\
t_{q}=\bar{t}_{p}}}^{j+1} \Delta\left(x^{2\left(q-p+t_{p}-N-2\right)}\right) \prod_{\substack{p=1 \\
t_{m}=\bar{t}_{p}}}^{m} \Delta\left(x^{2\left(m-p+t_{p}-N-1\right)}\right) \prod_{\substack{q=m+1 \\
t_{q}=\bar{t}_{m}}}^{j+1} \Delta\left(x^{2\left(q-m+t_{m}-N-2\right)}\right)
\end{aligned}
$$

and

$$
: \Lambda_{t_{m+1}}\left(x^{-j-1+2 m} z\right) \vec{\Lambda}_{\Omega_{j+1} \backslash\left\{t_{m+1}\right\}}(z):=: \Lambda_{t_{m}}\left(x^{-j-1+2 m} z\right) \vec{\Lambda}_{\Omega_{j+1} \backslash\left\{t_{m}\right\}}(z):
$$

yields $\bar{G}_{j+1-2 m}(z)=0,1 \leq m \leq\left[\frac{j}{2}\right], 1 \leq j \leq N$.
We define $G_{j+1-2 m}(z), 1 \leq m \leq\left[\frac{j}{2}\right], 1 \leq j \leq N$, as the coefficient of $\delta\left(x^{j+1-2 m} z_{2} / z_{1}\right)$ in (4.22). Adding the first term in (B.1), the second term in (B.1), the first term in (B.6), the
third term in (B.6), the first term in (B.7), and the third term in (B.7) yields

$$
\begin{aligned}
& G_{j+1-2 m}(z)=\sum_{\Omega_{j} \subset J_{N}} \sum_{\substack{s \in J_{N} \\
s_{j+1-m} \prec \prec \prec s_{j+2-m} \\
\bar{s} \notin \Omega_{j}}} d_{\Omega_{j}}(x, r): \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{s}\left(x^{j+1-2 m} z\right): \\
& -\sum_{\Omega_{j} \subset J_{N}} \sum_{\substack{s \in J_{N} \\
s_{j-m} \prec \prec s_{j-m+1} \\
\bar{s} \notin \Omega_{j}}} d_{\Omega_{j}}(x, r): \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{s}\left(x^{j+1-2 m} z\right): \\
& +\sum_{\Omega_{j} \subset J_{N}} \sum_{n=1}^{N} \sum_{\substack{l=j+2-m \\
s_{j}-m+2<s \\
\prec s_{j-m+2} \\
s_{l}=\bar{s}, s=n}}^{j} \Delta\left(x^{2(l+m+n-j-N-2)}\right) d_{\Omega_{j}}(x, r): \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{s}\left(x^{j+1-2 m} z\right): \\
& -\sum_{\Omega_{j} \subset J_{N}} \sum_{n=1}^{N} \sum_{\substack{l=j+1-m \\
s_{j-m} \prec s \\
\prec s_{j-m+1} \\
s_{l}=\bar{s}, s=n}}^{j} \Delta\left(x^{2(l+m+n-j-N-2)}\right) d_{\Omega_{j}}(x, r): \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{s}\left(x^{j+1-2 m} z\right): \\
& +\sum_{\Omega_{j} \subset J_{N}} \sum_{n=1}^{N} \sum_{\substack{k=1 \\
s_{k}=\bar{s}, s=\bar{n} \\
s_{j}+1-m \prec s \\
\prec s_{j}+2-m}} \Delta\left(x^{2(j-m-k+n-N)}\right) d_{\Omega_{j}}(x, r): \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{s}\left(x^{j+1-2 m} z\right): \\
& -\sum_{\Omega_{j} \subset J_{N}} \sum_{n=1}^{N} \sum_{\substack{k=1 \\
s_{k}=\bar{s}, s=\bar{n} \\
s_{j}-m \prec s \\
\prec s_{j-m+1}}}^{j-m} \Delta\left(x^{2(j-m-k+n-N)}\right) d_{\Omega_{j}}(x, r): \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{s}\left(x^{j+1-2 m} z\right): .
\end{aligned}
$$

Rewriting the sum, in the same way as the case of $\bar{G}_{j+1-2 m}(z)$, yields

$$
\begin{aligned}
& G_{j+1-2 m}(z)=\sum_{\substack{\Omega_{j+1} \subset J_{N} \\
\bar{t}_{j+2-m} \notin \Omega_{j+1} \backslash\left\{t_{j+2-m}\right\}}} d_{\left\{t_{1}, \ldots, t_{j+1-m}\right\}}(x, r) d_{\left\{t_{j+3-m}, \ldots, t_{j+1}\right\}}(x, r) \\
& \times \prod_{p=1}^{j+1-m} \prod_{\substack{ \\
q=m+j+3-m \\
t_{q}=\bar{t}_{p}}}^{j+1} \Delta\left(x^{2\left(q-p+t_{p}-N-2\right)}\right): \vec{\Lambda}_{\Omega_{j+1} \backslash\left\{t_{j+2-m}\right\}}(z) \Lambda_{t_{j+2-m}}\left(x^{j+1-2 m} z\right): \\
& -\sum_{\substack{\Omega_{j+1} \subset J_{N} \\
\bar{t}_{j+1-m} \notin \Omega_{j+1} \backslash\left\{t_{j+1-m}\right\}}} d_{\left\{t_{1}, \ldots, t_{j-m}\right\}}(x, r) d_{\left\{t_{j+2-m}, \ldots, t_{j+1}\right\}}(x, r) \\
& \times \prod_{p=1}^{j-m} \prod_{\substack{q=j+2-m \\
t_{q}=\bar{t}_{p}}}^{j+1} \Delta\left(x^{2\left(q-p+t_{p}-N-2\right)}\right): \vec{\Lambda}_{\Omega_{j+1} \backslash\left\{t_{j+1-m}\right\}}(z) \Lambda_{t_{j+1-m}}\left(x^{j+1-2 m} z\right): \\
& +\sum_{\Omega_{j+1} \subset J_{N}} \sum_{\substack{n=1 \\
t_{j-m+2}=n}}^{N} \sum_{\substack{l=j+2-m \\
t_{l}=\bar{n}}}^{j+1} d_{\left\{t_{1}, \ldots, t_{j+1-m}\right\}}(x, r) d_{\left\{t_{j+3-m}, \ldots, t_{j+1}\right\}}(x, r) \\
& \times \prod_{p=1}^{j+1-m} \prod_{\substack{q=j+3-m \\
t_{q}=\bar{t}_{p}}}^{j+1} \Delta\left(x^{2\left(q-p+t_{p}-N-2\right)}\right) \Delta\left(x^{2(l+m+n-j-N-2)}\right)
\end{aligned}
$$

$$
\begin{aligned}
& \times: \vec{\Lambda}_{\Omega_{j+1} \backslash\left\{t_{j+2-m}\right\}}(z) \Lambda_{t_{j+2-m}}\left(x^{j+1-2 m} z\right): \\
& -\sum_{\Omega_{j+1} \subset J_{N}} \sum_{\substack{n=1 \\
t_{j+1}=m}}^{N} \sum_{\substack{l=j+1-m \\
t_{l}=\bar{n}}}^{j} d_{\left\{t_{1}, \ldots, t_{j-m}\right\}}(x, r) d_{\left\{t_{j+2-m}, \ldots, t_{j+1}\right\}}(x, r) \\
& \times \prod_{p=1}^{j-m} \prod_{q=j+2-m}^{j+1} \Delta\left(x^{2\left(q-p+t_{p}-N-2\right)}\right) \Delta\left(x^{2(l+m+n-j-N-2)}\right) \\
& \times: \vec{\Lambda}_{\Omega_{j+1} \backslash\left\{t_{j+1-m}\right\}}(z) \Lambda_{t_{j+1-m}}\left(x^{j+1-2 m} z\right): \\
& +\sum_{\Omega_{j+1} \subset J_{N}} \sum_{\substack{n=1 \\
t_{j+2}=\bar{n}=\bar{n}}}^{N} \sum_{\substack{k=1 \\
t_{k}=n}}^{j+1-m} d_{\left\{t_{1}, \ldots, t_{j+1-m}\right\}}(x, r) d_{\left\{t_{j+3-m}, \ldots, t_{j+1}\right\}}(x, r) \\
& \times \prod_{p=1}^{j+1-m} \prod_{\substack{q=j+3-m \\
t_{q}=\bar{t}_{p}}}^{j+1} \Delta\left(x^{2\left(q-p+t_{p}-N-2\right)}\right) \Delta\left(x^{2(j+n-m-k-N)}\right) \\
& \times: \vec{\Lambda}_{\Omega_{j+1} \backslash\left\{t_{j+2-m}\right\}}(z) \Lambda_{t_{j+2-m}}\left(x^{j+1-2 m} z\right): \\
& \left.-\sum_{\Omega_{j+1} \subset J_{N}} \sum_{\substack{n=1 \\
t_{j+1}-\bar{m}=\bar{n} \\
t_{k}=n}}^{N} \sum_{\substack{k=1 \\
j-m}}^{\left\{t_{1}, \ldots, t_{j-m}\right\}}\right\}(x, r) d_{\left\{t_{j-m+1}, \ldots, t_{j+1}\right\}}(x, r) \\
& \times \prod_{p=1}^{j-m} \prod_{\substack{q=j+2-m \\
t_{q}=\bar{t}_{p}}}^{j+1} \Delta\left(x^{2\left(q-p+t_{p}-N-2\right)}\right) \Delta\left(x^{2(j+n-m-k-N)}\right) \\
& \times: \vec{\Lambda}_{\Omega_{j+1} \backslash\left\{t_{j+1-m}\right\}}(z) \Lambda_{t_{j+1-m}}\left(x^{j+1-2 m} z\right): .
\end{aligned}
$$

Rewriting the sum yields

$$
\begin{aligned}
G_{j+1-2 m}(z)= & \sum_{\Omega_{j+1} \subset J_{N}} d_{\left\{t_{1}, \ldots, t_{j+1-m}\right\}}(x, r) d_{\left\{t_{j+3-m}, \ldots, t_{j+1}\right\}}(x, r) \\
& \times \prod_{p=1}^{j+1-m} \prod_{\substack{q=j+3-m \\
t_{q}=\bar{t}_{p}}}^{j+1} \Delta\left(x^{2\left(q-p+t_{p}-N-2\right)}\right) \prod_{\substack{p=1 \\
t_{j+2-m}=\bar{t}_{p}}}^{j+1-m} \Delta\left(x^{2\left(j+t_{p}-m-p-N\right)}\right) \\
& \times \prod_{\substack{q=j+3-m \\
t_{q}=\bar{t}_{j+2-m}}}^{j+1} \Delta\left(x^{2\left(q+m+t_{j+2-m}-j-N-2\right)}\right): \vec{\Lambda}_{\Omega_{j+1} \backslash\left\{t_{j+2-m}\right\}}(z) \Lambda_{t_{j+2}-m}\left(x^{j+1-2 m} z\right): \\
& -\sum_{\Omega_{j+1} \subset J_{N}} d_{\left\{t_{1}, \ldots, t_{j-m}\right\}}(x, r) d_{\left\{t_{j+2-m}, \ldots, t_{j+1}\right\}}(x, r) \\
& \times \prod_{p=1}^{j-m} \prod_{q=j+2-m}^{j+1} \Delta\left(x^{2\left(q-p+t_{p}-N-2\right)}\right) \prod_{t_{p=1}}^{j-m} \Delta\left(x^{2\left(j+t_{p}-m-N-p\right)}\right) \\
& \times \prod_{t_{j+1}=\bar{t}_{p}=\bar{t}_{p}}^{j+1} \Delta\left(x^{2\left(q+m+t_{j+1-m}-j-N-3\right)}\right): \vec{\Lambda}_{\Omega_{j+1} \backslash\left\{t_{j+1-m}\right\}}(z) \Lambda_{t_{j+1-m}}\left(x^{j+1-2 m} z\right): . \\
&
\end{aligned}
$$

Using

$$
\begin{aligned}
& d_{\left\{t_{1}, \ldots, t_{j+1-m}\right\}}(x, r) d_{\left\{t_{j+3-m}, \ldots, t_{j+1}\right\}}(x, r) \prod_{p=1}^{j+1-m} \prod_{\substack{q=j+3-m \\
t_{q}=\bar{t}_{p}}}^{j+1} \Delta\left(x^{2\left(q-p+t_{p}-N-2\right)}\right) \\
& \times \prod_{\substack{ \\
t_{j+2-m}=\bar{t}_{p}}}^{j+1-m} \Delta\left(x^{2\left(j+t_{p}-m-p-N\right)}\right) \prod_{\substack{q=j+3-m \\
t_{q}=\bar{t}_{j+2-m}}}^{j+1} \Delta\left(x^{2\left(q+m+t_{j+2-m}-j-N-3\right)}\right) \\
& =d_{\left\{t_{1}, \ldots, t_{j-m}\right\}}(x, r) d_{\left\{t_{j-m+2}, \ldots, t_{j+1}\right\}}(x, r) \prod_{p=1}^{j-m} \prod_{q=j+2-m}^{j+1} \Delta\left(x^{2\left(q-p+t_{p}-N-2\right)}\right) \\
& \quad \times \prod_{p=1}^{j-t_{q}=\bar{t}_{p}} \\
& \quad \Delta\left(x^{2\left(j+t_{p}-m-p-N\right)}\right) \prod_{\substack{q=j+2-m}}^{j+1} \Delta\left(x^{2\left(q+m+t_{j+1-m}-j-N-3\right)}\right) \\
& t_{q}=\bar{t}_{j+1-m}
\end{aligned}
$$

and

$$
: \vec{\Lambda}_{\Omega_{j+1} \backslash\left\{t_{j+2-m}\right\}}(z) \Lambda_{t_{j+2-m}}\left(x^{j+1-2 m} z\right):=: \vec{\Lambda}_{\Omega_{j+1} \backslash\left\{t_{j+1-m}\right\}}(z) \Lambda_{t_{j+1-m}}\left(x^{j+1-2 m} z\right):
$$

yields $G_{j+1-2 m}(z)=0,1 \leq m \leq\left[\frac{j}{2}\right], 1 \leq j \leq N$.
We define $\bar{H}_{2 N-j+2-2 m}(z), 1 \leq j \leq N, 1 \leq m \leq N-\left[\frac{j-1}{2}\right]$, as the coefficient of $\delta\left(x^{-2 N+j-2+2 m} z_{2} / z_{1}\right)$ in (4.22). We set

$$
\bar{H}_{2 N-j+2-2 m}(z)=\sum_{\varepsilon= \pm} \varepsilon\left(\bar{\beta}_{\varepsilon}(z)+\bar{\gamma}_{\varepsilon}(z)+\bar{\delta}_{\varepsilon}(z)\right)
$$

where we give $\bar{\beta}_{+}(z), \bar{\beta}_{-}(z), \bar{\gamma}_{+}(z), \bar{\gamma}_{-}(z), \bar{\delta}_{+}(z)$, and $\bar{\delta}_{-}(z)$ in (4.25), (4.26), (4.27), (4.28), (4.29), and (4.30), respectively. Adding the first term in (B.3) and the fourth term in (B.6) yields

$$
\begin{align*}
& +\sum_{\Omega_{j} \subset J_{N}} \sum_{n=m+1}^{\operatorname{Min}(N, j+m)} \sum_{\substack{j=1 \\
s_{k}-1 \prec n \prec s_{k} \\
s_{l}=\bar{n}, l=j+m+1-n}}^{\substack{j+1-n}} d_{\Omega_{j}}(x, r) \Delta\left(x^{2(m+j-N-k))}\right. \\
& \times: \Lambda_{n}\left(x^{-2 N+j+2 m-2} z\right) \vec{\Lambda}_{\Omega_{j}}(z): . \tag{4.25}
\end{align*}
$$

Adding the second term in (B.3) and the second term in (B.6) yields

$$
\begin{align*}
\bar{\beta}_{-}(z)= & \sum_{\Omega_{j} \subset J_{N}} \sum_{n=m}^{\operatorname{Min}(N, j+m-2)} \sum_{\substack{k=1 \\
s=n \\
l=n-n \\
l=n, l=j+m-n}}^{j+m-n-1} d_{\Omega_{j}}(x, r): \Lambda_{n}\left(x^{-2 N+j+2 m-2} z\right) \vec{\Lambda}_{\Omega_{j}}(z): \\
& +\sum_{\Omega_{j} \subset J_{N}} \sum_{n=m}^{\operatorname{Min}(N, j+m-1)} \sum_{\substack{k=1 \\
s_{k-1}<n<s_{k} \\
s+m-n}} d_{\Omega_{j}}(x, r) \Delta\left(x^{2(m+j-N-k)}\right) \\
& \times: \Lambda_{n}\left(x^{-2 N+j+2 m-2} z\right) \vec{\Lambda}_{\Omega_{j}}(z): .
\end{align*}
$$

Adding the first term in (B.5) yields

$$
\begin{equation*}
\bar{\gamma}_{+}(z)=\sum_{\substack{\Omega_{j} \subset J_{N} \\ s_{k}=0, k=j+m-N}} d_{\Omega_{j}}(x, r): \Lambda_{0}\left(x^{-2 N+j-2+2 m} z\right) \vec{\Lambda}_{\Omega_{j}}(z): \tag{4.27}
\end{equation*}
$$

Adding the second term in (B.5) yields

$$
\begin{equation*}
\bar{\gamma}_{-}(z)=\sum_{\substack{\Omega_{j} \subset J_{N} \\ s_{k}=0, k=j+m-N-1}} d_{\Omega_{j}}(x, r): \Lambda_{0}\left(x^{-2 N+j-2+2 m} z\right) \vec{\Lambda}_{\Omega_{j}}(z): . \tag{4.28}
\end{equation*}
$$

Adding the first term in (B.4) and the fourth term in (B.7) yields

$$
\begin{align*}
\bar{\delta}_{+}(z)= & \sum_{\Omega_{j} \subset J_{N}} \sum_{n=2 N+2-j-m}^{N} \sum_{\substack{l=m+j+n-2 N \\
s_{k}=n, k=j+m+n-2 N-1 \\
s_{l}=\bar{n}}}^{j} d_{\Omega_{j}}(x, r): \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{\bar{n}}\left(x^{-2 N+j+2 m-2} z\right): \\
& +\sum_{\Omega_{j} \subset J_{N}} \sum_{n=2 N+2-j-m}^{N} \sum_{\substack{l=m+j+n-2 N-1 \\
s_{k}=n, k=j+m+n-2 N-1 \\
s_{l}<\bar{n}<s_{l}+1}}^{j} d_{\Omega_{j}}(x, r) \Delta\left(x^{2(N+l+1-m-j)}\right) \\
& \times: \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{\bar{n}}\left(x^{-2 N+j+2 m-2} z\right): . \tag{4.29}
\end{align*}
$$

Adding the second term in (B.4) and the second term in (B.7) yields

$$
\begin{align*}
\bar{\delta}_{-}(z)= & \sum_{\Omega_{j} \subset J_{N}} \sum_{n=2 N+3-j-m}^{N} \sum_{\substack{l=j+m+n-2 N-1 \\
s_{k}=n, k=j+m+n-2 N-2 \\
s_{l}=\bar{n}}}^{j} d_{\Omega_{j}}(x, r): \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{\bar{n}}\left(x^{-2 N+j+2 m-2} z\right): \\
& +\sum_{\Omega_{j} \subset J_{N}} \sum_{n=2 N+3-j-m}^{N} \sum_{\substack{l=j+m+n-2 N-2 \\
s_{k}=n, k=j+m+n-2 N-2 \\
s_{l} \prec n<s_{l}+1}}^{j} d_{\Omega_{j}}(x, r) \Delta\left(x^{2(N+l+1-m-j)}\right) \\
& \times: \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{\bar{n}}\left(x^{-2 N+j+2 m-2 z): .}\right. \tag{4.30}
\end{align*}
$$

We show $\bar{H}_{2 N-j+2-2 m}(z)=0,1 \leq j \leq N, 1 \leq m \leq N-\left[\frac{j-1}{2}\right]$. In this proof we frequently use relation (4.4). The proof is divided into three cases: $(i), 1 \leq m \leq N-j$, (ii), $m=N+1-j$, and (iii), $N+2-j \leq m \leq N-\left[\frac{j-1}{2}\right]$.

First, we study the case $(i), j+m \leq N$. In the case $(i), \bar{\gamma}_{ \pm}(z)$ and $\bar{\delta}_{ \pm}(z)$ vanish. Hence, we have $\bar{H}_{2 N-j+2-2 m}(z)=\bar{\beta}_{+}(z)-\bar{\beta}_{-}(z)$. We start from $\bar{\beta}_{+}(z)$. Rewriting the sum yields

$$
\begin{aligned}
& \bar{\beta}_{+}(z)=\left(\sum_{n=m+1}^{\operatorname{Min}(N, j+m-1)} \sum_{k=1}^{j+m-n} \sum_{\substack{\Omega_{j} \subset J_{N} \\
s_{k}=n}} \sum_{\substack{r=0 \\
\frac{s}{s}, s_{l+1}, \ldots, s_{l+r}=(\bar{n}, \overline{n-1}, \ldots, n-n) \\
n-r-1 \prec s_{l+r+1}, l=m+j+1-n}}^{n-m-1} \Delta\left(x^{2(m+j-N-k)}\right)\right. \\
& \left.+\sum_{n=m+1}^{\operatorname{Min}(N, j+m)} \sum_{k=1}^{j+m-n-1} \sum_{\substack{\Omega_{j} \subset J_{N} \\
s_{k-1} \prec n<s_{k}}} \sum_{\substack{r=0 \\
\left(s_{l}, s_{l+1}, \ldots, s_{l+r}\right)=(\bar{n}, n-1, \ldots, \overline{n-r}) \\
n-r-1 \prec s_{l+r+1}, l=m+j+1-n}}^{n-m-1} \Delta\left(x^{2(m+j-N-k)}\right)\right)
\end{aligned}
$$

$$
\begin{align*}
& \times \prod_{a=1}^{k-1} \prod_{\substack{b=1 \\
s_{a}=n-b}}^{r} \Delta\left(x^{2(m+j-N-a)}\right) \prod_{p=1}^{k-1} \prod_{\substack{q=l+r+1 \\
s_{q}=\bar{s}_{p}}}^{j} \Delta\left(x^{2\left(q-p+s_{p}-N-1\right)}\right) \\
& \times \prod_{\substack{ \\
k+1 \leq p<q \leq l-1 \\
s_{q}=\bar{s} p}} \Delta\left(x^{2\left(q-p+s_{p}-N-1\right)}\right): \Lambda_{n}\left(x^{-2 N+j+2 m-2} z\right) \vec{\Lambda}_{\Omega_{j}}(z): . \tag{4.31}
\end{align*}
$$

Using

$$
\begin{align*}
& : \Lambda_{n}\left(x^{-2 N+j+2 m-2} z\right) \vec{\Lambda}_{\left\{s_{1}, s_{2}, \ldots, s_{l-1}, \bar{n}, \overline{n-1}, \ldots, \overline{n-r}, s_{l+r+1}, s_{l+r+2}, \ldots, s_{j}\right\}}(z): \\
& \quad=: \Lambda_{n-r-1}\left(x^{-2 N+j+2 m-2} z\right) \vec{\Lambda}_{\left\{s_{1}, s_{2}, \ldots, s_{l-1}, \overline{n-1}, \overline{n-2}, \ldots, \overline{n-r-1}, s_{l+r+1}, s_{l+r+2}, \ldots, s_{j}\right\}}(z): \\
& 0 \leq r \leq n-m-1, \quad l=m+j+1-n \tag{4.32}
\end{align*}
$$

obtained from (4.4), and replacing $n$ by $n+1$ yields $\bar{\beta}_{+}(z)=\bar{\beta}_{-}(z)$. Hence, we obtain $\bar{H}_{2 N-j+2-2 m}(z)=\bar{\beta}_{+}(z)-\bar{\beta}_{-}(z)=0$ for $1 \leq m \leq N-j$.

Next, we examine the case $(i i), m=N+1-j$. In the case $(i i), \bar{\delta}_{ \pm}(z)$ and $\bar{\gamma}_{-}(z)$ vanish. Hence, we have $\bar{H}_{2 N-j+2-2 m}(z)=\bar{\beta}_{+}(z)+\bar{\gamma}_{+}(z)-\bar{\beta}_{-}(z)$. We start from $\bar{\beta}_{+}(z)+\bar{\gamma}_{+}(z)$. Using (4.4) yields

$$
\begin{equation*}
\bar{\gamma}_{+}(z)=\Delta(1) \sum_{k=1}^{j} \sum_{\substack{\Omega_{j} \subset J_{N} \\ N+1-k} s_{k+1} \prec \cdots \prec s_{j}}: \Lambda_{N+1-k}\left(x^{-j} z\right) \vec{\Lambda}_{\left\{\bar{N}, \overline{N-1}, \ldots, \overline{N+1-k}, s_{k+1}, s_{k+2}, \ldots, s_{j}\right\}}(z): \tag{4.33}
\end{equation*}
$$

Using (4.31), (4.32), and (4.33) yields $\bar{\beta}_{+}(z)+\bar{\gamma}_{+}(z)=\bar{\beta}_{-}(z)$. Hence, we obtain $\bar{H}_{2 N-j+2-2 m}(z)$ $=\bar{\beta}_{+}(z)+\bar{\gamma}_{+}(z)-\bar{\beta}_{-}(z)=0$ for $m=N-j+1$.

Next, we examine the case $(i i i), N+2-j \leq m \leq N-\left[\frac{j-1}{2}\right]$. Rewriting the sum yields

$$
\begin{align*}
& \bar{\delta}_{+}(z)=\left(\sum_{n=2 N+2-j-m}^{N} \sum_{l=j+m+n-2 N} \sum_{\substack{r=0 \\
\Omega_{j} \subset J_{N} \\
s_{l}=\bar{n}}}^{j} \sum_{\substack{\left(s_{k-r}, s_{k-r+1}, \ldots, s_{k}\right) \\
=(n-r, n-r+1, \ldots, n) \\
s_{k-r-1} \prec n-r-1 \\
k=m+j+n-2 N-1}}^{j+m+n-2 N-2} \Delta\left(x^{2(l-j-m+N)}\right)\right. \\
& \left.+\sum_{n=2 N+2-j-m}^{N} \sum_{l=m+j+n-2 N-1}^{j} \sum_{\substack{r=0 \\
\Omega_{j} \subset J_{N} \\
s_{l}=\bar{n}}}^{j} \sum_{\substack{\left(s_{k-r}, s_{k-r+1}, \ldots, s_{k}\right) \\
=(n-r, n-r+1, \ldots, n) \\
s_{k-r} \prec-1<n-r-1 \\
k=m+j+n-2 N-1}}^{j+m+n-2 N-2} \Delta\left(x^{2(l-j-m+N+1)}\right)\right) \\
& \times \prod_{a=1}^{r} \prod_{\substack{b=l+1 \\
\bar{s}_{b}=n-a}}^{j} \Delta\left(x^{2(N+b-j-m)}\right) \prod_{p=1}^{n-r-1} \prod_{\substack{q=l+1 \\
s_{q}=\bar{s}_{p}}}^{j} \Delta\left(x^{2\left(q-p+s_{p}-N-1\right)}\right) \\
& \times \prod_{\substack{ \\
k+1 \leq p<q \leq l-1 \\
s_{q}=\bar{s}_{p}}} \Delta\left(x^{2\left(q-p+s_{p}-N-1\right)}\right): \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{\bar{n}}\left(x^{-2 N+j+2 m-2} z\right): \tag{4.34}
\end{align*}
$$

and

$$
\bar{\gamma}_{+}(z)=\sum_{\substack{\Omega_{j} \subset J_{N} \\ s_{k}=0, k=j+m-N \\ \bar{N} \prec s_{k+1}}} \prod_{\substack{p=1}}^{k-1} \prod_{\substack{q=k+1 \\ s_{q}=\bar{s}_{p}}}^{j} \Delta\left(x^{2\left(q-p+s_{p}-N-1\right)}\right): \vec{\Omega}_{\Lambda_{j}}(z) \Lambda_{0}\left(x^{-2 N+j-2+2 m}\right):
$$

$$
\begin{align*}
& +\sum_{\substack{\Omega_{j} \subset J_{N} \\
s_{k}=0, k=j+m-N}} \sum_{\substack{r=0 \\
N \prec s_{k+1}}} \prod_{\substack{\left(s_{k+1}, \frac{\left(s_{k+2}, \ldots, s_{k+r+1}\right)}{\left.N-\frac{(N-1}{N-1}, \ldots, \frac{N-r}{}\right)}\right.}}^{k-1} \prod_{a=0}^{r-r-1 \prec s_{k+r+2}} \Delta\left(x^{2(j+m-N-a)}\right) \\
& \times \prod_{p=1}^{k-1} \prod_{\substack{q=k+r+2 \\
s_{q}=\bar{s}_{p}}}^{j} \Delta\left(x^{2\left(q-p+s_{p}-N-1\right)}\right): \vec{\Omega}_{\Lambda_{j}}(z) \Lambda_{0}\left(x^{-2 N+j-2+2 m}\right): \tag{4.35}
\end{align*}
$$

Using (4.2) and (4.4) yields

$$
\begin{align*}
& : \vec{\Lambda}_{\left\{s_{1}, s_{2}, \ldots, s_{k-r-1}, n-r, n-r+1, \ldots, n, s_{k+1}, s_{k+2}, \ldots, s_{j}\right\}}(z) \Lambda_{\bar{n}}\left(x^{-2 N+j+2 m-2} z\right): \\
& \quad=: \vec{\Lambda}_{\left\{s_{1}, s_{2}, \ldots, s_{k-r-1}, n-r-1, n-r, \ldots, n-1, s_{k+1}, s_{k+2}, \ldots, s_{j}\right\}}(z) \Lambda_{\overline{n-r-1}}\left(x^{-2 N+j+2 m-2} z\right) \\
& 0 \leq r \leq k-1, \quad k=j+m+n-2 N-1, \tag{4.36}
\end{align*}
$$

and

$$
\begin{align*}
& : \vec{\Lambda}_{\left\{s_{1}, s_{2}, \ldots, s_{k-1}, 0, \bar{N}, \overline{N-1}, \ldots, \overline{N-r}, s_{k+r+2}, s_{k+r+3}, \ldots, s_{j}\right\}}(z) \Lambda_{0}\left(x^{-2 N+j+2 m-2} z\right): \\
& \quad=\Delta(1): \vec{\Lambda}_{\left\{s_{1}, s_{2}, \ldots, s_{k-1}, \bar{N}, \overline{N-1}, \ldots, \overline{N-r-1}, s_{k+r+2}, s_{k+r+3}, \ldots, s_{j}\right\}}(z) \Lambda_{N-r-1}\left(x^{-2 N+j+2 m-2} z\right): \\
& 0 \leq r \leq 2 N+l+1-j-m, \quad k=j+m+n-2 N-1 . \tag{4.37}
\end{align*}
$$

Using (4.31), (4.32), (4.34), (4.35), (4.36), and (4.37) and replacing $n$ by $n+1$ yield $\bar{\beta}_{+}(z)+$ $\bar{\gamma}_{+}(z)+\bar{\delta}_{+}(z)=\bar{\beta}_{-}(z)+\bar{\gamma}_{-}(z)+\bar{\delta}_{-}(z)$. Hence, we obtain $\bar{H}_{2 N-j+2-2 m}(z)=\sum_{\varepsilon= \pm} \varepsilon\left(\bar{\beta}_{\varepsilon}(z)+\right.$ $\left.\bar{\gamma}_{\varepsilon}(z)+\bar{\delta}_{\varepsilon}(z)\right)=0$ for $N+2-j \leq m \leq N-\left[\frac{j-1}{2}\right]$. Finally, we obtain $\bar{H}_{2 N-j+2-2 m}(z)=0$ for $1 \leq m \leq N-\left[\frac{j-1}{2}\right]$.

We define $H_{2 N-j+2-2 m}(z), 1 \leq j \leq N, 1 \leq m \leq N-\left[\frac{j-1}{2}\right]$, as the coefficient of $\delta\left(x^{-2 N+j-2+2 m} z_{2} / z_{1}\right)$ in (4.22). We set

$$
H_{2 N-j+2-2 m}(z)= \begin{cases}\sum_{\varepsilon= \pm} \varepsilon\left(\beta_{\varepsilon}(z)+\gamma_{\varepsilon}(z)+\delta_{\varepsilon}(z)\right) & \text { otherwise, }  \tag{4.38}\\ 0 & \text { if } j \text { is even, } \quad m=N-\left[\frac{j-1}{2}\right]\end{cases}
$$

where we give $\beta_{+}(z), \beta_{-}(z), \gamma_{+}(z), \gamma_{-}(z), \delta_{+}(z)$, and $\delta_{-}(z)$ in (4.39), (4.40), (4.41), (4.42), (4.43), and (4.44), respectively. In (4.38) we define $H_{0}(z)=0$ to avoid ambiguity of $\bar{H}_{0}(z)$ and $H_{0}(z)$. In the case when $j$ is even, we have $\mathrm{LHS}_{1, j}=c(x, r)\left(\bar{H}_{0}\left(z_{2}\right)-H_{0}\left(z_{2}\right)\right) \delta\left(z_{2} / z_{1}\right)+$ $\left.\bar{H}_{2}\left(z_{2}\right) \delta\left(x^{-2} z_{2} / z_{1}\right)-H_{2}\left(z_{2}\right) \delta\left(x^{2} z_{2} / z_{1}\right)+\cdots\right)$. Adding the first term in (B.4) and the fourth term in (B.7) yields

$$
\begin{align*}
\beta_{+}(z)= & \sum_{\Omega_{j} \subset J_{N}} \sum_{n=m}^{\operatorname{Min}(N, j+m-2)} \sum_{\substack{l=n+2-m \\
s_{k}=n, k=n+1-m \\
=\bar{n}+1-m}}^{j} d_{\Omega_{j}}(x, r): \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{\bar{n}}\left(x^{2 N+2-j-2 m} z\right): \\
& +\sum_{\Omega_{j} \subset J_{N}} \sum_{n=m}^{\operatorname{Min}(N, j+m-1)} \sum_{\substack{l=n+1-m \\
s_{k}=n, k=n+1-m \\
\text { sl }<\bar{n} s_{l+1}}}^{j} d_{\Omega_{j}}(x, r) \Delta\left(x^{2(l+m-N-1)}\right) \\
& \times: \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{\bar{n}}\left(x^{2 N+2-j-2 m} z\right): . \tag{4.39}
\end{align*}
$$

Adding the second term in (B.4) and the second term in (B.7) yields

$$
\begin{align*}
\beta_{-}(z)= & \sum_{\Omega_{j} \subset J_{N}} \sum_{n=m+1}^{\operatorname{Min}(N, j+m-1)} \sum_{\substack{l=n+1-m \\
s_{k}=n, k=n-m \\
s_{l}=\bar{n}}}^{j} d_{\Omega_{j}}(x, r): \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{\bar{n}}\left(x^{2 N+2-j-2 m} z\right): \\
& +\sum_{\Omega_{j} \subset J_{N}} \sum_{n=m+1}^{\operatorname{Min}(N, j+m)} \sum_{\substack{l=n-m \\
s_{k}=n, k=n-m \\
s_{l} \prec n<s_{l}+1}}^{j} d_{\Omega_{j}}(x, r) \Delta\left(x^{2(l+m-N-1)}\right) \\
& \times: \vec{\Lambda}_{\Omega_{j}}(z) \Lambda_{\bar{n}}\left(x^{2 N+2-j-2 m} z\right): . \tag{4.40}
\end{align*}
$$

Adding the first term in (B.5) yields

$$
\begin{equation*}
\gamma_{+}(z)=\sum_{\substack{\Omega_{j} \subset J_{N} \\ s_{k}=0, k=N+2-m}} d_{\Omega_{j}}(x, r): \Lambda_{0}\left(x^{2 N+2-j-2 m} z\right) \vec{\Lambda}_{\Omega_{j}}(z): . \tag{4.41}
\end{equation*}
$$

Adding the second term in (B.5) yields

$$
\begin{equation*}
\gamma_{-}(z)=\sum_{\substack{\Omega_{j} \subset J_{N} \\ s_{k}=0, k=N+1-m}} d_{\Omega_{j}}(x, r): \Lambda_{0}\left(x^{2 N+2-j-2 m} z\right) \vec{\Lambda}_{\Omega_{j}}(z): \tag{4.42}
\end{equation*}
$$

Adding the first term in (B.3) and the fourth term in (B.6) yields

$$
\begin{align*}
\delta_{+}(z)= & \sum_{\Omega_{j} \subset J_{N}} \sum_{n=2 N+3-j-m}^{N} \sum_{\substack{k=1 \\
s_{k}=n \\
s_{l}=\bar{n}, l=2 N+3-m-n}}^{2 N+2-m-n} d_{\Omega_{j}}(x, r): \Lambda_{n}\left(x^{2 N+2-j-2 m} z\right) \vec{\Lambda}_{\Omega_{j}}(z): \\
& +\sum_{\Omega_{j} \subset J_{N}} \sum_{n=2 N+3-j-m}^{N} \sum_{\substack{k=1 \\
s_{k-1}<n<s_{k} \\
s_{l}=\bar{n}, l=2 N+3-m-n}}^{2 N+3-m-n} d_{\Omega_{j}}(x, r) \Delta\left(x^{2(N+2-m-k)}\right) \\
& \times: \Lambda_{n}\left(x^{2 N+2-j-2 m} z\right) \vec{\Lambda}_{\Omega_{j}}(z): . \tag{4.43}
\end{align*}
$$

Adding the second term in (B.3) and the second term in (B.6) yields

$$
\begin{align*}
\delta_{-}(z)= & \sum_{\Omega_{j} \subset J_{N}} \sum_{n=2 N+2-j-m}^{N} \sum_{\substack{k=1 \\
s_{k}=n \\
s_{l}=n, 2-m-n \\
s_{l}=2 N+2-m-n}}^{2 N+1-m-n} d_{\Omega_{j}}(x, r): \Lambda_{n}\left(x^{2 N+2-j-2 m} z\right) \vec{\Lambda}_{\Omega_{j}}(z): \\
& +\sum_{\Omega_{j} \subset J_{N}} \sum_{n=2 N+2-j-m}^{N} \sum_{\substack{k=1 \\
s_{k-1}<n<s_{k} \\
s_{l}=\bar{n}, l=2 N+2-m-n}}^{2 N+2-m-n} d_{\Omega_{j}}(x, r) \Delta\left(x^{2(N+2-m-k)}\right) \\
& \times: \Lambda_{n}\left(x^{2 N+2-j-2 m} z\right) \vec{\Lambda}_{\Omega_{j}}(z): . \tag{4.44}
\end{align*}
$$

The relation $H_{2 N+2-j-2 m}(z)=0$ is shown in the same way as $\bar{H}_{2 N+2-j-2 m}(z)=0$.
Lemma 4.7. The currents $T_{i}(z)$ satisfy the following fusion relation

$$
\begin{align*}
& \lim _{z_{1} \rightarrow x^{ \pm(i+j)} z_{2}}\left(1-\frac{x^{ \pm(i+j)} z_{2}}{z_{1}}\right) f_{i, j}\left(\frac{z_{2}}{z_{1}}\right) T_{i}\left(z_{1}\right) T_{j}\left(z_{2}\right) \\
& \quad=\mp c(x, r) \prod_{l=1}^{\operatorname{Min}(i, j)-1} \Delta\left(x^{2 l+1}\right) T_{i+j}\left(x^{ \pm i} z_{2}\right), \quad 1 \leq i, j \leq N .
\end{align*}
$$

Proof. For $\Omega_{i}^{(1)}=\left\{s_{1}, s_{2}, \ldots, s_{i}\right\} \subset J_{N}$ with $s_{1} \prec s_{2} \prec \cdots \prec s_{i}$ and $\Omega_{j}^{(2)}=\left\{t_{1}, t_{2}, \ldots, t_{j}\right\}$ $\subset J_{N}$ with $t_{1} \prec t_{2} \prec \cdots \prec t_{j}$, we set $\Omega_{i+j}=\Omega_{i}^{(1)} \cup \Omega_{j}^{(2)}$. From (4.1), the necessary and sufficient condition that $f_{i, j}\left(z_{2} / z_{1}\right) \vec{\Lambda}_{\Omega_{i}^{(1)}}\left(z_{1}\right) \vec{\Lambda}_{\Omega_{j}^{(2)}}\left(z_{2}\right)$ has a pole at $z_{1}=x^{-(i+j)} z_{2}$ (respectively $z_{1}=x^{i+j} z_{2}$ ) is $s_{i} \prec t_{1}$ (respectively $t_{j} \prec s_{1}$ ). In the case when $s_{i} \prec t_{1}$ or $t_{j} \prec s_{1}$, we obtain

$$
\begin{aligned}
& f_{i, j}\left(\frac{z_{2}}{z_{1}}\right) \vec{\Lambda}_{\Omega_{i}^{(1)}\left(z_{1}\right)} \vec{\Lambda}_{\Omega_{j}^{(2)}}\left(z_{2}\right)=\prod_{k=0}^{\operatorname{Min}(i, j)-1} \Delta\left(\frac{x^{ \pm(2 k+1-i-j)} z_{2}}{z_{1}}\right) \\
& \quad \times \prod_{\substack{p=1 \\
t_{q}=1 \\
t_{q}=\bar{s}_{p}}}^{j} \Delta\left(\frac{x^{ \pm\left\{2\left(q-p+s_{p}-N-1+i\right)-i-j\right\}} z_{2}}{z_{1}}\right): \vec{\Lambda}_{\Omega_{i}^{(1)}}\left(z_{1}\right) \vec{\Lambda}_{\Omega_{j}^{(2)}}\left(z_{2}\right): .
\end{aligned}
$$

The signs $\pm$ in the products in the above expression of $f_{i, j}\left(z_{2} / z_{1}\right) \vec{\Lambda}_{\Omega_{i}^{(1)}}\left(z_{1}\right) \vec{\Lambda}_{\Omega_{j}^{(2)}}\left(z_{2}\right)$ are in the same order. The upper sign is for $s_{i} \prec t_{1}$, and the lower sign is for $t_{j} \prec s_{1}$. Taking the limit yields

$$
\begin{align*}
& \quad \lim _{z_{1} \rightarrow x^{ \pm(i+j)} z_{2}}\left(1-\frac{x^{ \pm(i+j)} z_{2}}{z_{1}}\right) f_{i, j}\left(\frac{z_{2}}{z_{1}}\right) \vec{\Lambda}_{\Omega_{i}^{(1)}\left(z_{1}\right)} \vec{\Lambda}_{\Omega_{j}^{(2)}}\left(z_{2}\right)=\mp c(x, r) \prod_{l=1}^{\operatorname{Min}(i, j)-1} \Delta\left(x^{2 l+1}\right) \\
& \quad \times \prod_{p=1}^{i} \prod_{\substack{q=1 \\
t_{q}=\bar{s}_{p}}}^{j} \Delta\left(x^{2\left(q-p+s_{p}-N-1+i\right)}\right) \vec{\Lambda}_{\Omega_{i+j}}\left(x^{ \pm i} z_{2}\right), \quad 1 \leq i, j \leq N . \tag{4.46}
\end{align*}
$$

Here, we use $: \vec{\Lambda}_{\Omega_{i}^{(1)}}\left(x^{ \pm(i+j)} z\right) \vec{\Lambda}_{\Omega_{j}^{(2)}}(z):=\vec{\Lambda}_{\Omega_{i+j}}\left(x^{ \pm i} z\right)$. Adding (4.46) over all $\Omega_{i}^{(1)}$ and $\Omega_{j}^{(2)}$ yields (4.45).

Lemma 4.8. The currents $T_{i}(z)$ satisfy the following fusion relations:

$$
\begin{align*}
& \lim _{z_{1} \rightarrow x^{ \pm(2 N+1+i-j)} z_{2}}\left(1-\frac{x^{ \pm(2 N+1+i-j)} z_{2}}{z_{1}}\right) f_{i, j}\left(\frac{z_{2}}{z_{1}}\right) T_{i}\left(z_{1}\right) T_{j}\left(z_{2}\right) \\
& \quad=\mp c(x, r) \prod_{l=1}^{i-1} \Delta\left(x^{2 l+1}\right) \prod_{l=N+1-j}^{N+i-j} \Delta\left(x^{2 l}\right) T_{j-i}\left(x^{ \pm i} z_{2}\right), \quad 1 \leq i \leq j \leq N,  \tag{4.47}\\
& \lim _{z_{1} \rightarrow x^{ \pm(2 N+1-i+j)} z_{2}}\left(1-\frac{x^{ \pm(2 N+1-i+j)} z_{2}}{z_{1}}\right) f_{i, j}\left(\frac{z_{2}}{z_{1}}\right) T_{i}\left(z_{1}\right) T_{j}\left(z_{2}\right) \\
& =\mp c(x, r) \prod_{l=1}^{j-1} \Delta\left(x^{2 l+1}\right) \prod_{l=N+1-i}^{N+j-i} \Delta\left(x^{2 l}\right) T_{i-j}\left(x^{ \pm(2 N+1-i)} z_{2}\right), \quad 1 \leq j \leq i \leq N . \tag{4.48}
\end{align*}
$$

Proof. Using (3.3), (4.8), (4.9), and (4.45) yields (4.47) and (4.48).
Proof. Here we will give a proof of Theorem 3.2. We prove Theorem 3.2 by induction. Lemma 4.6 is the base for induction. We define $\operatorname{LHS}_{i, j}, \operatorname{RHS1}_{i, j}$ and $\operatorname{RHS2}_{i, j}(k)$ with $1 \leq k \leq i \leq$ $j \leq N$ as

$$
\begin{aligned}
\operatorname{LHS}_{i, j}= & f_{i, j}\left(\frac{z_{2}}{z_{1}}\right) T_{i}\left(z_{1}\right) T_{j}\left(z_{2}\right)-f_{j, i}\left(\frac{z_{1}}{z_{2}}\right) T_{j}\left(z_{2}\right) T_{i}\left(z_{1}\right), \\
\operatorname{RHS}_{i, j}= & c(x, r) \prod_{l=1}^{i-1} \Delta\left(x^{2 l+1}\right) \prod_{l=N+1-j}^{N+i-j} \Delta\left(x^{2 l}\right) \\
& \times\left(\delta\left(\frac{x^{-2 N+j-i-1} z_{2}}{z_{1}}\right) T_{j-i}\left(x^{-i} z_{2}\right)-\delta\left(\frac{x^{2 N-j+i+1} z_{2}}{z_{1}}\right) T_{j-i}\left(x^{i} z_{2}\right)\right),
\end{aligned}
$$

$$
\begin{aligned}
\operatorname{RHS}_{i, j}(k)= & c(x, r) \prod_{l=1}^{k-1} \Delta\left(x^{2 l+1}\right)\left(\delta\left(\frac{x^{-j+i-2 k} z_{2}}{z_{1}}\right) f_{i-k, j+k}\left(x^{j-i}\right) T_{i-k}\left(x^{k} z_{1}\right) T_{j+k}\left(x^{-k} z_{2}\right)\right. \\
& \left.-\delta\left(\frac{x^{j-i+2 k} z_{2}}{z_{1}}\right) f_{i-k, j+k}\left(x^{-j+i}\right) T_{i-k}\left(x^{-k} z_{1}\right) T_{j+k}\left(x^{k} z_{2}\right)\right), \quad 1 \leq k \leq i-1 \\
\operatorname{RHS}_{i, j}(i)= & c(x, r) \prod_{l=1}^{i-1} \Delta\left(x^{2 l+1}\right)\left(\delta\left(\frac{x^{-j-i} z_{2}}{z_{1}}\right) T_{j+i}\left(x^{-i} z_{2}\right)-\delta\left(\frac{x^{j+i} z_{2}}{z_{1}}\right) T_{j+i}\left(x^{i} z_{2}\right)\right)
\end{aligned}
$$

We prove the following relation by induction on $i, 1 \leq i \leq j \leq N$.

$$
\begin{equation*}
\mathrm{LHS}_{i, j}=\mathrm{RHS}_{i, j}+\sum_{k=1}^{i} \operatorname{RHS}_{i, j}(k) \tag{4.49}
\end{equation*}
$$

The base, $i=1 \leq j \leq N$ was proved previously in Lemma 4.6.
We assume that the relation (4.49) holds for some $i, 1 \leq i<j \leq N$, and show that $\mathrm{LHS}_{i+1, j}=$ $\operatorname{RHS1}_{i+1, j}+\sum_{k=1}^{i+1} \operatorname{RHS}_{i+1, j}(k)$ from this assumption. First, we summarize some relations. The assumption (4.49) yields

$$
\begin{align*}
& \lim _{w_{1} \rightarrow x^{2 N-i-j} w_{2}}\left(1-x^{2 N-i-j} \frac{w_{2}}{w_{1}}\right) f_{j-1, i}\left(\frac{w_{2}}{w_{1}}\right) T_{j-1}\left(w_{1}\right) T_{i}\left(w_{2}\right)=0  \tag{4.50}\\
& \lim _{w_{1} \rightarrow x^{2 N-i-j} w_{2}}\left(1-x^{2 N-i-j} \frac{w_{2}}{w_{1}}\right) f_{1, j-i}\left(\frac{w_{2}}{w_{1}}\right) T_{1}\left(w_{1}\right) T_{j-i}\left(w_{2}\right)=0  \tag{4.51}\\
& f_{i, j}\left(\frac{w_{2}}{w_{1}}\right) T_{i}\left(w_{1}\right) T_{j}\left(w_{2}\right)=f_{j, i}\left(\frac{w_{1}}{w_{2}}\right) T_{j}\left(w_{2}\right) T_{i}\left(w_{1}\right) \tag{4.52}
\end{align*}
$$

for $\frac{w_{2}}{w_{1}} \neq x^{ \pm(2 N-j+i+1)}, x^{ \pm(j-i+2 k)}, 1 \leq k \leq i$. Direct calculation yields

$$
\begin{equation*}
\lim _{w_{2} \rightarrow x^{-i-1} w_{1}}\left(1-x^{-i-1} \frac{w_{1}}{w_{2}}\right) \Delta\left(x^{-i} \frac{w_{1}}{w_{2}}\right)=c(x, r) \tag{4.53}
\end{equation*}
$$

Multiplying $\operatorname{LHS}_{i, j}$ by $f_{1, i}\left(z_{1} / z_{3}\right) f_{1, j}\left(z_{2} / z_{3}\right) T_{1}\left(z_{3}\right)$ on the left and using the quadratic relation (4.49) with $i=1$, along with the fusion relation (4.10) yields

$$
\begin{align*}
& f_{1, i}\left(\frac{z_{1}}{z_{3}}\right) f_{1, j}\left(\frac{z_{2}}{z_{3}}\right) T_{1}\left(z_{3}\right) \times \mathrm{LHS}_{i, j} \\
&= f_{1, j}\left(\frac{z_{2}}{z_{3}}\right) f_{i, j}\left(\frac{z_{2}}{z_{1}}\right) f_{1, i}\left(\frac{z_{1}}{z_{3}}\right) T_{1}\left(z_{3}\right) T_{i}\left(z_{1}\right) T_{j}\left(z_{2}\right) \\
&- f_{j, 1}\left(\frac{z_{3}}{z_{2}}\right) f_{j, i}\left(\frac{z_{1}}{z_{2}}\right) T_{j}\left(z_{2}\right) f_{1, i}\left(\frac{z_{1}}{z_{3}}\right) T_{1}\left(z_{3}\right) T_{i}\left(z_{1}\right) \\
&- c(x, r) \Delta\left(x^{2(N+1-j)}\right) \delta\left(\frac{x^{-2 N+j-2} z_{2}}{z_{3}}\right) \Delta\left(\frac{x^{-i} z_{1}}{z_{3}}\right) f_{j-1, i}\left(\frac{x^{-2 N+j-1} z_{1}}{z_{3}}\right) \\
& \times T_{j-1}\left(x^{2 N-j+1} z_{3}\right) T_{i}\left(z_{1}\right)+c(x, r) \Delta\left(x^{2(N+1-j)}\right) \\
& \times \delta\left(\frac{x^{2 N-j+2} z_{2}}{z_{3}}\right) \Delta\left(\frac{x^{i} z_{1}}{z_{3}}\right) f_{j-1, i}\left(\frac{x^{2 N-j+1} z_{1}}{z_{3}}\right) T_{j-1}\left(x^{-2 N+j-1} z_{3}\right) T_{i}\left(z_{1}\right) \\
&- c(x, r) \delta\left(\frac{x^{-j-1} z_{2}}{z_{3}}\right) \Delta\left(\frac{x^{-i} z_{1}}{z_{3}}\right) f_{j+1, i}\left(\frac{x^{-j} z_{1}}{z_{3}}\right) T_{j+1}\left(x^{j} z_{3}\right) T_{i}\left(z_{1}\right) \\
&+ c(x, r) \delta\left(\frac{x^{j+1} z_{2}}{z_{3}}\right) \Delta\left(\frac{x^{i} z_{1}}{z_{3}}\right) f_{j+1, i}\left(\frac{x^{j} z_{1}}{z_{3}}\right) T_{j+1}\left(x^{-j} z_{3}\right) T_{i}\left(z_{1}\right) . \tag{4.54}
\end{align*}
$$

Taking the limit $z_{3} \rightarrow x^{-i-1} z_{1}$ of (4.54) multiplied by $c(x, r)^{-1}\left(1-x^{-i-1} z_{1} / z_{3}\right)$ and using the relations (4.45) and (4.48), (4.50), (4.52), and (4.53) yields

$$
\begin{align*}
& \lim _{z_{3} \rightarrow x^{-i-1} z_{1}} \frac{1}{c(x, r)}\left(1-x^{-i-1} \frac{z_{1}}{z_{3}}\right) f_{1, i}\left(\frac{z_{1}}{z_{3}}\right) f_{1, j}\left(\frac{z_{2}}{z_{3}}\right) T_{1}\left(z_{3}\right) \times \operatorname{LHS}_{i, j} \\
& \quad=f_{i+1, j}\left(\frac{x z_{2}}{z_{1}}\right) T_{i+1}\left(x^{-1} z_{1}\right) T_{j}\left(z_{2}\right)-f_{j, i+1}\left(\frac{x^{-1} z_{1}}{z_{2}}\right) T_{j}\left(z_{2}\right) T_{i+1}\left(x^{-1} z_{1}\right) \\
& \quad+c(x, r) \prod_{l=1}^{i} \Delta\left(x^{2 l+1}\right) \prod_{l=N-j+1}^{N+i+1-j} \Delta\left(x^{2 l}\right) \delta\left(\frac{x^{2 N-j+i+3} z_{2}}{z_{1}}\right) T_{j-i-1}\left(x^{i+1} z_{2}\right) \\
& \quad-c(x, r) \delta\left(\frac{x^{i-j} z_{2}}{z_{1}}\right) f_{i, j+1}\left(x^{-i+j+1}\right) T_{i}\left(z_{1}\right) T_{j+1}\left(x^{-1} z_{2}\right) \\
& \quad+c(x, r) \delta\left(\frac{x^{i+j+2} z_{2}}{z_{1}}\right) \prod_{l=1}^{i} \Delta\left(x^{2 l+1}\right) T_{i+j+1}\left(x^{i+1} z_{2}\right) . \tag{4.55}
\end{align*}
$$

Multiplying $\operatorname{RHS1}_{i, j}$ by $f_{1, i}\left(z_{1} / z_{3}\right) f_{1, j}\left(z_{2} / z_{3}\right) T_{1}\left(z_{3}\right)$ from the left and using fusion relations (4.9) and (4.10) yields

$$
\begin{align*}
& f_{1, i}\left(\frac{z_{1}}{z_{3}}\right) f_{1, j}\left(\frac{z_{2}}{z_{3}}\right) T_{1}\left(z_{3}\right) \times \operatorname{RHS}_{i, j}=c(x, r) \prod_{l=1}^{i-1} \Delta\left(x^{2 l+1}\right) \prod_{l=N+1-j}^{N+i-j} \Delta\left(x^{2 l}\right) \\
& \quad \times\left\{\delta\left(\frac{x^{-2 N+j-i-1} z_{2}}{z_{1}}\right) \Delta\left(\frac{x^{i} z_{1}}{z_{3}}\right) f_{1, j-i}\left(\frac{x^{-i} z_{2}}{z_{3}}\right) T_{1}\left(z_{3}\right) T_{j-i}\left(x^{-i} z_{2}\right)\right. \\
& \left.\quad-\delta\left(\frac{x^{2 N-j+i+1} z_{2}}{z_{1}}\right) \Delta\left(\frac{x^{-i} z_{1}}{z_{3}}\right) f_{1, j-i}\left(\frac{x^{i} z_{2}}{z_{3}}\right) T_{1}\left(z_{3}\right) T_{j-i}\left(x^{i} z_{2}\right)\right\} . \tag{4.56}
\end{align*}
$$

Taking the limit $z_{3} \rightarrow x^{-i-1} z_{1}$ of (4.56) multiplied by $c(x, r)^{-1}\left(1-x^{-i-1} z_{1} / z_{3}\right)$ and using the relations (4.47) and (4.51) yields

$$
\begin{align*}
& \lim _{z_{3} \rightarrow x^{-i-1} z_{1}} \frac{1}{c(x, r)}\left(1-x^{-i-1} \frac{z_{1}}{z_{3}}\right) f_{1, i}\left(\frac{z_{1}}{z_{3}}\right) f_{1, j}\left(\frac{z_{2}}{z_{3}}\right) T_{1}\left(z_{3}\right) \times \mathrm{RHS}_{i, j} \\
& \quad=c(x, r) \prod_{l=1}^{i} \Delta\left(x^{2 l+1}\right) \prod_{l=N+1-j}^{N+i+1-j} \Delta\left(x^{2 l}\right) \delta\left(\frac{x^{-2 N+j-i-1} z_{2}}{z_{1}}\right) T_{j-i-1}\left(x^{-i-1} z_{2}\right) . \tag{4.57}
\end{align*}
$$

Multiplying $\operatorname{RHS}_{i, j}(i)$ by $f_{1, i}\left(z_{1} / z_{3}\right) f_{1, j}\left(z_{2} / z_{3}\right) T_{1}\left(z_{3}\right)$ from the left and using the fusion relation (4.11) yields

$$
\begin{align*}
& f_{1, i}\left(\frac{z_{1}}{z_{3}}\right) f_{1, j}\left(\frac{z_{2}}{z_{3}}\right) T_{1}\left(z_{3}\right) \times \operatorname{RHS}_{i, j}(i) \\
& \quad=c(x, r) \prod_{l=1}^{i-1} \Delta_{1}\left(x^{2 l+1}\right)\left(\delta\left(\frac{x^{-i-j} z_{2}}{z_{1}}\right) f_{1, i+1}\left(\frac{x^{j} z_{1}}{z_{3}}\right) \Delta\left(\frac{x^{i} z_{1}}{z_{3}}\right) T_{1}\left(z_{3}\right) T_{i+j}\left(x^{j} z_{1}\right)\right. \\
& \left.\quad-\delta\left(\frac{x^{i+j} z_{2}}{z_{1}}\right) f_{1, i+1}\left(\frac{x^{-j} z_{1}}{z_{3}}\right) \Delta\left(\frac{x^{-i} z_{1}}{z_{3}}\right) T_{1}\left(z_{3}\right) T_{i+j}\left(x^{-j} z_{1}\right)\right) . \tag{4.58}
\end{align*}
$$

Taking the limit $z_{3} \rightarrow x^{-i-1} z_{1}$ of (4.58) multiplied by $c(x, r)^{-1}\left(1-x^{-i-1} z_{1} / z_{3}\right)$ and using the relations (4.45) and (4.53) yields

$$
\lim _{z_{3} \rightarrow x^{-i-1} z_{1}} \frac{1}{c(x, r)}\left(1-x^{-i-1} \frac{z_{1}}{z_{3}}\right) f_{1, i}\left(\frac{z_{1}}{z_{3}}\right) f_{1, j}\left(\frac{z_{2}}{z_{3}}\right) T_{1}\left(z_{3}\right) \times \operatorname{RHS}_{i, j}(i)
$$

$$
\begin{align*}
= & c(x, r) \delta\left(\frac{x^{-i-j} z_{2}}{z_{1}}\right) \prod_{l=1}^{i} \Delta\left(x^{2 l+1}\right) T_{i+j+1}\left(x^{-i-1} z_{2}\right) \\
& -c(x, r) \delta\left(\frac{x^{i+j} z_{2}}{z_{1}}\right) \prod_{l=1}^{i-1} \Delta\left(x^{2 l+1}\right) f_{1, i+j}\left(x^{i-j+1}\right) T_{1}\left(x^{-i-1} z_{1}\right) T_{i+j}\left(x^{i} z_{2}\right) . \tag{4.59}
\end{align*}
$$

Multiplying $\operatorname{RHS}_{i, j}(k), 1 \leq k \leq i-1$, by $f_{1, i}\left(z_{1} / z_{3}\right) f_{1, j}\left(z_{2} / z_{3}\right) T_{1}\left(z_{3}\right)$ from the left and using relations (4.12) and (4.52) yields

$$
\begin{align*}
& f_{1, i}\left(\frac{z_{1}}{z_{3}}\right) f_{1, j}\left(\frac{z_{2}}{z_{3}}\right) T_{1}\left(z_{3}\right) \times \operatorname{RHS}_{i, j}(k)  \tag{4.60}\\
& \quad=c(x, r) \prod_{l=1}^{k-1} \Delta\left(x^{2 l+1}\right)\left(\delta\left(\frac{x^{-j+i-2 k} z_{2}}{z_{1}}\right) f_{1, i-k}\left(\frac{x^{k} z_{1}}{z_{3}}\right) f_{j+k, i-k}\left(x^{i-j}\right) f_{1, j+k}\right. \\
& \quad \times\left(\frac{x^{-i+j+k} z_{1}}{z_{3}}\right) T_{1}\left(z_{3}\right) T_{j+k}\left(x^{j-i+k} z_{1}\right) T_{i-k}\left(x^{k} z_{1}\right)-\delta\left(\frac{x^{j-i+2 k} z_{2}}{z_{1}}\right) f_{1, i-k}\left(\frac{x^{-k} z_{1}}{z_{3}}\right) \\
& \left.\quad \times f_{i-k, j+k}\left(x^{i-j}\right) f_{1, j+k}\left(\frac{x^{i-j-k} z_{1}}{z_{3}}\right) T_{1}\left(z_{3}\right) T_{i-k}\left(x^{-k} z_{1}\right) T_{j+k}\left(x^{k} z_{2}\right)\right), \quad 1 \leq k \leq i-1 .
\end{align*}
$$

Taking the limit $z_{3} \rightarrow x^{-i-1} z_{1}$ of (4.60) multiplied by $c(x, r)^{-1}\left(1-x^{-i-1} z_{1} / z_{3}\right)$, and using the fusion relations (4.7), (4.45), and (4.52) yields

$$
\begin{align*}
& \lim _{z_{3} \rightarrow x^{-i-1} z_{1}} \frac{1}{c(x, r)}\left(1-x^{-i-1} \frac{z_{1}}{z_{3}}\right) f_{1, i}\left(\frac{z_{1}}{z_{3}}\right) f_{1, j}\left(\frac{z_{2}}{z_{3}}\right) T_{1}\left(z_{3}\right) \times \operatorname{RHS}_{i, j}(k) \\
& =c(x, r) \prod_{l=1}^{k} \Delta\left(x^{2 l+1}\right) \delta\left(\frac{x^{-j+i-2 k} z_{2}}{z_{1}}\right) f_{j+k-1, i-k}\left(x^{i-j+1}\right) T_{i-k}\left(x^{k} z_{1}\right) T_{j+k+1}\left(x^{-k-1} z_{2}\right) \\
& \quad-c(x, r) \prod_{l=1}^{k-1} \Delta\left(x^{2 l+1}\right) \delta\left(\frac{x^{j-i+2 k} z_{2}}{z_{1}}\right) f_{i-k+1, j+k}\left(x^{i-j+1}\right) T_{i-k+1}\left(x^{-k-1} z_{1}\right) \\
& \quad \times T_{j+k}\left(x^{k} z_{2}\right), \quad 1 \leq k \leq i-1 . \tag{4.61}
\end{align*}
$$

Adding (4.55), (4.57), (4.59), and (4.61) for $1 \leq k \leq i-1$, and replacing $z_{1}$ by $x z_{1}$ yields $\operatorname{LHS}_{i+1, j}=\operatorname{RHS1}_{i+1, j}+\sum_{k=1}^{i+1} \mathrm{RHS}_{i+1, j}(k)$. By induction on $i$, we proved quadratic relation (3.4).

### 4.3 Proof of Lemma 3.4

Lemma 4.9. The current $T_{1}(z)$ commutes with the screening currents $S_{k}(w)$ as follows.

$$
\begin{equation*}
\left[T_{1}(z), S_{k}(w)\right]=C_{k}(z)\left(D_{x^{r}} \delta\right)\left(\frac{x^{k} w}{z}\right)+\bar{C}_{k}(z)\left(D_{x^{r}} \delta\right)\left(\frac{x^{2 N+1-k} w}{z}\right), \quad 1 \leq k \leq N \tag{4.62}
\end{equation*}
$$

Here we set $q$-difference

$$
\left(D_{q} \delta\right)(z)=\delta(q z)-\delta\left(q^{-1} z\right)
$$

the currents $C_{k}(z)$ and $\bar{C}_{k}(z), 1 \leq k \leq N$, are given by

$$
\begin{aligned}
& C_{k}(z)=x^{-r+1}\left(x^{r-1}-x^{-r+1}\right): \Lambda_{k}(z) S_{k}\left(x^{r-k} z\right): \\
& \bar{C}_{k}(z)=x^{r-1}\left(x^{r-1}-x^{-r+1}\right): \Lambda_{\bar{k}}(z) S_{k}\left(x^{-2 N-1+k-r} z\right):
\end{aligned}
$$

Proof. Adding (B.9) yields

$$
\begin{aligned}
{\left[T_{1}(z), S_{k}(w)\right]=} & \left(x^{r-1}-x^{-r+1}\right)\left(-x^{-r+1}: \Lambda_{k}(z) S_{k}(w): \delta\left(x^{k-r} \frac{w}{z}\right)\right. \\
& +x^{r-1}: \Lambda_{k+1}(z) S_{k}(w) \delta\left(x^{k+r} \frac{w}{z}\right)+x^{r-1}: \Lambda_{\bar{k}}(z) S_{k}(w) \delta\left(x^{2 N+1-k+r} \frac{w}{z}\right) \\
& \left.-x^{-r+1}: \Lambda_{\overline{k+1}}(z) S_{k}(w): \delta\left(x^{2 N+1-k-r} \frac{w}{z}\right)\right), \quad 1 \leq k \leq N-1,
\end{aligned}
$$

and

$$
\begin{aligned}
{\left[T_{1}(z), S_{N}(w)\right]=} & \left(x^{-r+1}-x^{r-1}\right)\left(x^{-r+1}: \Lambda_{N}(z) S_{N}(w): \delta\left(x^{N-r} \frac{w}{z}\right)\right. \\
& \left.-x^{r-1}: \Lambda_{N}(z) S_{\bar{N}}(w): \delta\left(x^{N+1+r} \frac{w}{z}\right)\right) \\
& +\frac{[r-1]_{x}\left[\frac{1}{2}\right]_{x}}{\left[r-\frac{1}{2}\right]_{x}}\left(x-x^{-1}\right)\left(\delta\left(x^{N+r} \frac{w}{z}\right)-\delta\left(x^{N+1-r} \frac{w}{z}\right)\right): \Lambda_{0}(z) S_{N}(w): .
\end{aligned}
$$

Using the relations

$$
\begin{aligned}
& x^{-r+1}: \Lambda_{k}(z) S_{k}\left(x^{r-k} z\right):=x^{r-1}: \Lambda_{k+1}(z) S_{k}\left(x^{-r-k} z\right):, \quad 1 \leq k \leq N-1, \\
& x^{-r+1}: \Lambda_{N}(z) S_{N}\left(x^{r-N} z\right):=\frac{\left[\frac{1}{2}\right]_{x}}{\left[r-\frac{1}{2}\right]_{x}}: \Lambda_{0}(z) S_{N}\left(x^{-r-N} z\right):, \\
& x^{r-1}: \Lambda_{\bar{k}}(z) S_{k}\left(x^{-2 N-1+k-r} z\right):=x^{-r+1}: \Lambda_{\overline{k+1}}(z) S_{k}\left(x^{-2 N-1+k+r} z\right):, \quad 1 \leq k \leq N-1, \\
& x^{r-1}: \Lambda_{\bar{N}}(z) S_{N}\left(x^{-r-N-1} z\right):=\frac{\left[\frac{1}{2}\right]_{x}}{\left[r-\frac{1}{2}\right]_{x}}: \Lambda_{0}(z) S_{N}\left(x^{r-N-1} z\right): .
\end{aligned}
$$

yields (4.62).
Corollary 4.10. The current $T_{1}(z)$ commutes with the screening operators $S_{k}$

$$
\begin{equation*}
\left[T_{1}(z), S_{k}\right]=0, \quad 1 \leq k \leq N . \tag{4.63}
\end{equation*}
$$

Proof. From (4.62), we obtain

$$
\left[T_{1}(z), S_{k}\right]=\oint \frac{\mathrm{d} w}{2 \pi \sqrt{-1} w}\left(C_{k}(z)\left(D_{x^{r}} \delta\right)\left(\frac{x^{k} w}{z}\right)+\bar{C}_{k}(z)\left(D_{x^{r}} \delta\right)\left(\frac{x^{2 N+1-k} w}{z}\right)\right) .
$$

Using $\oint \frac{\mathrm{d} w}{2 \pi \sqrt{-1} w}\left(D_{x^{r}} \delta\right)\left(\frac{x^{s} w}{z}\right)=0$ with $s=k, 2 N+1-k$ yields $\left[T_{1}(z), S_{k}\right]=0$.
Proof. Here we will give a proof of Lemma 3.4. Set $T_{j}(z)=\sum_{m \in \mathbb{Z}} T_{j}[m] z^{-m}, 1 \leq j \leq 2 N$ and $f_{i, j}(z)=\sum_{l=0}^{\infty} f_{i, j}^{l} z^{l}$. From (4.21), we obtain

$$
\begin{aligned}
& \left(x^{-(j+1) k+m}-x^{(j+1) k-m}\right) T_{j+1}[m] \\
& \quad=\Delta\left(x^{2 N+2-2 j}\right)\left(x^{(2 N-j+2) k-m}-x^{(-2 N+j-2) k+m}\right) T_{j-1}[m] \\
& \quad+c(x, r)^{-1} \sum_{l=0}^{\infty}\left(f_{1, j}^{l} T_{1}[k-l] T_{j}[l-k+m]-f_{j, 1}^{l} T_{j}[k-l-m] T_{1}[l-k]\right), \\
& m, k \in \mathbb{Z}, \quad 1 \leq j \leq N .
\end{aligned}
$$

Hence, $T_{j+1}[m], m \in \mathbb{Z}, 1 \leq j \leq N$, are expressed in terms of $T_{j}[n], T_{j-1}[n]$, and $T_{1}[n], n \in \mathbb{Z}$, $1 \leq j \leq N$. From duality (3.3), $T_{j}[m], m \in \mathbb{Z}, N+2 \leq j \leq 2 N$ are expressed in terms of $T_{2 N+1-j}[n], n \in \mathbb{Z}, N+2 \leq j \leq 2 N$. Finally, $T_{j}[m], m \in \mathbb{Z}, 1 \leq j \leq 2 N$ are expressed in terms of $T_{1}[n], n \in \mathbb{Z}$. Hence, we obtain (3.5) from (4.63).

## 5 Conclusion and discussion

In this paper, we obtained the free field construction of higher $W$-currents $T_{i}(z), i \geq 2$, of the deformed $W$-algebra $\mathcal{W}_{x, r}\left(A_{2 N}^{(2)}\right)$. We obtained a closed set of quadratic relations for the $W$ currents $T_{i}(z)$, which are completely different from those in types $A_{N}^{(1)}$ and $A(M, N)^{(1)}$. The quadratic relations of $\mathcal{W}_{x, r}\left(A_{2 N}^{(2)}\right)$ do not preserve "parity", though those of $\mathcal{W}_{x, r}\left(A_{N}^{(1)}\right)$ and $\mathcal{W}_{x, r}\left(A(M, N)^{(1)}\right)$ do. Here we define "parity" of $T_{i}(z) T_{j}(w)$ as $i+j$. We obtained the duality $T_{2 N+1-i}(z)=c_{i} T_{i}(z), 1 \leq i \leq N$, which is a new structure that does not occur in types $A_{2}^{(2)}, A_{N}^{(1)}$, and $A(M, N)^{(1)}$. This allowed us to define the deformed $W$-algebra $\mathcal{W}_{x, r}\left(A_{2 N}^{(2)}\right)$ using generators and relations similarly to the definition of the twisted affine Lie algebra of type $A_{2 N}^{(2)}$ given in Section 2.

We also justified our definition of the deformed $W$-algebra of type $A_{2 N}^{(2)}$. We compare Definition 3.3 with other definitions. In [8], the deformed $W$-algebras of types $A_{N}^{(1)}, B_{N}^{(1)}, C_{N}^{(1)}, D_{N}^{(1)}$, and $A_{2 N}^{(2)}$ were proposed as the intersection of the kernels of the screening operators. We recall the definition based on the screening operators for $A_{2 N}^{(2)}$. Let $\mathbf{H}_{x, r}$ be the vector space spanned by the formal power series currents of the form

$$
: \partial_{z}^{n_{1}} Y_{i_{1}}\left(x^{r j_{1}+k_{1}} z\right)^{\varepsilon_{1}} \cdots \partial_{z}^{n_{l}} Y_{i_{l}}\left(x^{r j_{l}+k_{l}} z\right)^{\varepsilon_{l}}:
$$

where $\varepsilon_{i}= \pm 1^{2}$. We define $\mathbf{W}_{x, r}$ as the vector subspace of $\mathbf{H}_{x, r}$ consisting of all currents that commute with the screening operators $S_{i}, 1 \leq i \leq N$, in (2.6). Let $\left\{F_{a}(z)=\sum_{m \in \mathbb{Z}} F_{a}[m] z^{-m}\right\}_{a \in A}$ be a basis of the vector space $\mathbf{W}_{x, r}$. Let $\mathcal{W}^{\mathrm{FR}}$ be the associative algebra generated by elements $F_{a}[m], m \in \mathbb{Z}, a \in A$. Let $J_{K}$ be the left ideal of $\mathcal{W}^{\mathrm{FR}}$ generated by elements $F_{a}[m]$, $m \geq K \in \mathbb{N}, a \in A$. We define the deformed $W$-algebra

$$
\mathcal{W}_{x, r}^{\mathrm{FR}}\left(A_{2 N}^{(2)}\right)=\lim _{\leftarrow} \mathcal{W}^{\mathrm{FR}} / J_{K}
$$

We propose another definition of the deformed $W$-algebra. From (3.5), the $W$-currents $T_{i}(z)=\sum_{m \in \mathbb{Z}} T_{i}[m] z^{-m}, 1 \leq i \leq 2 N$, commute with the screening operators. Let $\mathcal{W}^{\text {AKOS }}$ be the associative algebra generated by elements $T_{i}[m], m \in \mathbb{Z}, 1 \leq i \leq 2 N$. Let $L_{K}$ be the left ideal of $\mathcal{W}^{\text {AKOS }}$ generated by elements $T_{i}[m], m \geq K \in \mathbb{N}, 1 \leq i \leq 2 N$. We define the deformed $W$-algebra

$$
\mathcal{W}_{x, r}^{\mathrm{AKOS}}\left(A_{2 N}^{(2)}\right)=\lim _{\leftarrow} \mathcal{W}^{\mathrm{AKOS}} / L_{K} .
$$

In this study, our definitions $\mathcal{W}_{x, r}\left(A_{2 N}^{(2)}\right)$ were based on generators and relations. We have introduced three definitions of the deformed $W$-algebra for the twisted algebra of the type $A_{2 N}^{(2)}$. Conjecture 5.1. $\mathcal{W}_{x, r}\left(A_{2 N}^{(2)}\right)$, $\mathcal{W}_{x, r}^{\mathrm{AKOS}}\left(A_{2 N}^{(2)}\right)$, and $\mathcal{W}_{x, r}^{\mathrm{FR}}\left(A_{2 N}^{(2)}\right)$ are isomorphic as associative algebras

$$
\begin{equation*}
\mathcal{W}_{x, r}\left(A_{2 N}^{(2)}\right) \cong \mathcal{W}_{x, r}^{\mathrm{AKOS}}\left(A_{2 N}^{(2)}\right) \cong \mathcal{W}_{x, r}^{\mathrm{FR}}\left(A_{2 N}^{(2)}\right) \tag{5.1}
\end{equation*}
$$

The author believes that this conjecture can be extended to arbitrary affine Lie algebras. Some necessary conditions of isomorphism (5.1) in Conjecture 5.1 can be indicated immediately. From (3.5), we obtain the following inclusion:

$$
\mathcal{W}_{x, r}^{\mathrm{AKOS}}\left(A_{2 N}^{(2)}\right) \subseteq \mathcal{W}_{x, r}^{\mathrm{FR}}\left(A_{2 N}^{(2)}\right)
$$

[^1]We establish a homomorphism of associative algebras $\varphi \in \operatorname{Hom}_{\mathbf{C}}\left(\mathcal{W}_{x, r}\left(A_{2 N}^{(2)}\right), \mathcal{W}_{x, r}^{\mathrm{AKOS}}\left(A_{2 N}^{(2)}\right)\right)$ using $\varphi\left(\bar{T}_{i}[m]\right)=T_{i}[m] . \varphi$ is surjective,

$$
\varphi\left(\mathcal{W}_{x, r}\left(A_{2 N}^{(2)}\right)\right)=\mathcal{W}_{x, r}^{\mathrm{AKOS}}\left(A_{2 N}^{(2)}\right)
$$

If we assume that $\varphi$ is injective, the isomorphism on the left side in (5.1) is obtained. In other words, no independent relations other than (3.3) and (3.4) exist in $\mathcal{W}_{x, r}\left(A_{2 N}^{(2)}\right)$. We propose two results to support this claim. In the classical limit the second Hamiltonian structure $\{\cdot, \cdot\}$ of the $q$-Poisson algebra [7, 8, 9, 15] was obtained from the quadratic relations (see (3.6) and (3.7)). In the conformal limit all defining relations of the $W$-algebra $\mathcal{W}_{\beta}\left(A_{N}^{(1)}\right), N=1,2$, are obtained from the quadratic relations of $\mathcal{W}_{x, r}\left(A_{N}^{(1)}\right)$ upon the assumption that the currents $T_{i}(z)$ have the form of expansion for small parameter $\hbar$ (see [ 1 , Appendix]).

The definition of the deformed $W$-algebra $\mathcal{W}_{x, r}(\mathfrak{g})$ for non-twisted affine Lie algebra $\mathfrak{g}$ was formulated in terms of the quantum Drinfeld-Sokolov reduction in [16]. Formulating the definition of the deformed $W$-algebras $\mathcal{W}_{x, r}(\mathfrak{g})$ in terms of the quantum Drinfeld-Sokolov reduction for twisted affine Lie algebra or affine Lie superalgebra [4, 6, 10, 12, 13] is still a problem that needs to be solved.

It remains an open challenge to identify quadratic relations of the deformed $W$-algebras $\mathcal{W}_{x, r}(\mathfrak{g})$ for the affine Lie algebras $\mathfrak{g}$ except for types $A_{N}^{(1)}$ and $A_{2 N}^{(2)}$. We believe that this paper presents a key step towards extending our construction for general affine Lie algebras $\mathfrak{g}$. In [8] and [6] the free field construction of the basic $W$-current $T_{1}(z)$ of $\mathcal{W}_{x, r}(\mathfrak{g})$ was suggested in the case when the underlying simple finite-dimensional Lie algebra $\mathfrak{g}$ is of classical type,

$$
T_{1}(z)= \begin{cases}\Lambda_{1}(z)+\cdots \Lambda_{N+1}(z) & \text { for } \mathfrak{g} \text { of type } A_{N}^{(1)}, \\ \Lambda_{1}(z)+\cdots+\Lambda_{N}(z)+\Lambda_{0}(z) & \text { for } \mathfrak{g} \text { of types } B_{N}^{(1)}, A_{2 N}^{(2)}, D_{N+1}^{(2)}, \\ \quad+\Lambda_{\bar{N}}(z)+\cdots+\Lambda_{\overline{1}}(z) & \text { for } \mathfrak{g} \text { of types } C_{N}^{(1)}, D_{N}^{(1)}, A_{2 N-1}^{(2)} .\end{cases}
$$

Here we omit details of free field constructions of $\Lambda_{i}(z)$. The free field construction of $T_{1}(z)$ has similar form to that for $\mathfrak{g}$ of type $A_{2 N}^{(2)}$ except for the case of $A_{N}^{(1)}$. Therefore, we expect that a similar duality as (3.3) and similar quadratic relation (3.4) hold in all cases in types $B_{N}^{(1)}, C_{N}^{(1)}, D_{N}^{(1)}, A_{2 N-1}^{(2)}$, and $D_{N+1}^{(2)}$. We would like to draw your attention to the following analogy. Let $\mathfrak{g}$ be an affine Lie algebras of one of the types $B_{N}^{(1)}, C_{N}^{(1)}, D_{N}^{(1)}, A_{2 N-1}^{(2)}$, or $D_{N+1}^{(2)}$. Let $\mathfrak{g}$ be the underlying simple finite-dimensional Lie algebra. Let $\mathfrak{h}$ be a Cartan subalgebra of $\stackrel{\circ}{\mathfrak{g}}$. Let $\bar{\Lambda}_{1}, \bar{\Lambda}_{2}, \ldots, \bar{\Lambda}_{l}$ be the fundamental weights of $\stackrel{\circ}{\mathfrak{g}}$, where $l$ is the dimension of $\mathfrak{\circ}$. Let $V_{\bar{\Lambda}_{1}}$ be the integrable highest weight representation of $U_{q}(\mathfrak{\circ})$ with the highest weight $\bar{\Lambda}_{1}$. Let $V$ be the evaluation representation corresponding to $V_{\bar{\Lambda}_{1}}$ of the quantum affine algebra $U_{q}(\mathfrak{g})$ with a spectral parameter $z \in \mathbf{C}^{\times}$. Let $n$ be the dimension of $V_{\bar{\Lambda}_{1}}$. We have $\wedge^{n-i} V \simeq(\stackrel{i}{\wedge} V)^{*} \simeq \wedge^{i} V^{*}$, because $\wedge^{n} V \simeq \mathbf{C}$. The evaluation representation $V$ of $U_{q}(\mathfrak{g})$ is self-dual except for $\mathfrak{g}$ of type $A_{N}^{(1)}$. Hence, we obtain the duality of the representations of $U_{q}(\mathfrak{g})$,

$$
{ }^{n-i} V \simeq \wedge^{i} V \quad \text { if } \quad \mathfrak{g} \text { is not of type } A_{N}^{(1)},
$$

which is similar as that in (3.3). As an analogy, we expect the duality of the $W$-currents,

$$
T_{n-i}(z)=c_{i} T_{i}(z) \quad \text { if } \quad \mathfrak{g} \text { is not of type } A_{N}^{(1)},
$$

for the deformed $W$ algebras $\mathcal{W}_{x, r}(\mathfrak{g})$. Here $c_{i}, 0 \leq i \leq n$, are constants.

It remains an open challenge to identify quadratic relations of the deformed $W$-algebras $\mathcal{W}_{x, r}(\mathfrak{g})$ for affine superalgebra $\mathfrak{g}$ except for those of type $A(M, N)^{(1)}$. Recently the deformed $W$-superalgebra $\mathcal{W}_{x, r}(\mathfrak{g})$ has appeared in the study of D-branes and physical interest is growing to this subject, see, e.g., [10]. As revealed in [10, 12, 13], it is expected that, in cases of superalgebras $\mathfrak{g}$, infinite number of higher $W$-currents $T_{i}(z), i=1,2,3, \ldots$, satisfy a closed set of infinite number of quadratic relations. It is interesting to understand how duality will be extended to the case of superalgebras. We expect to report on quadratic relations and duality for more general deformed $W$-algebras $\mathcal{W}_{x, r}(\mathfrak{g})$ associated with affine Lie algebras and affine Lie superalgebras in the near future.

## A Normal ordering rules

We list the normal ordering rules. For operators $V(z)$ and $W(w)$ we use the notation

$$
\begin{equation*}
V(z) W(w)=\langle V(z) W(w)\rangle: V(z) W(w) \tag{A.1}
\end{equation*}
$$

and write down only the part $\langle V(z) W(w)\rangle$ in the formulas below. Using the standard formula

$$
\mathrm{e}^{A} \mathrm{e}^{B}=\mathrm{e}^{[A, B]} \mathrm{e}^{B} \mathrm{e}^{A}, \quad[[A, B], A]=0 \quad \text { and } \quad[[A, B], B]=0
$$

we obtain the normal ordering rules.

## A. $1 \quad A_{i}(z)$ and $S_{i}(z)$

$$
\begin{align*}
& \left\langle A_{i}\left(z_{1}\right) A_{i}\left(z_{2}\right)\right\rangle=\left(\Delta\left(\frac{x z_{2}}{z_{1}}\right) \Delta\left(\frac{x^{-1} z_{2}}{z_{1}}\right)\right)^{-1}, \quad 1 \leq i \leq N-1, \\
& \left\langle A_{N}\left(z_{1}\right) A_{N}\left(z_{2}\right)\right\rangle=\Delta\left(\frac{z_{2}}{z_{1}}\right)\left(\Delta\left(\frac{x z_{2}}{z_{1}}\right) \Delta\left(\frac{x^{-1} z_{2}}{z_{1}}\right)\right)^{-1}, \\
& \left\langle A_{i}\left(z_{1}\right) A_{j}\left(z_{2}\right)\right\rangle=\Delta\left(\frac{z_{2}}{z_{1}}\right), \quad|i-j|=1, \quad 1 \leq i, j \leq N, \\
& \left\langle A_{i}\left(z_{1}\right) A_{j}\left(z_{2}\right)\right\rangle=1, \quad|i-j| \geq 2, \quad 1 \leq i, j \leq N,  \tag{A.2}\\
& \left\langle S_{i}\left(z_{1}\right) S_{i}\left(z_{2}\right)\right\rangle=z_{1}^{\frac{2(r-1)}{r}}\left(1-\frac{z_{2}}{z_{1}}\right) \frac{\left(x^{2} z_{2} / z_{1} ; x^{2 r}\right)_{\infty}}{\left(x^{2 r-2} z_{2} / z_{1} ; x^{2 r}\right)_{\infty}}, \quad 1 \leq i \leq N-1, \\
& \left\langle S_{N}\left(z_{1}\right) S_{N}\left(z_{2}\right)\right\rangle=z_{1}^{\frac{r-1}{r}}\left(1-\frac{z_{2}}{z_{1}}\right) \frac{\left(x^{2} z_{2} / z_{1} ; x^{2 r}\right)_{\infty}\left(x^{2 r-2} z_{2} / z_{1} ; x^{2 r}\right)_{\infty}}{\left(x z_{2} / z_{1} ; x^{2 r}\right)_{\infty}\left(x^{2 r-1} z_{2} / z_{1} ; x^{2 r}\right)_{\infty}} \\
& \left\langle S_{i}\left(z_{1}\right) S_{j}\left(z_{2}\right)\right\rangle=z_{1}^{-\frac{r-1}{r}} \frac{\left(x^{2 r-1} z_{2} / z_{1} ; x^{2 r}\right)_{\infty},}{\left(x z_{2} / z_{1} ; x^{2 r}\right)_{\infty}}, \quad|i-j|=1, \quad 1 \leq i, j \leq N, \\
& \left\langle S_{i}\left(z_{1}\right) S_{j}\left(z_{2}\right)\right\rangle=1, \quad|i-j| \geq 2, \quad 1 \leq i, j \leq N,  \tag{A.3}\\
& \left\langle A_{i}\left(z_{1}\right) S_{i}\left(z_{2}\right)\right\rangle=x^{-4(r-1)} \frac{\left(1-x^{r} \frac{z_{2}}{z_{1}}\right)\left(1-x^{r-2} \frac{z_{2}}{z_{1}}\right)}{\left(1-x^{-r} \frac{z_{2}}{z_{1}}\right)\left(1-x^{2-r} \frac{z_{2}}{z_{1}}\right)}, \quad 1 \leq i \leq N-1, \\
& \left\langle A_{N}\left(z_{1}\right) S_{N}\left(z_{2}\right)\right\rangle=x^{-2(r-1)} \frac{\left(1-x^{r} \frac{z_{2}}{z_{1}}\right)\left(1-x^{r-2} \frac{z_{2}}{z_{1}}\right)\left(1-x^{1-r} \frac{z_{2}}{z_{1}}\right)}{\left(1-x^{-r} \frac{z_{2}}{z_{1}}\right)\left(1-x^{2-r} \frac{z_{2}}{z_{1}}\right)\left(1-x^{r-1} \frac{z_{2}}{z_{1}}\right)},  \tag{A.4}\\
& \left\langle A_{i}\left(z_{1}\right) S_{j}\left(z_{2}\right)\right\rangle=x^{2(r-1)} \frac{\left(1-x^{1-r} \frac{z_{2}}{z_{1}}\right)}{\left(1-x^{r-1} \frac{z_{2}}{z_{1}}\right)},|i-j|=1, \quad 1 \leq i, j \leq N, \\
& \left\langle A_{i}\left(z_{1}\right) S_{j}\left(z_{2}\right)\right\rangle=1, \quad|i-j| \geq 2, \quad 1 \leq i, j \leq N, \tag{A.5}
\end{align*}
$$

$$
\begin{align*}
& \left\langle S_{i}\left(z_{1}\right) A_{i}\left(z_{2}\right)\right\rangle=\frac{\left(1-x^{2-r} \frac{z_{2}}{z_{1}}\right)}{\left(1-x^{r} \frac{z_{2}}{z_{1}}\right)\left(1-x^{r-2} \frac{z_{2}}{z_{1}}\right)}, \quad 1 \leq i \leq N-1, \\
& \left\langle S_{N}\left(z_{1}\right) A_{N}\left(z_{2}\right)\right\rangle=\frac{\left(1-x^{-r} \frac{z_{2}}{z_{1}}\right)\left(1-x^{2-r} \frac{z_{2}}{z_{1}}\right)\left(1-x^{r-1} \frac{z_{2}}{z_{1}}\right)}{\left(1-x^{r} \frac{z_{2}}{z_{1}}\right)\left(1-x^{r-2} \frac{z_{2}}{z_{1}}\right)\left(1-x^{1-r} \frac{z_{2}}{z_{1}}\right)}, \quad|i-j|=1, \quad 1 \leq i, j \leq N, \\
& \left\langle S_{j}\left(z_{1}\right) A_{i}\left(z_{2}\right)\right\rangle=\frac{\left(1-x^{r-1} \frac{z_{2}}{z_{1}}\right)}{\left(1-x^{1-r \frac{z_{2}}{z_{1}}}\right)}, \quad|i-j| \geq 2, \quad 1 \leq i, j \leq N .
\end{align*}
$$

## A. $2 \quad Y_{i}(z), A_{i}(z)$ and $S_{i}(z)$

The symmetric matrix $I(m)=\left(I_{i, j}(m)\right)_{i, j=1}^{N}$ is the inverse matrix of $B(m)$. The elements $I_{i, j}(m)=I_{j, i}(m), 1 \leq i \leq j \leq N$, are written as

$$
I_{i, j}(m)=\frac{1}{[(N+1) m]_{x}-[N m]_{x}} \times \begin{cases}{[(N+1-j) m]_{x}-[(N-j) m]_{x},} & i=1,1 \leq j \leq N \\ (-1)^{N-j+i} \sum_{k=i-1}^{N-j+i}(-1)^{k}[k m]_{x}, & 2 \leq i \leq j \leq N-1 \\ {[i m]_{x},} & 1 \leq i \leq N, j=N\end{cases}
$$

The generators $y_{i}(m), 1 \leq i \leq N$, are written as

$$
\begin{equation*}
y_{i}(m)=\sum_{j=1}^{N} I_{i, j}(m) a_{j}(m), \quad Q_{i}^{y}=\sum_{j=1}^{N} I_{i, j}(0) Q_{j} . \tag{A.7}
\end{equation*}
$$

From (2.2), (2.3) and (A.2) we obtain

$$
\begin{align*}
& \left\langle Y_{1}\left(z_{1}\right) Y_{1}\left(z_{2}\right)\right\rangle=f_{1,1}\left(\frac{z_{2}}{z_{1}}\right)^{-1}, \\
& \left\langle Y_{1}\left(z_{1}\right) A_{1}\left(z_{2}\right)\right\rangle=\Delta\left(\frac{z_{2}}{z_{1}}\right)^{-1}, \quad\left\langle Y_{1}\left(z_{1}\right) A_{i}\left(z_{2}\right)\right\rangle=1, \quad 2 \leq i \leq N, \\
& \left\langle A_{1}\left(z_{1}\right) Y_{1}\left(z_{2}\right)\right\rangle=\Delta\left(\frac{z_{2}}{z_{1}}\right)^{-1}, \quad\left\langle A_{i}\left(z_{1}\right) Y_{1}\left(z_{2}\right)\right\rangle=1, \quad 2 \leq i \leq N, \\
& \left\langle Y_{1}\left(z_{1}\right) S_{1}\left(z_{2}\right)\right\rangle=x^{-2(r-1)} \frac{\left(1-x^{r-1} \frac{z_{2}}{z_{1}}\right)}{\left(1-x^{1-r} \frac{z_{2}}{z_{1}}\right)}, \quad\left\langle Y_{1}\left(z_{1}\right) S_{i}\left(z_{2}\right)\right\rangle=1, \quad 2 \leq i \leq N, \\
& \left\langle S_{1}\left(z_{1}\right) Y_{1}\left(z_{2}\right)\right\rangle=\frac{\left(1-x^{\left.1-r \frac{z_{2}}{z_{1}}\right)}\right.}{\left(1-x^{\left.r-\frac{z_{2}}{z_{1}}\right)}, \quad\left\langle S_{i}\left(z_{1}\right) Y_{1}\left(z_{2}\right)\right\rangle=1, \quad 2 \leq i \leq N .\right.} . \tag{A.8}
\end{align*}
$$

## B Exchange relations

In this appendix we list the exchange relations.

## B. $1 \quad \Lambda_{i}(z)$

We give the exchange relations of $\Lambda_{j}(z)$ and $\vec{\Lambda}_{\Omega_{i}}(z)$, which are obtained from (4.1). We set $s \in$ $J_{N}=\{1,2, \ldots, N, 0, \bar{N}, \ldots, \overline{2}, \overline{1}\}$. For an element $s \in J_{N}$ and a subset $\Omega_{i}=\left\{s_{1}, s_{2}, \ldots, s_{i}\right\} \subset J_{N}$ with $s_{1} \prec s_{2} \prec \cdots \prec s_{i}$, we calculate

$$
X_{\Omega_{i}, s}\left(z_{1}, z_{2}\right)=f_{1, i}\left(\frac{z_{2}}{z_{1}}\right) \Lambda_{s}\left(z_{1}\right) \vec{\Lambda}_{\Omega_{i}}\left(z_{2}\right)-f_{i, 1}\left(\frac{z_{1}}{z_{2}}\right) \vec{\Lambda}_{\Omega_{i}}\left(z_{2}\right) \Lambda_{s}\left(z_{1}\right)
$$

- In the case of $s, \bar{s} \notin \Omega_{i}$, we obtain

$$
\begin{equation*}
X_{\Omega_{i}, s}\left(z_{1}, z_{2}\right)=c(x, r): \Lambda_{s}\left(z_{1}\right) \vec{\Lambda}_{\Omega_{i}}\left(z_{2}\right):\left(\delta\left(\frac{x^{-i-3+2 k} z_{2}}{z_{1}}\right)-\delta\left(\frac{x^{-i-1+2 k} z_{2}}{z_{1}}\right)\right) \tag{B.1}
\end{equation*}
$$

Here we set $k, 1 \leq k \leq i+1$, by $k=\left\{\begin{array}{lll}1 & \text { if } \quad s \prec s_{1}, \\ q & \text { if } & s_{q-1} \prec s \prec s_{q}, \quad 2 \leq q \leq i, \\ i+1 & \text { if } \quad s_{i} \prec s .\end{array}\right.$

- In the case of $s \in \Omega_{i}$ and $\bar{s} \notin \Omega_{i}$, we obtain

$$
\begin{equation*}
f_{1, i}\left(\frac{z_{2}}{z_{1}}\right) \Lambda_{s}\left(z_{1}\right) \vec{\Lambda}_{\Omega_{i}}\left(z_{2}\right)-f_{i, 1}\left(\frac{z_{1}}{z_{2}}\right) \vec{\Lambda}_{\Omega_{i}}\left(z_{2}\right) \Lambda_{s}\left(z_{1}\right)=0 \tag{B.2}
\end{equation*}
$$

- In the case of $s, \bar{s} \in \Omega_{i}$ and $s=n, 1 \leq n \leq N$, we obtain

$$
\begin{align*}
X_{\Omega_{i}, s}\left(z_{1}, z_{2}\right)= & c(x, r): \Lambda_{n}\left(z_{1}\right) \vec{\Lambda}_{\Omega_{i}}\left(z_{2}\right): \\
& \times\left(\delta\left(\frac{x^{-2 N-i+2 n+2 l-4} z_{2}}{z_{1}}\right)-\delta\left(\frac{x^{-2 N-i+2 n+2 l-2} z_{2}}{z_{1}}\right)\right) . \tag{B.3}
\end{align*}
$$

Here we set $k, l, 1 \leq k<l \leq i$, by $s=n=s_{k}$ and $\bar{s}=\bar{n}=s_{l}$.

- In the case of $s, \bar{s} \in \Omega_{i}$ and $s=\bar{n}, 1 \leq n \leq N$, we obtain

$$
\begin{align*}
X_{\Omega_{i}, s}\left(z_{1}, z_{2}\right)= & c(x, r): \Lambda_{\bar{n}}\left(z_{1}\right) \vec{\Lambda}_{\Omega_{i}}\left(z_{2}\right): \\
& \times\left(\delta\left(\frac{x^{2 N-i-2 n+2 k} z_{2}}{z_{1}}\right)-\delta\left(\frac{x^{2 N-i-2 n+2 k+2} z_{2}}{z_{1}}\right)\right) . \tag{B.4}
\end{align*}
$$

Here we set $k, l, 1 \leq k<l \leq i$, by $\bar{s}=n=s_{k}$ and $s=\bar{n}=s_{l}$.

- In the case of $s=0 \in \Omega_{i}$, we obtain

$$
\begin{equation*}
X_{\Omega_{i}, s}\left(z_{1}, z_{2}\right)=c(x, r): \Lambda_{0}\left(z_{1}\right) \vec{\Lambda}_{\Omega_{i}}\left(z_{2}\right):\left(\delta\left(\frac{x^{-i-2+2 k} z_{2}}{z_{1}}\right)-\delta\left(\frac{x^{-i+2 k} z_{2}}{z_{1}}\right)\right) \tag{B.5}
\end{equation*}
$$

Here we set $k, 1 \leq k \leq i$, by $s_{k}=0$.

- In the case of $s \notin \Omega_{i}$ and $\bar{s} \in \Omega_{i}$ and $s=n, 1 \leq n \leq N$, we obtain

$$
\begin{align*}
X_{\Omega_{i}, s}\left(z_{1}, z_{2}\right)= & c(x, r) \Delta\left(x^{2(l-k+n-N)}\right): \Lambda_{n}\left(z_{1}\right) \vec{\Lambda}_{\Omega_{i}}\left(z_{2}\right): \\
& \times\left(\delta\left(\frac{x^{-i+2 k-3} z_{2}}{z_{1}}\right)-\delta\left(\frac{x^{-2 N+2 n+2 l-i-2} z_{2}}{z_{1}}\right)\right) \\
+ & c(x, r) \Delta\left(x^{2(l-k+n-N-1)}\right): \Lambda_{n}\left(z_{1}\right) \vec{\Lambda}_{\Omega_{i}}\left(z_{2}\right): \\
& \times\left(-\delta\left(\frac{x^{-i+2 k-1} z_{2}}{z_{1}}\right)+\delta\left(\frac{x^{-2 N+2 n+2 l-i-4} z_{2}}{z_{1}}\right)\right) . \tag{B.6}
\end{align*}
$$

Here we set $k, l, 1 \leq k \leq l \leq i$, by $s_{l}=\bar{s}=\bar{n}$ and $k=\left\{\begin{array}{lll}1 & \text { if } & s=n \prec s_{1}, \\ q & \text { if } & s_{q-1} \prec s=n \prec s_{q}, \quad 2 \leq q \leq i\end{array}\right.$

- In the case of $s \notin \Omega_{i}$ and $\bar{s} \in \Omega_{i}$ and $s=\bar{n}, 1 \leq n \leq N$, we obtain

$$
\begin{align*}
X_{\Omega_{i}, s}\left(z_{1}, z_{2}\right)= & c(x, r) \Delta\left(x^{2(l-k+n-N-1)}\right): \Lambda_{\bar{n}}\left(z_{1}\right) \vec{\Lambda}_{\Omega_{i}}\left(z_{2}\right): \\
& \times\left(\delta\left(\frac{x^{-i+2 l-1} z_{2}}{z_{1}}\right)-\delta\left(\frac{x^{2 N-2 n-i+2 k+2} z_{2}}{z_{1}}\right)\right) \\
+ & c(x, r) \Delta\left(x^{2(l-k+n-N)}\right): \Lambda_{\bar{n}}\left(z_{1}\right) \vec{\Lambda}_{\Omega_{i}}\left(z_{2}\right): \\
& \times\left(-\delta\left(\frac{x^{-i+2 l+1} z_{2}}{z_{1}}\right)+\delta\left(\frac{x^{2 N-2 n-i+2 k} z_{2}}{z_{1}}\right)\right) . \tag{B.7}
\end{align*}
$$

Here we set $k, l, 1 \leq k \leq l \leq i$, by $s_{k}=\bar{s}=n$ and $l=\left\{\begin{array}{lll}q & \text { if } & s_{q} \prec s=\bar{n} \prec s_{q+1}, \quad 1 \leq q \leq i-1, \\ i & \text { if } & s_{i} \prec s=\bar{n} .\end{array}\right.$

## B. $2 S_{i}(z)$

From (A.3) we obtain

$$
\begin{align*}
& S_{i}\left(z_{1}\right) S_{i}\left(z_{2}\right)=-\frac{\left[u_{2}-u_{1}+1\right]}{\left[u_{1}-u_{2}+1\right]} S_{i}\left(z_{2}\right) S_{i}\left(z_{1}\right), \quad 1 \leq i \leq N-1 \\
& S_{N}\left(z_{1}\right) S_{N}\left(z_{2}\right)=-\frac{\left[u_{1}-u_{2}+\frac{1}{2}\right]\left[u_{2}-u_{1}+1\right]}{\left[u_{2}-u_{1}+\frac{1}{2}\right]\left[u_{1}-u_{2}+1\right]} S_{N}\left(z_{2}\right) S_{N}\left(z_{1}\right) \\
& S_{i}\left(z_{1}\right) S_{j}\left(z_{2}\right)=\frac{\left[u_{1}-u_{2}+\frac{1}{2}\right]}{\left[u_{2}-u_{1}+\frac{1}{2}\right]} S_{j}\left(z_{2}\right) S_{i}\left(z_{1}\right), \quad|i-j|=1, \quad 1 \leq i, j \leq N \\
& S_{i}\left(z_{1}\right) S_{j}\left(z_{2}\right)=S_{j}\left(z_{2}\right) S_{i}\left(z_{1}\right), \quad|i-j| \geq 2, \quad 1 \leq i, j \leq N \tag{B.8}
\end{align*}
$$

Here we set $z_{i}=x^{2 u_{i}}, i=1,2$ and $[u]=x^{\frac{u^{2}}{r}-2 u} \Theta_{x^{2 r}}(z)$.

## B. $3 \quad \Lambda_{i}(z)$ and $S_{i}(z)$

From (A.5), (A.6) and (A.8) we obtain

$$
\begin{align*}
& {\left[\Lambda_{k}\left(z_{1}\right), S_{k}\left(z_{2}\right)\right]=\left(x^{-2 r+2}-1\right): \Lambda_{k}\left(z_{1}\right) S_{k}\left(z_{2}\right): \delta\left(\frac{x^{k-r} z_{2}}{z_{1}}\right), \quad 1 \leq k \leq N} \\
& {\left[\Lambda_{k+1}\left(z_{1}\right), S_{k}\left(z_{2}\right)\right]=\left(x^{2 r-2}-1\right): \Lambda_{k+1}\left(z_{1}\right) S_{k}\left(z_{2}\right): \delta\left(\frac{x^{k+r} z_{2}}{z_{1}}\right), \quad 1 \leq k \leq N-1} \\
& {\left[\Lambda_{\bar{k}}\left(z_{1}\right), S_{k}\left(z_{2}\right)\right]=\left(x^{2 r-2}-1\right): \Lambda_{\bar{k}}\left(z_{1}\right) S_{k}\left(z_{2}\right): \delta\left(\frac{x^{2 N+1-k+r} z_{2}}{z_{1}}\right), \quad 1 \leq k \leq N} \\
& {\left[\Lambda_{\overline{k+1}}\left(z_{1}\right), S_{k}\left(z_{2}\right)\right]=\left(x^{-2 r+2}-1\right): \Lambda_{\frac{1}{k+1}}\left(z_{1}\right) S_{k}\left(z_{2}\right): \delta\left(\frac{x^{2 N+1-k-r} z_{2}}{z_{1}}\right), \quad 1 \leq k \leq N-1} \\
& {\left[\Lambda_{0}\left(z_{1}\right), S_{N}\left(z_{2}\right)\right]=\left(x-x^{-1}\right) \frac{[r-1]_{x}\left[\frac{1}{2}\right]_{x}}{\left[r-\frac{1}{2}\right]_{x}}\left(\delta\left(\frac{x^{r+N} z_{2}}{z_{1}}\right)-\delta\left(\frac{x^{-r+N+1} z_{2}}{z_{1}}\right)\right)} \\
& \quad \times: \Lambda_{0}\left(z_{1}\right) S_{N}\left(z_{2}\right): \tag{B.9}
\end{align*}
$$

Other commutators on the type $\left[\Lambda_{i}\left(z_{1}\right), S_{k}\left(z_{2}\right)\right]$ that are used in the proof of Lemma 4.9 are zeroes.

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## References

[1] Awata H., Kubo H., Odake S., Shiraishi J., Quantum deformation of the $\mathcal{W}_{N}$ algebras, arXiv:q-alg/9612001.
[2] Awata H., Kubo H., Odake S., Shiraishi J., Quantum $\mathcal{W}_{N}$ algebras and Macdonald polynomials, Comm. Math. Phys. 179 (1996), 401-416, arXiv:q-alg/9508011.
[3] Brazhnikov V., Lukyanov S., Angular quantization and form factors in massive integrable models, Nuclear Phys. B 512 (1998), 616-636, arXiv:hep-th/9707091.
[4] Ding J., Feigin B., Quantized $W$-algebra of $\mathfrak{s l}(2,1)$ : a construction from the quantization of screening operators, in Recent Developments in Quantum Affine Algebras and Related Topics (Raleigh, NC, 1998), Contemp. Math., Vol. 248, Amer. Math. Soc., Providence, RI, 1999, 83-108, arXiv:math.QA/9801084.
[5] Feigin B., Frenkel E., Quantum $\mathcal{W}$-algebras and elliptic algebras, Comm. Math. Phys. 178 (1996), 653-678, arXiv:q-alg/9508009.
[6] Feigin B., Jimbo M., Mukhin E., Vilkoviskiy I., Deformations of $\mathcal{W}$ algebras via quantum toroidal algebras, Selecta Math. (N.S.) 27 (2021), 52, 62 pages, arXiv:2003.04234.
[7] Frenkel E., Reshetikhin N., Quantum affine algebras and deformations of the Virasoro and $\mathcal{W}$-algebras, Comm. Math. Phys. 178 (1996), 237-264, arXiv:q-alg/9505025.
[8] Frenkel E., Reshetikhin N., Deformations of $\mathcal{W}$-algebras associated to simple Lie algebras, Comm. Math. Phys. 197 (1998), 1-32, arXiv:q-alg/9708006.
[9] Frenkel E., Reshetikhin N., Semenov-Tian-Shansky M.A., Drinfeld-Sokolov reduction for difference operators and deformations of $\mathcal{W}$-algebras. I. The case of Virasoro algebra, Comm. Math. Phys. 192 (1998), 605-629, arXiv:q-alg/9704011.
[10] Harada K., Matsuo Y., Noshita G., Watanabe A., $q$-Deformation of corner vertex operator algebras by Miura transformation, J. High Energy Phys. 2021 (2021), no. 4, 202, 49 pages, arXiv:2101.03953.
[11] Kac V.G., Infinite-dimensional Lie algebras, 3rd ed., Cambridge University Press, Cambridge, 1990.
[12] Kojima T., Quadratic relations of the deformed $W$-superalgebra $\mathcal{W}_{q, t}(A(M, N))$, J. Phys. A 54 (2021), 335201, 37 pages, arXiv:2101.01110.
[13] Kojima T., Quadratic relations of the deformed $W$-superalgebra $\mathcal{W}_{q, t}(\mathfrak{s l}(2 \mid 1))$, J. Math. Phys. 62 (2021), 051702, 19 pages, arXiv:1912.03096.
[14] Odake S., Comments on the deformed $W_{N}$ algebra, Internat. J. Modern Phys. B 16 (2002), 2055-2064, arXiv:math.QA/0111230.
[15] Semenov-Tian-Shansky M.A., Sevostyanov A.V., Drinfeld-Sokolov reduction for difference operators and deformations of $\mathcal{W}$-algebras. II. The general semisimple case, Comm. Math. Phys. 192 (1998), 631-647, arXiv:q-alg/9702016.
[16] Sevostyanov A., Drinfeld-Sokolov reduction for quantum groups and deformations of $W$-algebras, Selecta Math. (N.S.) 8 (2002), 637-703, arXiv:math.QA/0107215.
[17] Shiraishi J., Kubo H., Awata H., Odake S., A quantum deformation of the Virasoro algebra and the Macdonald symmetric functions, Lett. Math. Phys. 38 (1996), 33-51, arXiv:q-alg/9507034.
[18] van de Leur J.W., Contragredient Lie superalgebras of finite growth, Ph.D. Thesis, Utrecht University, 1986.
[19] van de Leur J.W., A classification of contragredient Lie superalgebras of finite growth, Comm. Algebra $\mathbf{1 7}$ (1989), 1815-1841.


[^0]:    ${ }^{1}$ We use two types of symbols, $\mathcal{W}_{x, r}(\mathfrak{g})$ and $\mathcal{W}_{x, r}\left(X_{n}^{(r)}\right)$, for the deformed $W$-algebra associated with the affine Lie algebra $\mathfrak{g}$ of type $X_{n}^{(r)}$.

[^1]:    ${ }^{2}$ We define $Y_{i}(z)^{-1}$ as the inverse element of $Y_{i}(z)$, that is, $Y_{i}(z) Y_{i}(z)^{-1}=Y_{i}(z)^{-1} Y_{i}(z)=1$. Specifically, we obtain $Y_{i}(z)^{-1}=x^{-r y_{i}(0)}\left\langle Y_{i}(z) Y_{i}(z)\right\rangle: \exp \left(-\sum_{m \neq 0} y_{i}(m) z^{-m}\right)$ :, where we used the symbol $\rangle$ defined in (A.1).

