# A Characterization of Invariant Connections 

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#### Abstract

Given a principal fibre bundle with structure group $S$ and a fibre transitive Lie group $G$ of automorphisms thereon, Wang's theorem identifies the invariant connections with certain linear maps $\psi: \mathfrak{g} \rightarrow \mathfrak{s}$. In the present paper we prove an extension of this theorem that applies to the general situation where $G$ acts non-transitively on the base manifold. We consider several special cases of the general theorem including the result of Harnad, Shnider and Vinet which applies to the situation where $G$ admits only one orbit type. Along the way we give applications to loop quantum gravity.


Key words: invariant connections; principal fibre bundles; loop quantum gravity; symmetry reduction
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## 1 Introduction

The set of connections on a principal fibre bundle $(P, \pi, M, S)$ is closed under pullback by automorphisms and it is natural to search for connections that do not change under this operation. Especially, connections invariant under a Lie group $(G, \Phi)$ of automorphisms are of particular interest as they reflect the symmetry of the whole group and, for this reason, find their applications in the symmetry reduction of (quantum) gauge field theories [1, 4, 5]. The first classification theorem for such connections was given by Wang [8], cf. Case 5.7. This applies to the case where the induced action ${ }^{1} \varphi$ acts transitively on the base manifold and states that each point in the bundle gives rise to a bijection between the set of $\Phi$-invariant connections and certain linear maps $\psi: \mathfrak{g} \rightarrow \mathfrak{s}$. In [6] the authors generalize this to the situation where $\varphi$ admits only one orbit type. More precisely, they discuss a variation ${ }^{2}$ of the case where the bundle admits a submanifold $P_{0}$ with $\pi\left(P_{0}\right)$ intersecting each $\varphi$-orbit in a unique point, see Case 4.5 and Example 4.6. Here the $\Phi$-invariant connections are in bijection with such smooth maps $\psi: \mathfrak{g} \times P_{0} \rightarrow \mathfrak{s}$ for which the restrictions $\left.\psi\right|_{\mathfrak{g} \times T_{p_{0}} P_{0}}$ are linear for all $p_{0} \in P_{0}$ and that fulfil additional consistency conditions.

Now, in the general case we consider $\Phi$-coverings of $P$. These are families $\left\{P_{\alpha}\right\}_{\alpha \in I}$ of immersed submanifolds ${ }^{3} P_{\alpha}$ of $P$ such that each $\varphi$-orbit has non-empty intersection with $\bigcup_{\alpha \in I} \pi\left(P_{\alpha}\right)$ and for which $T_{p} P=T_{p} P_{\alpha}+\mathrm{d}_{e} \Phi_{p}(\mathfrak{g})+T v_{p} P$ holds whenever $p \in P_{\alpha}$ for some $\alpha \in I$. Here $T v_{p} P \subseteq T_{p} P$ denotes the vertical tangent space at $p \in P$ and $e$ the identity of $G$. Observe that the intersection properties of the sets $\pi\left(P_{\alpha}\right)$ with the $\varphi$ - orbits in the base manifold need not to be convenient in any sense. Here one might think of situations in which $\varphi$ admits dense orbits, or of the almost-fibre transitive case, cf. Case 5.4.

[^0]If $\Theta:(G \times S) \times P \rightarrow P$ is defined by $((g, s), p) \mapsto \Phi(g, p) \cdot s^{-1}$, then the main result of the present paper can be stated as follows:

Theorem. Let $(P, \pi, M, S)$ be a principal fibre bundle and $(G, \Phi)$ a Lie group of automorphisms thereon. Then each $\Phi$-covering $\left\{P_{\alpha}\right\}_{\alpha \in I}$ admits a bijection between the $\Phi$-invariant connections on $P$ and the families $\left\{\psi_{\alpha}\right\}_{\alpha \in I}$ of smooth maps $\psi_{\alpha}: \mathfrak{g} \times T P_{\alpha} \rightarrow \mathfrak{s}$ for which $\left.\psi_{\alpha}\right|_{\mathfrak{g} \times T_{p_{\alpha}} P_{\alpha}}$ is linear for all $p_{\alpha} \in P_{\alpha}$ and that fulfil the following two (generalized Wang) conditions:

$$
\begin{aligned}
& \text { - } \widetilde{g}\left(p_{\beta}\right)+\vec{w}_{p_{\beta}}-\widetilde{s}\left(p_{\beta}\right)=\mathrm{d} L_{q} \vec{w}_{p_{\alpha}} \quad \Longrightarrow \quad \psi_{\beta}\left(\vec{g}, \vec{w}_{p_{\beta}}\right)-\vec{s}=\rho(q) \circ \psi_{\alpha}\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{w}_{p_{\alpha}}\right), \\
& \text { - } \psi_{\beta}\left(\operatorname{Ad}_{q}(\vec{g}), \overrightarrow{0}_{p_{\beta}}\right)=\rho(q) \circ \psi_{\alpha}\left(\vec{g}, \overrightarrow{0}_{p_{\alpha}}\right)
\end{aligned}
$$

Here $q \in G \times S, p_{\alpha} \in P_{\alpha}, p_{\beta} \in P_{\beta}$ with $p_{\beta}=q \cdot p_{\alpha}$ and $\vec{w}_{p_{\alpha}} \in T_{p_{\alpha}} P_{\alpha}, \vec{w}_{p_{\beta}} \in T_{p_{\beta}} P_{\beta}$. Moreover, $\widetilde{g}, \widetilde{s}$ denote the fundamental vector fields assigned to the elements $\vec{g} \in \mathfrak{g}$ and $\vec{s} \in \mathfrak{s}$, respectively.

Using this theorem the calculation of invariant connections reduces to identifying a $\Phi$-covering that makes the above conditions as easy as possible. Here one has to find the balance between quantity and complexity of these conditions. Of course, the more submanifolds there are, the more conditions we have, so that usually it is convenient to use as few of them as possible. For instance, in the situation where $\varphi$ is transitive it suggests itself to choose a $\Phi$-covering that consists of one single point which, in turn, has to be chosen appropriately. Also if there is some $m \in M$ contained in the closure of each $\varphi$-orbit, one single submanifold is sufficient, see Case 5.4 and Example 5.5. The same example shows that sometimes pointwise ${ }^{4}$ evaluation of the above conditions proves non-existence of $\Phi$-invariant connections.

In any case, one can use the inverse function theorem to construct a $\Phi$-covering $\left\{P_{\alpha}\right\}_{\alpha \in I}$ of $P$ such that the submanifolds $P_{\alpha}$ have minimal dimension in a certain sense, see see Lemma 3.4 and Corollary 5.1. This reproduces the description of connections by means of local 1-forms on $M$ provided that $G$ acts trivially or, more generally, via gauge transformations on $P$, see Case 5.2.

Finally, since orbit structures can depend very sensitively on the action or the group, one cannot expect to have a general concept for finding the $\Phi$-covering optimal for calculations. Indeed, sometimes these calculations become easier if one uses coverings that seem less optimal at a first sight ${ }^{5}$.

The present paper is organized as follows: In Section 2 we fix the notations. In Section 3 we introduce the notion of a $\Phi$-covering, the central object of this paper. In Section 4 we prove the main theorem and deduce a slightly more general version of the result from [6]. In Section 5 we show how to construct $\Phi$-coverings to be used in special situations. In particular, we consider the (almost-)fibre transitive case, trivial principal fibre bundles and Lie groups of gauge transformations. Along the way we give applications to loop quantum gravity.

## 2 Preliminaries

We start with fixing the notations.

### 2.1 Notations

Manifolds are always assumed to be smooth. If $M, N$ are manifolds and $f: M \rightarrow N$ is a smooth map, then $\mathrm{d} f: T M \rightarrow T N$ denotes the differential map between their tangent manifolds. The map $f$ is said to be an immersion iff for each $x \in M$ the restriction $\mathrm{d}_{x} f:=\left.\mathrm{d} f\right|_{T_{x} M}: T_{x} M \rightarrow$ $T_{f(x)} N$ is injective.

[^1]Let $V$ be a finite dimensional vector space. A $V$-valued 1-form $\omega$ on the manifold $N$ is a smooth map $\omega: T N \rightarrow V$ whose restriction $\omega_{y}:=\left.\omega\right|_{T_{y} N}$ is linear for all $y \in N$. The pullback of $\omega$ by $f$ is the $V$-valued 1-form $f^{*} \omega: T M \rightarrow V, \vec{v}_{x} \rightarrow \omega_{f(x)}\left(\mathrm{d}_{x} f\left(\vec{v}_{x}\right)\right)$.

Let $G$ be a Lie group and $\mathfrak{g}$ its Lie algebra. For $g \in G$ we define the corresponding conjugation map by $\alpha_{g}: G \rightarrow G, h \mapsto g h g^{-1}$. Its differential $\mathrm{d}_{e} \alpha_{g}: \mathfrak{g} \rightarrow \mathfrak{g}$ at the unit element $e \in G$ is denoted by $\operatorname{Ad}_{g}$ in the following.

Let $\Psi$ be a (left) action of the Lie group $G$ on the manifold $M$. If $g \in G$, then $\Psi_{g}: M \rightarrow M$ denotes the map $\Psi_{g}: x \mapsto \Psi(g, x)$. We often write $L_{g}$ instead of $\Psi_{g}$ as well as $g \cdot x$ or $g x$ instead of $\Psi_{g}(x)$ if it is clear, which action is meant. If $x \in M$, let $\Psi_{x}: G \rightarrow M, g \mapsto \Psi(g, x)$. Then for $\vec{g} \in \mathfrak{g}$ and $x \in M$ the map $\widetilde{g}(x):=\left.\frac{\mathrm{d}}{\mathrm{d} t}\right|_{t=0} \Psi_{x}(\exp (t \vec{g}))$ is called the fundamental vector field w.r.t. $\vec{g}$. The Lie subgroup $G_{x}:=\{g \in G \mid g \cdot x=x\}$ is called the stabilizer of $x \in M$ (w.r.t. $\Psi$ ) and its Lie algebra $\mathfrak{g}_{x}$ equals $\operatorname{ker}\left[\mathrm{d}_{x} \Psi\right]$, see e.g. [3]. The orbit of $x$ under $G$ is the set $G x:=\operatorname{im}\left[\Psi_{x}\right]$, and $\Psi$ is said to be transitive iff $G x=M$ for one and then each $x \in M$. Analogous conventions also hold for right actions.

### 2.2 Invariant connections

Let $\pi: P \rightarrow M$ be a smooth (surjective) map between manifolds $P$ and $M$, and denote by $F_{x}:=\pi^{-1}(x) \subseteq P$ the fibre over $x \in M$ in $P$. Assume that $(S, R)$ is a Lie group that acts from the right on $P$. If there is an open covering $\left\{U_{\alpha}\right\}_{\alpha \in I}$ of $M$ and a family $\left\{\phi_{\alpha}\right\}_{\alpha \in I}$ of diffeomorphisms $\phi_{\alpha}: \pi^{-1}\left(U_{\alpha}\right) \rightarrow U_{\alpha} \times S$ with $^{6}$

$$
\begin{equation*}
\phi_{\alpha}(p \cdot s)=\left(\pi(p),\left[\operatorname{pr}_{2} \circ \phi_{\alpha}\right](p) \cdot s\right) \quad \forall p \in \pi^{-1}\left(U_{\alpha}\right), \quad \forall s \in S \tag{2.1}
\end{equation*}
$$

then $(P, \pi, M, S)$ is called principal fibre bundle with total space $P$, projection map $\pi$, base manifold $M$ and structure group $S$. It follows from (2.1) that

- $R_{s}\left(F_{x}\right) \subseteq F_{x}$ for all $x \in M$ and all $s \in S$,
- if $x \in M$ and $p, p^{\prime} \in F_{x}$, then $p^{\prime}=p \cdot s$ for a unique element $s \in S$.

The subspace $T v_{p} P:=\operatorname{ker}\left[d_{p} \pi\right] \subseteq T_{p} P$ is called vertical tangent space at $p \in P$ and

$$
\widetilde{s}(p):=\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} p \cdot \exp (t \vec{s}) \in T v_{p} P \quad \forall p \in P,
$$

denotes the fundamental vector field of $\vec{s}$ w.r.t. the right action of $S$ on $P$. The map $\mathfrak{s} \ni \vec{s} \rightarrow$ $\widetilde{s}(p) \in T v_{p} P$ is a vector space isomorphism for all $p \in P$.

Complementary to that, a connection $\omega$ is an $\mathfrak{s}$-valued 1 -form on $P$ with

- $R_{s}^{*} \omega=\operatorname{Ad}_{s^{-1}} \circ \omega$ for all $s \in S$,
- $\omega_{p}(\widetilde{s}(p))=\vec{s}$ for all $\vec{s} \in \mathfrak{s}$.

The subspace $T h_{p} P:=\operatorname{ker}\left[\omega_{p}\right] \subseteq T_{p} P$ is called the horizontal tangent space at $p$ (w.r.t. $\omega$ ). We have $\mathrm{d} R_{s}\left(T h_{p} P\right)=T h_{p \cdot s} P$ for all $s \in S$ and one can show that $T_{p} P=T v_{p} P \oplus T h_{p} P$ for all $p \in P$.

A diffeomorphism $\kappa: P \rightarrow P$ is said to be an automorphism iff $\kappa(p \cdot s)=\kappa(p) \cdot s$ for all $p \in P$ and all $s \in S$. It is straightforward to see that an $\mathfrak{s}$-valued 1 -form $\omega$ on $P$ is a connection iff this is true for the pullback $\kappa^{*} \omega$. A Lie group of automorphisms $(G, \Phi)$ of $P$ is a Lie group $G$ together with a left action $\Phi$ of $G$ on $P$ such that the map $\Phi_{g}$ is an automorphism for each $g \in G$. This is equivalent to say that $\Phi(g, p \cdot s)=\Phi(g, p) \cdot s$ for all $p \in P, g \in G$ and all $s \in S$. In this situation we often write $g p s$ instead of $(g \cdot p) \cdot s=g \cdot(p \cdot s)$. Each such a left action $\Phi$ gives rise to two further actions:

[^2]- The induced action $\varphi$ is defined by

$$
\begin{array}{ll}
\varphi: \quad & G \times M \rightarrow M \\
& (g, m) \mapsto(\pi \circ \Phi)\left(g, p_{m}\right), \tag{2.2}
\end{array}
$$

where $p_{m} \in \pi^{-1}(m)$ is arbitrary. $\Phi$ is called fibre transitive iff $\varphi$ is transitive.

- We equip $Q=G \times S$ with the canonical Lie group structure and define [8]

$$
\begin{array}{ll}
\Theta: \quad & Q \times P \rightarrow P \\
& ((g, s), p) \mapsto \Phi\left(g, p \cdot s^{-1}\right) \tag{2.3}
\end{array}
$$

A connection $\omega$ is called $\Phi$-invariant iff $\Phi_{g}^{*} \omega=\omega$ for all $g \in G$. This is equivalent to require that for each $p \in P$ and $g \in G$ the differential $\mathrm{d}_{p} L_{g}$ induces an isomorphism between the horizontal tangent spaces $T h_{p} P$ and $T h_{g p} P .{ }^{7}$

We conclude this subsection with the following straightforward facts, see also [8]:

- Consider the representation $\rho: Q \rightarrow \operatorname{Aut}(\mathfrak{s}),(g, s) \mapsto \operatorname{Ad}_{s}$. Then it is immediate that each $\Phi$-invariant connection $\omega$ is of type $\rho$, i.e., $\omega$ is an $\mathfrak{s}$-valued 1-form on $P$ with $L_{q}^{*} \omega=\rho(q) \circ \omega$ for all $q \in Q$.
- An $\mathfrak{s}$-valued 1 -form $\omega$ on $P$ with $\omega(\widetilde{s}(p))=\vec{s}$ for all $\vec{s} \in \mathfrak{s}$ is a $\Phi$-invariant connection iff it is of type $\rho$.
- Let $Q_{p}$ denote the stabilizer of $p \in P$ w.r.t. $\Theta$ and $G_{\pi(p)}$ the stabilizer of $\pi(p)$ w.r.t. $\varphi$. Then $G_{\pi(p)}=\left\{g \in G \mid L_{g}: F_{p} \rightarrow F_{p}\right\}$ and we obtain a smooth homomorphism (Lie group homomorphism) $\phi_{p}: G_{\pi(p)} \rightarrow S$ by requiring that $\Phi(j, p)=p \cdot \phi_{p}(j)$ for all $j \in G_{\pi(p)}$. If $\mathfrak{q}_{p}$ and $\mathfrak{g}_{\pi(p)}$ denote the Lie algebras of $Q_{p}$ and $G_{\pi(p)}$, respectively, then

$$
\begin{equation*}
Q_{p}=\left\{\left(j, \phi_{p}(j)\right) \mid j \in G_{\pi(p)}\right\} \quad \text { and } \quad \mathfrak{q}_{p}=\left\{\left(\vec{j}, \mathrm{~d}_{e} \phi_{p}(\vec{j})\right) \mid \vec{j} \in \mathfrak{g}_{\pi(p)}\right\} \tag{2.4}
\end{equation*}
$$

## 3 -coverings

We start this section with some facts and conventions concerning submanifolds. Then we give the definition of a $\Phi$-covering and discuss some its properties.

Convention 3.1. Let $M$ be a manifold.

1. A pair $\left(N, \tau_{N}\right)$ consisting of a manifold $N$ and an injective immersion $\tau: N \rightarrow M$ is called submanifold of $M$.
2. If $\left(N, \tau_{N}\right)$ is a submanifold of $M$, we tacitly identify $N$ and $T N$ with their images $\tau_{N}(N) \subseteq$ $M$ and $\mathrm{d} \tau_{N}(T N) \subseteq T M$, respectively. In particular, this means that:

- If $M^{\prime}$ is a manifold and $\kappa: M \rightarrow M^{\prime}$ is a smooth map, then for $x \in N$ and $\vec{v} \in T N$ we write $\kappa(x)$ and $\mathrm{d} \kappa(\vec{v})$ instead of $\kappa\left(\tau_{N}(x)\right)$ and $\mathrm{d} \kappa(\mathrm{d} \tau(\vec{v}))$, respectively.
- If $\Psi: G \times M \rightarrow M$ is a left action of the Lie group $G$ and $\left(H, \tau_{H}\right)$ a submanifold of $G$, then the restriction of $\Psi$ to $H \times N$ is defined by

$$
\left.\Psi\right|_{H \times N}(h, x):=\Psi\left(\tau_{H}(h), \tau_{N}(x)\right) \quad \forall(h, x) \in H \times N .
$$

- If $\omega: T M \rightarrow V$ is a $V$-valued 1-form on $M$, then

$$
\left.\left(\Psi^{*} \omega\right)\right|_{T G \times T N}(\vec{m}, \vec{v}):=\left(\Psi^{*} \omega\right)(\vec{m}, \mathrm{~d} \tau(\vec{v})) \quad \forall(\vec{m}, \vec{v}) \in T G \times T N .
$$

- We will not explicitly refer to the maps $\tau_{N}$ and $\tau_{H}$ in the following.

[^3]3. Open subsets $U \subseteq M$ are equipped with the canonical manifold structure making the inclusion map an embedding.
4. If $L$ is a submanifold of $N$ and $N$ is a submanifold of $M$, we consider $L$ as a submanifold of $M$ in the canonical way.

Definition 3.2. A submanifold $N \subseteq M$ is called $\Psi$-patch iff for each $x \in N$ there is an open neighbourhood $N^{\prime} \subseteq N$ of $x$ and a submanifold $H$ of $G$ through $e$ such that the restriction $\left.\Psi\right|_{H \times N^{\prime}}$ is a diffeomorphism to an open subset $U \subseteq M$.

## Remark 3.3.

1. It follows from the inverse function theorem and ${ }^{8}$

$$
\mathrm{d}_{(e, x)} \Psi\left(\mathfrak{g} \times T_{x} N\right)=\mathrm{d}_{e} \Psi_{x}(\mathfrak{g})+\mathrm{d}_{x} \Psi_{e}\left(T_{x} N\right)=\mathrm{d}_{e} \Psi_{x}(\mathfrak{g})+T_{x} N \quad \forall x \in N
$$

that $N$ is a $\Psi$-patch iff for each $x \in N$ we have $T_{x} M=\mathrm{d}_{e} \Psi_{x}(\mathfrak{g})+T_{x} N .{ }^{9}$
2. Open subsets $U \subseteq M$ are always $\Psi$-patches. They are of maximal dimension which, for instance, is necessary if there is a point in $U$ whose stabilizer equals $G$, see Lemma 3.4.1.
3. We allow zero-dimensional patches, i.e., $N=\{x\}$ for $x \in M$. Necessarily, then $\mathrm{d}_{e} \Psi_{x}(\mathfrak{g})=$ $T_{x} M$ and $\left.\Psi\right|_{H \times N}=\left.\Psi_{x}\right|_{H}$ for each submanifold $H$ of $G$.

The second part of the next elementary lemma equals Lemma 2.1.1 in [3].
Lemma 3.4. Let $(G, \Psi)$ be a Lie group that acts on the manifold $M$ and let $x \in M$.

1. If $N$ is a $\Psi$-patch with $x \in N$, then $\operatorname{dim}[N] \geq \operatorname{dim}[M]-\operatorname{dim}[G]+\operatorname{dim}\left[G_{x}\right]$.
2. Let $V$ and $W$ be algebraic complements of $\mathrm{d}_{e} \Psi_{x}(\mathfrak{g})$ in $T_{x} M$ and of $\mathfrak{g}_{x}$ in $\mathfrak{g}$, respectively. Then there are submanifolds $N$ of $M$ through $x$ and $H$ of $G$ through e such that $T_{x} N=V$, $T_{e} H=W$. In particular, $N$ is a $\Psi$-patch and $\operatorname{dim}[N]=\operatorname{dim}[M]-\operatorname{dim}[G]+\operatorname{dim}\left[G_{x}\right]$.

Proof. 1. By Remark 3.3.1 and since $\operatorname{ker}\left[\mathrm{d}_{e} \Psi_{x}\right]=\mathfrak{g}_{x}$, we have

$$
\begin{equation*}
\operatorname{dim}[M] \leq \operatorname{dim}\left[\mathrm{d}_{e} \Psi_{x}(\mathfrak{g})\right]+\operatorname{dim}\left[T_{x} N\right]=\operatorname{dim}[G]-\operatorname{dim}\left[G_{x}\right]+\operatorname{dim}[N] \tag{3.1}
\end{equation*}
$$

2. Of course, we find submanifolds $N^{\prime}$ of $M$ through $x$ and $H^{\prime}$ of $G$ through $e$ such that $T_{x} N^{\prime}=V$ and $T_{e} H^{\prime}=W$. So, if $\vec{g} \in \mathfrak{g}$ and $\vec{v}_{x} \in T_{x} N^{\prime}$, then $0=\mathrm{d}_{(e, x)} \Psi\left(\vec{g}, \vec{v}_{x}\right)=\mathrm{d}_{e} \Psi_{x}(\vec{g})+\vec{v}_{x}$ implies $\mathrm{d}_{e} \Psi_{x}(\vec{g})=0$ and $\vec{v}_{x}=0$. Hence, $\vec{g} \in \operatorname{ker}\left[\mathrm{~d}_{e} \Psi_{x}\right]=\mathfrak{g}_{x}$ so that $\left.{ }^{10} \mathrm{~d}_{(e, x)} \Psi\right|_{T_{e} H^{\prime} \times T_{e} N^{\prime}}$ is injective. It is immediate from the definitions that this map is surjective so that by the inverse function theorem we find open neighbourhoods $N \subseteq N^{\prime}$ of $x$ and $H \subseteq G$ of $e$ such that $\left.\Psi\right|_{H \times N}$ is a diffeomorphism to an open subset $U \subseteq M$. Then $N$ is a $\Psi$-patch and since in (3.1) equality holds, also the last claim is clear.

Definition 3.5. Let $(G, \Phi)$ be a Lie group of automorphisms of the principal fibre bundle $P$ and recall the actions $\varphi$ and $\Theta$ defined by (2.2) and (2.3), respectively. A family of $\Theta$ patches $\left\{P_{\alpha}\right\}_{\alpha \in I}$ is said to be a $\Phi$-covering of $P$ iff each $\varphi$-orbit intersects at least one of the sets $\pi\left(P_{\alpha}\right)$.

[^4]
## Remark 3.6.

1. If $O \subseteq P$ is a $\Theta$-patch, then Lemma 3.4.1 and (2.4) yield

$$
\operatorname{dim}[O] \geq \operatorname{dim}[P]-\operatorname{dim}[Q]+\operatorname{dim}\left[Q_{p}\right] \stackrel{(2.4)}{=} \operatorname{dim}[M]-\operatorname{dim}[G]+\operatorname{dim}\left[G_{\pi(p)}\right]
$$

2. It follows from Remark 3.3.1 and $\mathrm{d}_{e} \Theta_{p}(\mathfrak{q})=\mathrm{d}_{e} \Phi_{p}(\mathfrak{g})+T v_{p} P$ that $O$ is a $\Theta$-patch iff

$$
\begin{equation*}
T_{p} P=T_{p} O+\mathrm{d}_{e} \Phi_{p}(\mathfrak{g})+T v_{p} P \quad \forall p \in O . \tag{3.2}
\end{equation*}
$$

As a consequence

- each $\Phi$-patch is a $\Theta$-patch,
- $P$ is always a $\Phi$-covering by itself and if $P=M \times S$ is trivial, then $M \times\{e\}$ is a $\Phi$-covering.

3. If $N$ is a $\varphi$-patch and $s_{0}: N \rightarrow P$ a smooth section, i.e., $\pi \circ s_{0}=\operatorname{id}_{N}$, then $O:=s_{0}(N)$ is a $\Theta$-patch as Lemma 3.7.2 shows.

Lemma 3.7. Let $(G, \Phi)$ be a Lie group of automorphisms of the principal bundle $(P, \pi, M, S)$.

1. If $O \subseteq P$ is a $\Theta$-patch, then for each $p \in O$ and $q \in Q$ the differential $\mathrm{d}_{(q, p)} \Theta: T_{q} Q \times T_{p} O \rightarrow$ $T_{q \cdot p} P$ is surjective.
2. If $N$ is a $\varphi$-patch and $s_{0}: N \rightarrow P$ a smooth section, then $O:=s_{0}(N)$ is a $\Theta$-patch.

Proof. 1. Since $O$ is a $\Theta$-patch, the claim is clear for $q=e$. If $q$ is arbitrary, then for each $\vec{m}_{q} \in T_{q} Q$ we find some $\vec{q} \in \mathfrak{q}$ such that $\vec{m}_{q}=\mathrm{d} L_{q} \vec{q}$. Consequently, for $\vec{w}_{p} \in T_{p} P$ we have

$$
\mathrm{d}_{(q, p)} \Theta\left(\vec{m}_{q}, \vec{w}_{p}\right)=\mathrm{d}_{(q, p)} \Theta\left(\mathrm{d} L_{q} \vec{q}, \vec{w}_{p}\right)=\mathrm{d}_{p} L_{q}\left(\mathrm{~d}_{(e, p)} \Theta\left(\vec{q}, \vec{w}_{p}\right)\right) .
$$

But, left translation w.r.t. $\Theta$ is a diffeomorphism so that $\mathrm{d}_{p} L_{q}$ is surjective.
2. First observe that $O$ is a submanifold of $P$ because $s_{0}$ is an injective immersion. By Remark 3.6.2 it suffices to show that

$$
\operatorname{dim}\left[T_{s_{0}(x)} O+\mathrm{d}_{e} \Phi_{s_{0}(x)}(\mathfrak{g})+T v_{s_{0}(x)} P\right] \geq \operatorname{dim}\left[T_{s_{0}(x)} P\right] \quad \forall x \in N
$$

For this, let $x \in N$ and $V^{\prime} \subseteq \mathfrak{g}$ be a linear subspace such that $T_{x} M=T_{x} N \oplus \mathrm{~d}_{e} \varphi_{x}\left(V^{\prime}\right)$. Then $T_{s_{0}(x)} O \oplus \mathrm{~d}_{e} \Phi_{s_{0}(x)}\left(V^{\prime}\right) \oplus T v_{s_{0}(x)} P$ because if $\mathrm{d}_{x} s_{0}\left(\vec{v}_{x}\right)+\mathrm{d}_{e} \Phi_{s_{0}(x)}\left(\vec{g}^{\prime}\right)+\vec{v}_{v}=0$ for $\vec{v}_{x} \in T_{x} N$, $\vec{g}^{\prime} \in V^{\prime}$ and $\vec{v}_{v} \in T v_{s_{0}(x)} P$, then

$$
0=\mathrm{d}_{s_{0}(x)} \pi\left(\mathrm{d}_{x} s_{0}\left(\vec{v}_{x}\right)+\mathrm{d}_{e} \Phi_{s_{0}(x)}\left(\vec{g}^{\prime}\right)+\vec{v}_{v}\right)=\vec{v}_{x}+\mathrm{d}_{e} \varphi_{x}\left(\vec{g}^{\prime}\right)
$$

so that $\vec{v}_{x}=0$ and $\vec{g}^{\prime}=0$ by assumption, hence $\vec{v}_{v}=0$. In particular, this shows $\operatorname{dim}\left[\mathrm{d}_{e} \Phi_{s_{0}(x)}\left(V^{\prime}\right)\right]$ $\geq \operatorname{dim}\left[\mathrm{d}_{e} \varphi_{x}\left(V^{\prime}\right)\right]$ and we obtain

$$
\begin{aligned}
\operatorname{dim} & {\left[T_{s_{0}(x)} O+\mathrm{d}_{e} \Phi_{s_{0}(x)}(\mathfrak{g})+T v_{s_{0}(x)} P\right] } \\
& \geq \operatorname{dim}\left[T_{s_{0}(x)} O \oplus \mathrm{~d}_{e} \Phi_{s_{0}(x)}\left(V^{\prime}\right) \oplus T v_{s_{0}(x)} P\right] \\
& =\operatorname{dim}\left[T_{x} N\right]+\operatorname{dim}\left[\mathrm{d}_{e} \Phi_{s_{0}(x)}\left(V^{\prime}\right)\right]+\operatorname{dim}[S] \\
& \geq \operatorname{dim}\left[T_{x} N\right]+\operatorname{dim}\left[\mathrm{d}_{e} \varphi_{x}\left(V^{\prime}\right)\right]+\operatorname{dim}[S] \\
& =\operatorname{dim}[P] .
\end{aligned}
$$

## 4 Characterization of invariant connections

In this section we use $\Phi$-coverings $\left\{P_{\alpha}\right\}_{\alpha \in I}$ of the bundle $P$ to characterize the set of $\Phi$-invariant connections by families $\left\{\psi_{\alpha}\right\}_{\alpha \in I}$ of smooth maps $\psi_{\alpha}: \mathfrak{g} \times T P_{\alpha} \rightarrow \mathfrak{s}$ whose restrictions $\left.\psi_{\alpha}\right|_{\mathfrak{g} \times T_{p_{\alpha}} P_{\alpha}}$ are linear and that fulfil two additional compatibility conditions. Here we follow the lines of Wang's original approach, which means that we generalize the proofs from [8] to the nontransitive case. We will proceed in two steps where the first one is done in the next subsection. Here we show that a $\Phi$-invariant connection gives rise to a consistent family $\left\{\psi_{\alpha}\right\}_{\alpha \in I}$ of smooth maps as described above. We also discuss the situation in [6] in order to make the two conditions more intuitive. Then, in Subsection 4.2, we verify that such families $\left\{\psi_{\alpha}\right\}_{\alpha \in I}$ glue together to a $\Phi$-invariant connection on $P$.

### 4.1 Reduction of invariant connections

In the following let $\left\{P_{\alpha}\right\}_{\alpha \in I}$ be a fixed $\Phi$-covering of $P$ and $\omega$ a $\Phi$-invariant connection on $P$. We define $\omega_{\alpha}:=\left.\left(\Theta^{*} \omega\right)\right|_{T Q \times T P_{\alpha}}$ and $\psi_{\alpha}:=\left.\omega_{\alpha}\right|_{\mathfrak{g} \times T P_{\alpha}}$. For $q^{\prime} \in Q$ we let $\alpha_{q^{\prime}}: Q \times P \rightarrow Q \times P$ denote the map $\alpha_{q^{\prime}}(q, p):=\left(\alpha_{q^{\prime}}(q), p\right)$ for $\alpha_{q^{\prime}}: Q \mapsto Q$ the conjugation map w.r.t. $q^{\prime}$ as defined in Section 2.1.

Lemma 4.1. Let $q \in Q$, $p_{\alpha} \in P_{\alpha}, p_{\beta} \in P_{\beta}$ with $^{11} p_{\beta}=q \cdot p_{\alpha}$ and $\vec{w}_{p_{\alpha}} \in T_{p_{\alpha}} P_{\alpha}$. Then

1) $\omega_{\beta}(\vec{\eta})=\rho(q) \circ \omega_{\alpha}\left(\overrightarrow{0}_{\mathfrak{q}}, \vec{w}_{p_{\alpha}}\right)$ for all $\vec{\eta} \in T Q \times T P_{\beta}$ with $\mathrm{d} \Theta(\vec{\eta})=\mathrm{d} L_{q} \vec{w}_{p_{\alpha}}$,
2) $\left(\alpha_{q}^{*} \omega_{\beta}\right)\left(\vec{m}, \overrightarrow{0}_{p_{\beta}}\right)=\rho(q) \circ \omega_{\alpha}\left(\vec{m}, \overrightarrow{0}_{p_{\alpha}}\right)$ for all $\vec{m} \in T Q$.

Proof. 1. Let $\vec{\eta} \in T_{q^{\prime}} Q \times T_{p} P_{\beta}$ for $q^{\prime} \in Q$. Then, since ${ }^{12} L_{q}^{*} \omega=\rho(q) \circ \omega$ for each $q \in Q$ and $q^{\prime} \cdot p=q \cdot p_{\alpha}=p_{\beta}$, we have

$$
\begin{aligned}
\omega_{\beta}(\vec{\eta}) & =\omega_{q^{\prime} \cdot p}\left(\mathrm{~d}_{\left(q^{\prime}, p\right)} \Theta(\vec{\eta})\right)=\omega_{p_{\beta}}\left(\mathrm{d} L_{q} \vec{w}_{p_{\alpha}}\right)=\left(L_{q}^{*} \omega\right)_{p_{\alpha}}\left(\vec{w}_{p_{\alpha}}\right) \\
& =\rho(q) \circ \omega_{p_{\alpha}}\left(\vec{w}_{p_{\alpha}}\right)=\rho(q) \circ \omega_{p_{\alpha}}\left(\mathrm{d}_{\left(e, p_{\alpha}\right)} \Theta\left(\vec{w}_{p_{\alpha}}\right)\right)=\rho(q) \circ \omega_{\alpha}\left(\overrightarrow{0}_{q}, \vec{w}_{p_{\alpha}}\right) .
\end{aligned}
$$

2. For $\vec{m}_{q^{\prime}} \in T_{q^{\prime}} Q$ let $\gamma:(-\epsilon, \epsilon) \rightarrow Q$ be smooth with $\dot{\gamma}(0)=\vec{m}_{q^{\prime}}$. Then

$$
\begin{aligned}
\left(\alpha_{q}^{*} \omega_{\beta}\right)_{\left(q^{\prime}, p_{\beta}\right)}\left(\vec{m}_{q^{\prime}}, \overrightarrow{0}_{p_{\beta}}\right) & =\omega_{\beta\left(\alpha_{q}\left(q^{\prime}\right), p_{\beta}\right)}\left(\operatorname{Ad}_{q}\left(\vec{m}_{q^{\prime}}\right), \overrightarrow{0}_{p_{\beta}}\right)=\omega_{q q^{\prime} q^{-1} q \cdot p_{\alpha}}\left(\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} q \gamma(t) q^{-1} q \cdot p_{\alpha}\right) \\
& =\left(L_{q}^{*} \omega\right)_{q^{\prime} \cdot p_{\alpha}}\left(\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} \gamma(t) \cdot p_{\alpha}\right)=\rho(q) \circ \omega_{q^{\prime} \cdot p_{\alpha}}\left(\mathrm{d}_{\left(q^{\prime}, p_{\alpha}\right)} \Theta\left(\vec{m}_{q^{\prime}}\right)\right) \\
& =\rho(q) \circ \omega_{\alpha\left(q^{\prime}, p_{\alpha}\right)}\left(\vec{m}_{q^{\prime}}, \overrightarrow{0}_{p_{\alpha}}\right) .
\end{aligned}
$$

Corollary 4.2. Let $q \in Q, p_{\alpha} \in P_{\alpha}, p_{\beta} \in P_{\beta}$ with $p_{\beta}=q \cdot p_{\alpha}$ and $\vec{w}_{p_{\alpha}} \in T_{p_{\alpha}} P_{\alpha}$. Then for $\vec{w}_{p_{\beta}} \in T_{p_{\beta}} P_{\beta}, \vec{g} \in \mathfrak{g}$ and $\vec{s} \in \mathfrak{s}$ we have
i) $\widetilde{g}\left(p_{\beta}\right)+\vec{w}_{p_{\beta}}-\widetilde{s}\left(p_{\beta}\right)=\mathrm{d} L_{q}{\overrightarrow{w_{p_{\alpha}}}}^{\Longrightarrow} \quad \psi_{\beta}\left(\vec{g}, \vec{w}_{p_{\beta}}\right)-\vec{s}=\rho(q) \circ \psi_{\alpha}\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{w}_{p_{\alpha}}\right)$,
ii) $\psi_{\beta}\left(\operatorname{Ad}_{q}(\vec{g}), \overrightarrow{0}_{p_{\beta}}\right)=\rho(q) \circ \psi_{\alpha}\left(\vec{g}, \overrightarrow{0}_{p_{\alpha}}\right)$.

Proof. i) In general, for $\vec{w}_{p} \in T_{p} P, \vec{g} \in \mathfrak{g}$ and $\vec{s} \in \mathfrak{s}$ we have

$$
\begin{equation*}
\mathrm{d}_{(e, p)} \Theta\left((\vec{g}, \vec{s}), \vec{w}_{p}\right)=\mathrm{d}_{(e, p)} \Phi\left(\vec{g}, \vec{w}_{p}\right)-\widetilde{s}(p)=\widetilde{g}(p)+\vec{w}_{p}-\widetilde{s}(p) \tag{4.1}
\end{equation*}
$$

[^5]and, since $\omega$ is a connection, for $\left((\vec{g}, \vec{s}), \vec{w}_{p_{\alpha}}\right) \in \mathfrak{q} \times T P_{\alpha}$ we obtain
\[

$$
\begin{align*}
\omega_{\alpha}\left((\vec{g}, \vec{s}), \vec{w}_{p_{\alpha}}\right) & =\omega\left(\mathrm{d}_{\left(e, p_{\alpha}\right)} \Phi\left(\vec{g}, \vec{w}_{p_{\alpha}}\right)-\widetilde{s}\left(p_{\alpha}\right)\right)=\omega\left(\mathrm{d}_{\left(e, p_{\alpha}\right)} \Phi\left(\vec{g}, \vec{w}_{p_{\alpha}}\right)\right)-\vec{s} \\
& =\omega_{\alpha}\left(\vec{g}, \vec{w}_{p_{\alpha}}\right)-\vec{s}=\psi_{\alpha}\left(\vec{g}, \vec{w}_{p_{\alpha}}\right)-\vec{s} . \tag{4.2}
\end{align*}
$$
\]

Now, assume that $\mathrm{d}_{e} \Phi_{p_{\beta}}(\vec{g})+\vec{w}_{p_{\beta}}-\widetilde{s}(p)=\mathrm{d} L_{q} \vec{w}_{p_{\alpha}}$. Then $\mathrm{d}_{\left(e, p_{\beta}\right)} \Theta\left((\vec{g}, \vec{s}), \vec{w}_{p_{\beta}}\right)=\mathrm{d} L_{q} \vec{w}_{p_{\alpha}}$ by (4.1) so that $\omega_{\beta}\left((\vec{g}, \vec{s}), \vec{w}_{p_{\beta}}\right)=\rho(q) \circ \omega_{\alpha}\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{w}_{p_{\alpha}}\right)$ by Lemma 4.1.1. Consequently,

$$
\psi_{\beta}\left(\vec{g}, \vec{w}_{p_{\beta}}\right)-\vec{s} \stackrel{(4.2)}{=} \omega_{\beta}\left((\vec{g}, \vec{s}), \vec{w}_{p_{\beta}}\right)=\rho(q) \circ \omega_{\alpha}\left(\overrightarrow{0}_{\mathfrak{q}}, \vec{w}_{p_{\alpha}}\right) \stackrel{(4.2)}{=} \rho(q) \circ \psi_{\alpha}\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{w}_{p_{\alpha}}\right) .
$$

ii) Lemma 4.1.2 yields

$$
\psi_{\beta}\left(\operatorname{Ad}_{q}(\vec{g}), \overrightarrow{0}_{p_{\beta}}\right)=\left(\alpha_{q}^{*} \omega_{\beta}\right)_{\left(e, p_{\beta}\right)}\left(\vec{g}, \overrightarrow{0}_{p_{\beta}}\right)=\rho(q) \circ\left(\omega_{\alpha}\right)_{\left(e, p_{\alpha}\right)}\left(\vec{g}, \overrightarrow{0}_{p_{\alpha}}\right)=\rho(q) \circ \psi_{\alpha}\left(\vec{g}, \overrightarrow{0}_{p_{\alpha}}\right)
$$

Definition 4.3. A family $\left\{\psi_{\alpha}\right\}_{\alpha \in I}$ of smooth maps $\psi_{\alpha}: \mathfrak{g} \times T P_{\alpha} \rightarrow \mathfrak{s}$ that are linear in the sense that $\left.\psi_{\alpha}\right|_{\mathfrak{g} \times T_{p_{\alpha}} P_{\alpha}}$ is linear for all $p_{\alpha} \in P_{\alpha}$ is called reduced connection w.r.t. $\left\{P_{\alpha}\right\}_{\alpha \in I}$ iff it fulfils the conditions $i$ ) and $i i$ ) from Corollary 4.2.

## Remark 4.4.

1) In particular, Corollary 4.2.i) encodes the following condition
a) For all $\beta \in I,(\vec{g}, \vec{s}) \in \mathfrak{q}$ and $\vec{w}_{p_{\beta}} \in T_{p_{\beta}} P_{\beta}$ we have

$$
\widetilde{g}\left(p_{\beta}\right)+\vec{w}_{p_{\beta}}-\widetilde{s}\left(p_{\beta}\right)=0 \quad \Longrightarrow \quad \psi_{\beta}\left(\vec{g}, \vec{w}_{p_{\beta}}\right)-\vec{s}=0 .
$$

2) Assume that $a$ ) is true and let $q \in Q, p_{\alpha} \in P_{\alpha}, p_{\beta} \in P_{\beta}$ with $p_{\beta}=q \cdot p_{\alpha}$. Moreover, assume that we find elements $\vec{w}_{p_{\alpha}} \in T_{p_{\alpha}} P_{\alpha}$ and $\left((\vec{g}, \vec{s}), \vec{w}_{p_{\beta}}\right) \in \mathfrak{q} \times T_{p_{\beta}} P_{\beta}$ such that

$$
\mathrm{d}_{\left(e, p_{\beta}\right)} \Theta\left((\vec{g}, \vec{s}), \vec{w}_{p_{\beta}}\right)=\mathrm{d} L_{q} \vec{w}_{p_{\alpha}} \quad \text { and } \quad \psi_{\beta}\left(\vec{g}, \vec{w}_{p_{\beta}}\right)-\vec{s}=\rho(q) \circ \psi_{\alpha}\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{w}_{p_{\alpha}}\right)
$$

holds. Then $\psi_{\beta}\left(\vec{g}^{\prime}, \vec{w}_{p_{\beta}}^{\prime}\right)-\vec{s}^{\prime}=\rho(q) \circ \psi_{\alpha}\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{w}_{p_{\alpha}}\right)$ holds for each element ${ }^{13}\left(\left(\vec{g}^{\prime}, \vec{s}^{\prime}\right), \vec{w}_{p_{\beta}}^{\prime}\right) \in$ $\mathfrak{q} \times T_{p_{\beta}} P_{\beta}$ with $^{14} \mathrm{~d}_{\left(e, p_{\beta}\right)} \Theta\left(\left(\vec{g}^{\prime}, \vec{s}^{\prime}\right), \vec{w}_{p_{\beta}}^{\prime}\right)=\mathrm{d} L_{q} \vec{w}_{p_{\alpha}}$. In fact, we have

$$
\mathrm{d}_{\left(e, p_{\beta}\right)} \Theta\left(\left(\vec{g}-\vec{g}^{\prime}, \vec{s}-\vec{s}^{\prime}\right), \vec{w}_{p_{\beta}}-\vec{w}_{p_{\beta}}^{\prime}\right)=0
$$

so that $a$ ) gives

$$
\begin{aligned}
& \left.0 \stackrel{a)}{=} \psi_{\beta}\left(\vec{g}-\vec{g}^{\prime}, \vec{w}_{p_{\beta}}-\vec{w}_{p_{\beta}}^{\prime}\right)-\left(\vec{s}-\vec{s}^{\prime}\right)\right)=\left[\psi_{\beta}\left(\vec{g}, \vec{w}_{p_{\beta}}\right)-\vec{s}\right]-\left[\psi_{\beta}\left(\vec{g}^{\prime}, \vec{w}_{p_{\beta}}^{\prime}\right)-\vec{s}^{\prime}\right] \\
& \quad=\rho(q) \circ \psi_{\alpha}\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{w}_{p_{\alpha}}\right)-\left[\psi_{\beta}\left(\vec{g}^{\prime}, \vec{w}_{p_{\beta}}^{\prime}\right)-\vec{s}^{\prime}\right] .
\end{aligned}
$$

3) Assume that $\mathrm{d} L_{q} \vec{w}_{p_{\alpha}} \in T_{p_{\beta}} P_{\beta}$ holds for all $q \in Q, p_{\alpha} \in P_{\alpha}, p_{\beta} \in P_{\beta}$ with $p_{\beta}=q \cdot p_{\alpha}$ and all $\vec{w}_{p_{\alpha}} \in T_{p_{\alpha}} P_{\alpha}$. Then $\mathrm{d}_{\left(e, p_{\beta}\right)} \Theta\left(\mathrm{d} L_{q}{\overrightarrow{w_{p}}}\right)=\mathrm{d} L_{q} \vec{w}_{p_{\alpha}}$ so that it follows from 2) that in this case we can substitute $i$ ) by $a$ ) and condition
b) Let $q \in Q, p_{\alpha} \in P_{\alpha}, p_{\beta} \in P_{\beta}$ with $p_{\beta}=q \cdot p_{\alpha}$. Then

$$
\psi_{\beta}\left(\overrightarrow{0}_{\mathfrak{g}}, \mathrm{d} L_{q} \vec{w}_{p_{\alpha}}\right)=\rho(q) \circ \psi_{\alpha}\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{w}_{p_{\alpha}}\right) \quad \forall \vec{w}_{p_{\alpha}} \in T_{p_{\alpha}} P_{\alpha} .
$$

[^6]Now, b) looks similar to $i i$ ) and makes it plausible that the conditions $i$ ) and $i i$ ) from Corollary 4.2 encode the $\rho$-invariance of the corresponding connection $\omega$. However, usually there is no reason for $\mathrm{d} L_{q} \vec{w}_{p_{\alpha}}$ to be an element of $T_{p_{\beta}} P_{\beta}$. Even for $p_{\alpha}=p_{\beta}$ and $q \in Q_{p_{\alpha}}$ this is not true in general. So, typically there is no way to split up $i$ ) into parts whose meaning is more intuitive.

## Remark 4.4 immediately proves

Case 4.5 (gauge fixing). Let $P_{0}$ be a $\Theta$-patch of the bundle $P$ such that $\pi\left(P_{0}\right)$ intersects each $\varphi$-orbit in a unique point and $\mathrm{d} L_{q}\left(T_{p} P_{0}\right) \subseteq T_{p} P_{0}$ for all $p \in P_{0}$ and all $q \in Q_{p}$. Then a corresponding reduced connection consists of one single smooth map $\psi: \mathfrak{g} \times T P_{0} \rightarrow \mathfrak{s}$ and we have $p=q \cdot p^{\prime}$ for $q \in Q, p, p^{\prime} \in P_{0}$ iff $p=p^{\prime}$ and $q \in Q_{p}$. Then, by Remark 4.4 the two conditions from Corollary 4.2 are equivalent to:

Let $p \in P_{0}, q=\left(j, \phi_{p}(j)\right) \in Q_{p}, \vec{w}_{p} \in T_{p} P_{0}$ and $\vec{g} \in \mathfrak{g}, \vec{s} \in \mathfrak{s}$. Then
$\left.i^{\prime}\right) \widetilde{g}(p)+\vec{w}_{p}-\widetilde{s}(p)=0 \quad \Longrightarrow \quad \psi\left(\vec{g}, \vec{w}_{p}\right)-\vec{s}=0$,
ii') $\psi\left(\overrightarrow{0}_{\mathfrak{g}}, \mathrm{d} L_{q} \vec{w}_{p}\right)=\rho(q) \circ \psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{w}_{p}\right)$,
$\left.i i i^{\prime}\right) \psi\left(\operatorname{Ad}_{j}(\vec{g}), \overrightarrow{0}_{p}\right)=\operatorname{Ad}_{\phi_{p}(j)} \circ \psi\left(\vec{g}, \overrightarrow{0}_{p}\right)$.
The next example is a slight generalization of Theorem 2 in [6]. Here the authors assume that $\varphi$ admits only one orbit type so that $\operatorname{dim}\left[G_{x}\right]=l$ for all $x \in M$. Then they restrict to the situation where we find a triple $\left(U_{0}, \tau_{0}, s_{0}\right)$ consisting of an open subset $U_{0} \subseteq \mathbb{R}^{k}$ for $k=\operatorname{dim}[M]-[\operatorname{dim}[G]-l]$, an embedding $\tau_{0}: U_{0} \rightarrow M$ and a smooth map $s_{0}: U_{0} \rightarrow P$ with $\pi \circ s_{0}=\tau_{0}$ and the addition property that $Q_{p}$ is the same for all $p \in \operatorname{im}\left[s_{0}\right]$. More precisely, they assume that $G_{x}$ and the structure group of the bundle are compact. Then they show the non-trivial fact that $s_{0}$ can be modified in such a way that in addition $Q_{p}$ is the same for all $p \in \operatorname{im}\left[s_{0}\right]$.

Observe that the authors omitted to require that $\operatorname{im}\left[\mathrm{d}_{x} \tau_{0}\right]+\operatorname{im}\left[\mathrm{d}_{e} \varphi_{\tau_{0}(x)}\right]=T_{\tau_{0}(x)} M$ holds for all $x \in U_{0}$, i.e., that $\tau_{0}\left(U_{0}\right)$ is a $\varphi$-patch (so that $s_{0}\left(U_{0}\right)$ is a $\Theta$-patch). Indeed, Example 4.10.2 shows that this additional condition is crucial. The next example is a slight modification of the result [6] in the sense that we do not assume $G_{x}$ and the structure group to be compact but make the ad hoc requirement that $Q_{p}$ is the same for all $p \in P_{0}$.
Example 4.6 (Harnad, Shnider, Vinet). Let $P_{0}$ be a $\Theta$-patch of the bundle $P$ such that $\pi\left(P_{0}\right)$ intersects each $\varphi$-orbit in a unique point, and assume that the $\Theta$-stabilizer $L:=Q_{p}$ is the same for all $p \in P_{0}$. Then it is clear from (2.4) that $H:=G_{\pi(p)}$ and $\phi:=\phi_{p}: H \rightarrow S$ are independent of the choice of $p \in P_{0}$. Finally, we require that

$$
\begin{equation*}
\operatorname{dim}\left[P_{0}\right]=\operatorname{dim}[M]-[\operatorname{dim}[G]-\operatorname{dim}[H]]=\operatorname{dim}[P]-[\operatorname{dim}[Q]-\operatorname{dim}[H]] . \tag{4.3}
\end{equation*}
$$

Now, let $p \in P_{0}$ and $q=(j, \phi(j)) \in Q_{p}$. Then for $\vec{w}_{p} \in T_{p} P_{0}$ we have

$$
\begin{aligned}
\mathrm{d} L_{q} \vec{w}_{p} & =\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} \Phi(j, \gamma(t)) \cdot \phi_{p}^{-1}(j)=\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0}\left[\gamma(t) \cdot \phi_{\gamma(t)}(j)\right] \cdot \phi_{p}^{-1}(j) \\
& =\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0}\left[\gamma(t) \cdot \phi_{p}(j)\right] \cdot \phi_{p}^{-1}(j)=\vec{w}_{p},
\end{aligned}
$$

where $\gamma:(-\epsilon, \epsilon) \rightarrow P_{0}$ is some smooth curve with $\dot{\gamma}(0)=\vec{w}_{p}$. Consequently, $\mathrm{d} L_{q}\left(T_{p} P_{0}\right) \subseteq T_{p} P_{0}$ so that we are in the situation of Case 4.5. Here $\left.i i^{\prime}\right)$ now reads $\psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{w}_{p}\right)=\operatorname{Ad}_{\phi(j)} \circ \psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{w}_{p}\right)$ for all $j \in H$ and $i i i^{\prime}$ ) does not change. For $i^{\prime}$ ) observe that the Lie algebra $\mathfrak{l}$ of $L$ is contained in the kernel of $\mathrm{d}_{\left(e, p_{0}\right)} \Theta$. But $\mathrm{d}_{\left(e, p_{0}\right)} \Theta$ is surjective since $P_{0}$ is a $\Theta$-patch ${ }^{15}$ so that

$$
\operatorname{dim}\left[\operatorname{ker}\left[\mathrm{d}_{\left(e, p_{0}\right)} \Theta\right]\right]=\operatorname{dim}[Q]+\operatorname{dim}\left[P_{0}\right]-\operatorname{dim}[P] \stackrel{(4.3)}{=} \operatorname{dim}[H] .
$$

This shows $\operatorname{ker}\left[\mathrm{d}_{(e, p)} \Theta\right]=\mathfrak{l}$ for all $p \in P_{0}$. Altogether, it follows that a reduced connection w.r.t. $P_{0}$ is a smooth, linear ${ }^{16}$ map $\psi: \mathfrak{g} \times T P_{0} \rightarrow \mathfrak{s}$ which fulfils the following three conditions:

[^7]\[

$$
\begin{array}{rll}
\left.i^{\prime \prime}\right) \psi\left(\vec{j}, \overrightarrow{0}_{p}\right) \stackrel{(4.2)}{=} \mathrm{d}_{e} \phi(\vec{j}) & \forall \vec{j} \in \mathfrak{h}, & \forall p \in P_{0}, \\
\left.i i^{\prime \prime}\right) \psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{w}\right)=\operatorname{Ad}_{\phi(j)} \circ \psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{w}\right) & \forall j \in H, & \forall \vec{w} \in T P_{0}, \\
\left.i i i^{\prime \prime}\right) \psi\left(\operatorname{Ad}_{j}(\vec{g}), \overrightarrow{0}_{p}\right)=\operatorname{Ad}_{\phi(j)} \circ \psi\left(\vec{g}, \overrightarrow{0}_{p}\right) & \forall \vec{g} \in \mathfrak{g}, & \forall j \in H, \quad \forall p \in P_{0} .
\end{array}
$$
\]

Then $\mu:=\left.\psi\right|_{T P_{0}}$ and $A_{p_{0}}(\vec{g}):=\psi\left(\vec{g}, \overrightarrow{0}_{p_{0}}\right)$ are the maps that are used for the characterization in Theorem 2 in [6].

### 4.2 Reconstruction of invariant connections

Let $\left\{P_{\alpha}\right\}_{\alpha \in I}$ be a $\Phi$-covering of $P$. We now show that each corresponding reduced connection $\left\{\psi_{\alpha}\right\}_{\alpha \in I}$ gives rise to a unique $\Phi$-invariant connection on $P$. To this end, for each $\alpha \in I$ we define the maps $\lambda_{\alpha}: \mathfrak{q} \times T P_{\alpha} \rightarrow \mathfrak{s},((\vec{g}, \vec{s}), \vec{w}) \mapsto \psi_{\alpha}(\vec{g}, \vec{w})-\vec{s}$ and

$$
\begin{aligned}
\omega_{\alpha}: & T Q \times T P_{\alpha} \rightarrow \mathfrak{s} \\
& \left(\vec{m}_{q}, \vec{w}_{p_{\alpha}}\right) \mapsto \rho(q) \circ \lambda_{\alpha}\left(\mathrm{d} L_{q^{-1}} \vec{m}_{q}, \vec{w}_{p_{\alpha}}\right)
\end{aligned}
$$

where $\vec{m}_{q} \in T_{q} Q$ and $\vec{w}_{p_{\alpha}} \in T_{p_{\alpha}} P_{\alpha}$.
Lemma 4.7. Let $q \in Q, p_{\alpha} \in P_{\alpha}, p_{\beta} \in P_{\beta}$ with $p_{\beta}=q \cdot p_{\alpha}$ and $\vec{w}_{p_{\alpha}} \in T_{p_{\alpha}} P_{\alpha}$. Then

1) $\lambda_{\beta}(\vec{\eta})=\rho(q) \circ \lambda_{\alpha}\left(\overrightarrow{0}_{\mathfrak{q}}, \vec{w}_{p_{\alpha}}\right)$ for all $\vec{\eta} \in \mathfrak{q} \times T_{p_{\beta}} P$ with $\mathrm{d} \Theta_{\left(e, p_{\beta}\right)}(\vec{\eta})=\mathrm{d} L_{q} \vec{w}_{p_{\alpha}}$,
2) $\lambda_{\beta}\left(\operatorname{Ad}_{q}(\vec{q}), \overrightarrow{0}_{p_{\beta}}\right)=\rho(q) \circ \lambda_{\alpha}\left(\vec{q}, \overrightarrow{0}_{p_{\alpha}}\right)$ for all $\vec{q} \in \mathfrak{q}$.

For all $\alpha \in I$ we have
3) $\operatorname{ker}\left[\left.\lambda_{\alpha}\right|_{\mathfrak{q} \times T_{p_{\alpha}} P_{\alpha}}\right] \subseteq \operatorname{ker}\left[\mathrm{d}_{\left(e, p_{\alpha}\right)} \Theta\right]$ for all $p_{\alpha} \in P_{\alpha}$,
4) the map $\omega_{\alpha}$ is the unique $\mathfrak{s}$-valued 1-form on $Q \times P_{\alpha}$ that extends $\lambda_{\alpha}$ and for which we have $L_{q}^{*} \omega_{\alpha}=\rho(q) \circ \omega_{\alpha}$ for all $q \in Q$.
Proof. 1. Write $\vec{\eta}=\left((\vec{g}, \vec{s}), \vec{w}_{p_{\beta}}\right)$ for $\vec{g} \in \mathfrak{g}, \vec{s} \in \mathfrak{s}$ and $\vec{w}_{p_{\beta}} \in T_{p_{\beta}} P_{\beta}$. Then

$$
\widetilde{g}\left(p_{\beta}\right)+\vec{w}_{p_{\beta}}-\widetilde{s}\left(p_{\beta}\right) \stackrel{(4.1)}{=} \mathrm{d} \Theta_{\left(e, p_{\beta}\right)}(\vec{\eta})=\mathrm{d} L_{q} \vec{w}_{p_{\alpha}}
$$

so that from condition $i$ ) in Corollary 4.2 we obtain

$$
\lambda_{\beta}(\vec{\eta})=\psi_{\beta}\left(\vec{g}, \vec{w}_{p_{\beta}}\right)-\vec{s}=\rho(q) \circ \psi_{\alpha}\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{w}_{p_{\alpha}}\right)=\rho(q) \circ \lambda_{\alpha}\left(\overrightarrow{0}_{\mathfrak{q}}, \vec{w}_{p_{\alpha}}\right)
$$

2. Let $\vec{q}=(\vec{g}, \vec{s})$ for $\vec{g} \in \mathfrak{g}$ and $\vec{s} \in \mathfrak{s}$. Then by Corollary 4.2.ii) we have

$$
\lambda_{\beta}\left(\operatorname{Ad}_{q}(\vec{q}), \overrightarrow{0}_{p_{\beta}}\right)=\psi_{\beta}\left(\operatorname{Ad}_{q}(\vec{g}), \overrightarrow{0}_{p_{\beta}}\right)-\operatorname{Ad}_{q}(\vec{s})=\rho(q) \circ\left[\psi_{\alpha}\left(\vec{g}, \overrightarrow{0}_{p_{\alpha}}\right)-\vec{s}\right]=\rho(q) \circ \lambda_{\alpha}\left(\vec{q}, \overrightarrow{0}_{p_{\alpha}}\right)
$$

3. This follows from the first part for $\alpha=\beta, q=e$ and $\vec{w}_{p_{\alpha}}=\overrightarrow{0}_{p_{\alpha}}$.
4. By definition we have $\left.\omega_{\alpha}\right|_{\mathfrak{q} \times T P_{\alpha}}=\lambda_{\alpha}$ and for the pullback property we calculate

$$
\begin{aligned}
\left(L_{q^{\prime}}^{*} \omega_{\alpha}\right)_{\left(q, p_{\alpha}\right)}\left(\vec{m}_{q}, \vec{w}_{p_{\alpha}}\right) & =\omega_{\alpha\left(q^{\prime} q, p_{\alpha}\right)}\left(\mathrm{d} L_{q^{\prime}} \vec{m}_{q}, \vec{w}_{p_{\alpha}}\right)=\rho\left(q^{\prime} q\right) \circ \lambda_{\alpha}\left(\mathrm{d} L_{\left.q^{-1}{q^{\prime-1}} \mathrm{~d} L_{q^{\prime}} \vec{m}_{q}, \vec{w}_{p_{\alpha}}\right)}\right. \\
& =\rho\left(q^{\prime}\right) \circ \rho(q) \circ \lambda_{\alpha}\left(\mathrm{d} L_{q^{-1}} \vec{m}_{q}, \vec{w}_{p_{\alpha}}\right)=\rho\left(q^{\prime}\right) \circ \omega_{\alpha\left(q, p_{\alpha}\right)}\left(\vec{m}_{q}, \vec{w}_{p_{\alpha}}\right)
\end{aligned}
$$

where $q, q^{\prime} \in Q$ and $\vec{m}_{q} \in T_{q} Q$. For uniqueness let $\omega$ be another $\mathfrak{s}$-valued 1-form on $Q \times P_{\alpha}$ whose restriction to $\mathfrak{q} \times T P_{\alpha}$ is $\lambda_{\alpha}$ and that fulfils $L_{q}^{*} \omega=\rho(q) \circ \omega$ for all $q \in Q$. Then

$$
\begin{aligned}
\omega_{\left(q, p_{\alpha}\right)}\left(\vec{m}_{q}, \vec{w}_{p_{\alpha}}\right) & =\omega_{\left(q, p_{\alpha}\right)}\left(\mathrm{d} L_{q} \circ \mathrm{~d} L_{q^{-1}} \vec{m}_{q}, \vec{w}_{p_{\alpha}}\right)=\left(L_{q}^{*} \omega\right)_{\left(e, p_{\alpha}\right)}\left(\mathrm{d} L_{q^{-1}} \vec{m}_{q}, \vec{w}_{p_{\alpha}}\right) \\
& =\rho(q) \circ \omega_{\left(e, p_{\alpha}\right)}\left(\mathrm{d} L_{q^{-1}} \vec{m}_{q}, \vec{w}_{p_{\alpha}}\right)=\rho(q) \circ \lambda_{\alpha}\left(\mathrm{d} L_{q^{-1}} \vec{m}_{q}, \vec{w}_{p_{\alpha}}\right) \\
& =\omega_{\alpha}\left(\mathrm{d} L_{q^{-1}} \vec{m}_{q}, \vec{w}_{p_{\alpha}}\right)
\end{aligned}
$$

Finally, smoothness of $\omega_{\alpha}$ is an easy consequence of smoothness of the maps $\rho, \lambda_{\alpha}$ and $\mu: T Q \rightarrow$ $\mathfrak{q}, \vec{m}_{q} \mapsto \mathrm{~d} L_{q^{-1}} \vec{m}_{q}$ with $\vec{m}_{q} \in T_{q} Q$. For this observe that $\mu=\mathrm{d} \tau \circ \kappa$ for $\tau: Q \times Q \rightarrow Q$, $\left(q, q^{\prime}\right) \mapsto q^{-1} q^{\prime}$ and $\kappa: T Q \rightarrow T Q \times T Q, \vec{m}_{q} \mapsto\left(\overrightarrow{0}_{q}, \vec{m}_{q}\right)$ for $\vec{m}_{q} \in T_{q} Q$.

So far, we have shown that each reduced connection $\left\{\psi_{\alpha}\right\}_{\alpha \in I}$ gives rise to uniquely determined maps $\left\{\lambda_{\alpha}\right\}_{\alpha \in I}$ and $\left\{\omega_{\alpha}\right\}_{\alpha \in I}$. In the final step we will construct a unique $\Phi$-invariant connection $\omega$ out of the data $\left\{\left(P_{\alpha}, \lambda_{\alpha}\right)\right\}_{\alpha \in I}$. Here, uniqueness and smoothness of $\omega$ will follow from uniqueness and smoothness of the maps $\omega_{\alpha}$.

Proposition 4.8. There is one and only one $\mathfrak{s}$-valued 1 -form $\omega$ on $P$ such that

$$
\omega_{\alpha}=\left.\left(\Theta^{*} \omega\right)\right|_{T Q \times T P_{\alpha}}
$$

holds for all $\alpha \in I$. Moreover, $\omega$ is a $\Phi$-invariant connection on $P$.
Proof. For uniqueness we have to show that the values of such an $\omega$ are uniquely determined by the maps $\omega_{\alpha}$. To this end, let $p \in P, \alpha \in I$ and $p_{\alpha} \in P_{\alpha}$ such that $p=q \cdot p_{\alpha}$ for some $q \in Q$. By Lemma 3.7.1 for $\vec{w}_{p} \in T_{p} P$ we find some $\vec{\eta} \in T_{q} Q \times T_{p_{\alpha}} P_{\alpha}$ with $\vec{w}_{p}=\mathrm{d}_{\left(q, p_{\alpha}\right)} \Theta(\vec{\eta})$, so that uniqueness follows from

$$
\omega_{p}\left(\vec{w}_{p}\right)=\omega_{q \cdot p_{\alpha}}\left(\mathrm{d}_{\left(q, p_{\alpha}\right)} \Theta(\vec{\eta})\right)=\left(\Theta^{*} \omega\right)_{\left(q, p_{\alpha}\right)}(\vec{\eta})=\omega_{\alpha}(\vec{\eta}) .
$$

For existence let $\alpha \in I$ and $p_{\alpha} \in P_{\alpha}$. Due to surjectivity of $\mathrm{d}_{\left(e, p_{\alpha}\right)} \Theta$ and Lemma 4.7.3 there is a (unique) map $\widehat{\lambda}_{p_{\alpha}}: T_{p_{\alpha}} P \rightarrow \mathfrak{s}$ with

$$
\begin{equation*}
\widehat{\lambda}_{p_{\alpha}} \circ \mathrm{d}_{\left(e, p_{\alpha}\right)} \Theta=\left.\lambda_{\alpha}\right|_{\mathfrak{q} \times T_{p_{\alpha}} P_{\alpha}} . \tag{4.4}
\end{equation*}
$$

Let $\widehat{\lambda}_{\alpha}: \bigsqcup_{p_{\alpha} \in P_{\alpha}} T_{p_{\alpha}} P \rightarrow \mathfrak{s}$ denote the (unique) map whose restriction to $T_{p_{\alpha}} P$ is $\widehat{\lambda}_{p_{\alpha}}$ for each $p_{\alpha} \in P_{\alpha}$. Then $\lambda_{\alpha}=\left.\widehat{\lambda}_{\alpha} \circ \mathrm{d} \Theta\right|_{\mathfrak{q} \times T P_{\alpha}}$ and we construct the connection $\omega$ as follows. For $p \in P$ we choose some $\alpha \in I$ and $\left(q, p_{\alpha}\right) \in Q \times P_{\alpha}$ such that $q \cdot p_{\alpha}=p$ and define

$$
\begin{equation*}
\omega_{p}\left(\vec{w}_{p}\right):=\rho(q) \circ \widehat{\lambda}_{\alpha}\left(\mathrm{d} L_{q^{-1}}\left(\vec{w}_{p}\right)\right) \quad \forall \vec{w}_{p} \in T_{p} P . \tag{4.5}
\end{equation*}
$$

We have to show that this depends neither on $\alpha \in I$ nor on the choice of $\left(q, p_{\alpha}\right) \in Q \times P_{\alpha}$. For this, let $p_{\alpha} \in P_{\alpha}, p_{\beta} \in P_{\beta}$ and $q \in Q$ with $p_{\beta}=q \cdot p_{\alpha}$. Then for $\vec{w} \in T_{p_{\alpha}} P$ we have $\vec{w}=\mathrm{d} \Theta\left(\vec{q}, \vec{w}_{p_{\alpha}}\right)$ for some $\left(\vec{q}, \vec{w}_{p_{\alpha}}\right) \in \mathfrak{q} \times T_{p_{\alpha}} P_{\alpha}$, and since $\mathrm{d} L_{q} \vec{w}_{p_{\alpha}} \in T_{p_{\beta}} P$, there is $\vec{\eta} \in \mathfrak{q} \times T_{p_{\beta}} P_{\beta}$ such that $\mathrm{d}_{\left(e, p_{\beta}\right)} \Theta(\vec{\eta})=\mathrm{d} L_{q} \vec{w}_{p_{\alpha}}$. It follows from the conditions 1 and 2 in Lemma 4.7 that

$$
\begin{align*}
\hat{\lambda}_{\beta}\left(\mathrm{d} L_{q} \vec{w}\right) & =\widehat{\lambda}_{\beta}\left(\left(\mathrm{d} L_{q} \circ \mathrm{~d} \Theta\right)\left(\vec{q}, \vec{w}_{p_{\alpha}}\right)\right)=\widehat{\lambda}_{\beta}\left(\left(\mathrm{d} L_{q} \circ \mathrm{~d} \Theta\right)\left(\vec{q}, \overrightarrow{0}_{p_{\alpha}}\right)\right)+\widehat{\lambda}_{\beta}\left(\mathrm{d} L_{q} \vec{w}_{p_{\alpha}}\right) \\
& \stackrel{(4.7)}{=} \hat{\lambda}_{\beta} \circ \mathrm{d} \Theta\left(\operatorname{Ad}_{q}(\vec{q}), \overrightarrow{0}_{p_{\beta}}\right)+\widehat{\lambda}_{\beta} \circ \mathrm{d} \Theta(\vec{\eta})  \tag{4.6}\\
& \stackrel{(4.4)}{=} \lambda_{\beta}\left(\operatorname{Ad}_{q}(\vec{q}), \overrightarrow{0}_{p_{\beta}}\right)+\lambda_{\beta}(\vec{\eta})=\rho(q) \circ \lambda_{\alpha}\left(\vec{q}, \overrightarrow{0}_{p_{\alpha}}\right)+\rho(q) \circ \lambda_{\alpha}\left(\overrightarrow{0}_{q}, \vec{w}_{p_{\alpha}}\right) \\
& =\rho(q) \circ \lambda_{\alpha}\left(\vec{q}, \vec{w}_{p_{\alpha}}\right)=\rho(q) \circ \hat{\lambda}_{\alpha} \circ \mathrm{d} \Theta\left(\vec{q}, \vec{w}_{p_{\alpha}}\right)=\rho(q) \circ \widehat{\lambda}_{\alpha}(\vec{w}),
\end{align*}
$$

where for the third equality we have used that

$$
\begin{align*}
\left(\mathrm{d} L_{q} \circ \mathrm{~d} \Theta\right)\left(\vec{q}, \overrightarrow{0}_{p_{\alpha}}\right) & =\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} q \cdot\left(\exp (t \vec{q}) \cdot p_{\alpha}\right)  \tag{4.7}\\
& =\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} \alpha_{q}(\exp (t \vec{q})) \cdot p_{\beta}=\mathrm{d} \Theta\left(\operatorname{Ad}_{q}(\vec{q}), \overrightarrow{0}_{p_{\beta}}\right) .
\end{align*}
$$

Consequently, if $\widetilde{q} \cdot p_{\beta}=p$ with $\left(\widetilde{q}, p_{\beta}\right) \in Q \times P_{\beta}$ for some $\beta \in I$, then $p_{\beta}=\left(q^{-1} \widetilde{q}\right)^{-1} \cdot p_{\alpha}$ and well-definedness follows from

$$
\begin{aligned}
\rho(\widetilde{q}) \circ \widehat{\lambda}_{\beta}\left(\mathrm{d} L_{\widetilde{q}^{-1}}\left(\vec{w}_{p}\right)\right) & =\rho(q) \circ \rho\left(q^{-1} \widetilde{q}\right) \circ \widehat{\lambda}_{\beta}\left(\mathrm{d} L_{\left(q^{-1} \widetilde{q}\right)^{-1}}\left(\mathrm{~d} L_{q^{-1}} \vec{w}_{p}\right)\right) \\
& =\rho(q) \circ \widehat{\lambda}_{\alpha}\left(\mathrm{d} L_{q^{-1}} \vec{w}_{p}\right),
\end{aligned}
$$

where the last step is due to (4.6) with $\vec{w}=\mathrm{d} L_{q^{-1}} \vec{w}_{p} \in T_{p_{\alpha}} P$. Next, we show that $\omega$ fulfils the pullback property. For this, let $\left(\vec{m}, \vec{w}_{p_{\alpha}}\right) \in T_{q} Q \times T_{p_{\alpha}} P_{\alpha}$. Then

$$
\begin{aligned}
\left(\Theta^{*} \omega\right)\left(\vec{m}_{q}, \vec{w}_{p_{\alpha}}\right) & =\omega_{q \cdot p_{\alpha}}\left(\mathrm{d} \Theta\left(\vec{m}_{q}, \vec{w}_{p_{\alpha}}\right)\right) \stackrel{(4.5)}{=} \rho(q) \circ \widehat{\lambda}_{\alpha}\left(\mathrm{d} L_{q^{-1}} \mathrm{~d} \Theta\left(\vec{m}_{q}, \vec{w}_{p_{\alpha}}\right)\right) \\
& =\rho(q) \circ \widehat{\lambda}_{\alpha} \circ \mathrm{d} \Theta\left(\mathrm{~d} L_{q^{-1}} \vec{m}_{q}, \vec{w}_{p_{\alpha}}\right) \stackrel{(4.4)}{=} \rho(q) \circ \lambda_{\alpha}\left(\mathrm{d} L_{q^{-1}} \vec{m}_{q}, \vec{w}_{p_{\alpha}}\right) \\
& =\omega_{\alpha}\left(\vec{m}_{q}, \vec{w}_{p_{\alpha}}\right) .
\end{aligned}
$$

In the third step we have used that $L_{q^{-1}} \circ \Theta=\Theta\left(L_{q^{-1}}(\cdot), \cdot\right)$. Finally, we have to verify that $\omega$ is a $\Phi$-invariant, smooth connection. For this let $p \in P$ and $\left(\widetilde{q}, p_{\alpha}\right) \in Q \times P_{\alpha}$ with $p=\widetilde{q} \cdot p_{\alpha}$. Then for $q \in Q$ and $\vec{w}_{p} \in T_{p} P$ we have

$$
\begin{aligned}
\left(L_{q}^{*} \omega\right)_{p}\left(\vec{w}_{p}\right) & =\omega_{q \cdot p}\left(\mathrm{~d} L_{q} \vec{w}_{p}\right)=\omega_{\left(q \widetilde{q} \cdot p_{\alpha}\right.}\left(\mathrm{d} L_{q} \vec{w}_{p}\right) \\
& =\rho(q) \circ \rho(\widetilde{q}) \circ \widehat{\lambda}_{\alpha}\left(\mathrm{d} L_{\widetilde{q}^{-1}} \vec{w}_{p}\right)=\rho(q) \circ \omega_{p}\left(\vec{w}_{p}\right),
\end{aligned}
$$

hence

$$
\begin{aligned}
& R_{s}^{*} \omega=L_{\left(e, s^{-1}\right)}^{*} \omega=\rho\left(\left(e, s^{-1}\right)\right) \circ \omega=\operatorname{Ad}_{s^{-1}} \circ \omega, \\
& L_{g}^{*} \omega=L_{(g, e)}^{*} \omega=\rho((g, e)) \circ \omega=\omega .
\end{aligned}
$$

So, it remains to show smoothness of $\omega$ and that $\omega_{p}(\widetilde{s}(p))=\vec{s}$ holds for all $p \in P$ and all $\vec{s} \in \mathfrak{s}$. For the second property let $p=q \cdot p_{\alpha}$ for $\left(q, p_{\alpha}\right) \in Q \times P_{\alpha}$. Then $q=(g, s)$ for some $g \in G$ and $s \in S$ and we obtain

$$
\begin{aligned}
\omega_{p}(\widetilde{s}(p)) & =\rho(q) \circ \widehat{\lambda}_{\alpha}\left(\mathrm{d} L_{q^{-1}} \widetilde{s}\left(q \cdot p_{\alpha}\right)\right)=\rho(q) \circ \widehat{\lambda}_{\alpha}\left(\left.\frac{\mathrm{d}}{\mathrm{~d} \mathrm{~d}}\right|_{t=0} p_{\alpha} \cdot\left(\alpha_{s^{-1}}(\exp (t \vec{s}))\right)\right. \\
& =\rho(q) \circ \widehat{\lambda}_{\alpha}\left(\mathrm{d} \Theta\left(\operatorname{Ad}_{s^{-1}}(\vec{s}), \overrightarrow{0}_{p_{\alpha}}\right)\right)=\operatorname{Ad}_{s^{\prime}} \circ \lambda_{\alpha}\left(\operatorname{Ad}_{s^{-1}}(\vec{s}), \overrightarrow{0}_{p_{\alpha}}\right)=\operatorname{Ad}_{s} \circ \operatorname{Ad}_{s^{-1}}(\vec{s})=\vec{s} .
\end{aligned}
$$

For smoothness let $p_{\alpha} \in P_{\alpha}$ and choose a submanifold $Q^{\prime}$ of $Q$ through $e$, an open neighbourhood $P_{\alpha}^{\prime} \subseteq P_{\alpha}$ of $p_{\alpha}$ and an open subset $U \subseteq P$ such that the restriction $\widehat{\Theta}:=\left.\Theta\right|_{Q^{\prime} \times P_{\alpha}^{\prime}}$ is a diffeomorphism to $U$. Then $p_{\alpha} \in U$ because $e \in Q^{\prime}$, hence

$$
\left.\omega\right|_{U}=\widehat{\Theta}^{-1 *}\left[\widehat{\Theta}^{*} \omega\right]=\widehat{\Theta}^{-1 *}\left[\left.\left(\Theta^{*} \omega\right)\right|_{T Q \times T P_{\alpha}}\right]=\widehat{\Theta}^{-1 *} \omega_{\alpha} .
$$

Since $\omega_{\alpha}$ is smooth and $\widehat{\Theta}$ is a diffeomorphism, $\left.\omega\right|_{U}$ is smooth as well. Finally, if $p=q \cdot p_{\alpha}$ for $q \in Q$, then $L_{q}(U)$ is an open neighbourhood of $p$ and

$$
\left.\omega\right|_{L_{q}(U)}=\left.\left(L_{q^{-1}}^{*}\left(L_{q}^{*} \omega\right)\right)\right|_{L_{q}(U)}=\left.\rho(q) \circ\left(L_{q^{-1}}^{*} \omega\right)\right|_{L_{q}(U)}=\rho(q) \circ L_{q^{-1}}^{*}\left(\left.\omega\right|_{U}\right)
$$

is smooth because $\left.\omega\right|_{U}$ and $L_{q^{-1}}$ are smooth.
Corollary 4.2 and Proposition 4.8 now prove
Theorem 4.9. Let $G$ be a Lie group of automorphisms of the principal fibre bundle $P$. Then for each $\Phi$-covering $\left\{P_{\alpha}\right\}_{\alpha \in I}$ of $P$ there is a bijection between the corresponding set of reduced connections and the $\Phi$-invariant connections on $P$.

As already mentioned in the preliminary remarks to Example 4.6, the second part of the next example shows the importance of the transversality condition

$$
\operatorname{im}\left[\mathrm{d}_{x} \tau_{0}\right]+\operatorname{im}\left[\mathrm{d}_{e} \varphi_{\tau_{0}(x)}\right]=T_{\tau_{0}(x)} M \quad \forall x \in U_{0}
$$

for the formulation in [6].

Example 4.10 ((semi-)homogeneous connections).

1. Let $P=X \times S$ for an $n$-dimensional $\mathbb{R}$-vector space $X$ and an arbitrary structure group $S$. Moreover, let $G \subseteq X$ be a linear subspace of dimension $1 \leq k \leq n$ acting via $\Phi: G \times P \rightarrow P$, $(g,(x, \sigma)) \mapsto(g+x, \sigma)$. If $W$ is an algebraic complement of $G$ in $X$ and $P_{0}:=W \times\left\{e_{S}\right\} \subseteq P$, then $P_{0}$ is a $\Phi$-covering since $\Theta:(G \times S) \times P_{0} \rightarrow P$ is a diffeomorphism and each $\varphi$-orbit intersects $W$ in a unique point. Consequently, the $\Phi$-invariant connections on $P$ are in bijection with the smooth maps $\psi: G \times T W \rightarrow \mathfrak{s}$ such that $\psi_{w}:=\left.\psi\right|_{G \times T_{w} W}$ is linear for all $w \in W$. This is because the conditions $i$ ) and $i i$ ) from Corollary 4.2 give no further restrictions in this case. The $\Phi$-invariant connection that corresponds to $\psi$ is given by

$$
\begin{equation*}
\omega_{(x, s)}^{\psi}\left(\vec{v}_{x}, \vec{\sigma}_{s}\right)=\operatorname{Ad}_{s^{-1}} \circ \psi_{\operatorname{pr}_{W}(x)}\left(\operatorname{pr}_{G}\left(\vec{v}_{x}\right), \operatorname{pr}_{W}\left(\vec{v}_{x}\right)\right)+\mathrm{d} L_{s^{-1}}\left(\vec{\sigma}_{s}\right) \tag{4.8}
\end{equation*}
$$

for $\left(\vec{v}_{x}, \vec{\sigma}_{s}\right) \in T_{(x, s)} P$.
2. In the situation of the previous part let $X=\mathbb{R}^{2}, G=\operatorname{span}_{\mathbb{R}}\left(\vec{e}_{1}\right), W=\operatorname{span}_{\mathbb{R}}\left(\vec{e}_{2}\right)$ and $P_{0}=W \times\{e\}$. We fix $0 \neq \vec{s} \in \mathfrak{s}$ and define $\psi: \mathfrak{g} \times T P_{0} \rightarrow \mathfrak{s}$ by

$$
\psi_{y}\left(\lambda \vec{e}_{1}, \mu \vec{e}_{2}\right):=\mu f(y) \cdot \vec{s} \quad \text { for } \quad\left(\lambda \vec{e}_{1}, \mu \vec{e}_{2}\right) \in \mathfrak{g} \times T_{\left(y \cdot \vec{e}_{2}, e\right)} P_{0},
$$

where $f(0):=0$ and $f(y):=1 / \sqrt[3]{y}$ for $y \neq 0$. Then $\omega^{\psi}$ defined by (4.8) is not smooth at the origin because

$$
\omega_{((x, y), e)}^{\psi}\left(\left(\overrightarrow{0}, \vec{e}_{2}\right), \overrightarrow{0}_{\mathfrak{s}}\right)=\psi_{y}\left(\overrightarrow{0}, \vec{e}_{2}\right)=f(y) \cdot \vec{s} \quad \forall y \in \mathbb{R}
$$

Now, let $U_{0}=\mathbb{R}, \tau_{0}: U_{0} \rightarrow \mathbb{R}^{2}, t \mapsto\left(t, t^{3}\right)$ and $s_{0}: t \mapsto\left(\tau_{0}(t), e\right)$. Then $\left(U_{0}, \tau_{0}, s_{0}\right)$ fulfils the conditions in [6] but we have $\operatorname{im}\left[\mathrm{d}_{0} \tau_{0}\right]+\operatorname{im}\left[\mathrm{d}_{e} \varphi_{\tau_{0}(0)}\right]=\operatorname{span}_{\mathbb{R}}\left(\vec{e}_{1}\right) \neq T_{0} X=T_{0} \mathbb{R}^{2}=\mathbb{R}^{2} .{ }^{17}$ As a consequence, $\psi^{\prime}: \mathfrak{g} \times T U_{0} \rightarrow \mathfrak{s}$ defined by $\psi_{t}^{\prime}:=\left.\left(\Phi^{*} \omega^{\psi}\right)\right|_{\mathfrak{g} \times T_{t} U_{0}}$ is smooth because for $t \neq 0$ and $r \in T_{t} U_{0}=\mathbb{R}$ we have

$$
\begin{aligned}
\psi_{t}^{\prime}\left(\lambda \vec{e}_{1}, r\right) & =\left(\Phi^{*} \omega^{\psi}\right)\left(\lambda \vec{e}_{1}, r \cdot \vec{e}_{1}+3 t^{2} r \cdot \vec{e}_{2}\right)=\omega_{\left(\left(t, t^{3}\right), e\right)}^{\psi}\left((\lambda+r) \cdot \vec{e}_{1}+3 t^{2} r \cdot \vec{e}_{2}\right) \\
& =\psi_{t^{3}}\left((\lambda+r) \cdot \vec{e}_{1}, 3 t^{2} r \cdot \vec{e}_{2}\right)=3 t r \cdot \vec{s},
\end{aligned}
$$

as well as $\psi_{0}^{\prime}\left(\lambda \vec{e}_{1}, r\right)=0$ if $t=0$. For the first step keep in mind that $\left.\left(\Phi^{*} \omega^{\psi}\right)\right|_{\mathfrak{g} \times T_{t} U_{0}}(\vec{g}, r)=$ $\left(\Phi^{*} \omega^{\psi}\right)\left(\vec{g}, \mathrm{~d}_{t} s_{0}(r)\right)$ by Convention 3.1.2. Then the maps $\mu:=\left.\psi^{\prime}\right|_{T U_{0}}$ and $A_{t_{0}}(\vec{g}):=$ $\psi\left(\vec{g}, \overrightarrow{0}_{t_{0}}\right)$ fulfil the conditions from Theorem 2 in [6] because $\psi^{\prime}$ fulfils the three algebraic conditions in Example 4.6 ${ }^{18}$. This, however, contradicts that $\omega^{\psi}$ is not a smooth connection.

## 5 Particular cases and applications

In the first part of this section we consider $\Phi$-coverings of $P$ that arise from the induced action $\varphi$ on the base manifold $M$ of $P$. Then we discuss the case where $\Phi$ acts via gauge transformations on $P$. This leads to a straightforward generalization of the description of connections by consistent families of local 1 -forms on $M$. In the second part we discuss the (almost-)fibre transitive case and deduce Wang's original theorem [8] from Theorem 4.9. Finally, we consider the situation where $P$ is trivial and give examples in loop quantum gravity.

[^8]
## 5.1 $\quad \Phi$-coverings and the induced action

Let $(G, \Phi)$ be a Lie group of automorphisms of the principal fibre bundle $P$. According to Lemma 3.4 for each $x \in M$ there is a $\varphi$-patch (with minimal dimension) $M_{x}$ with $x \in M$. Consequently, there is an open neighbourhood $M_{x}^{\prime} \subseteq M_{x}$ of $x$ and a local section $s_{x}: U \rightarrow P$ such that $M_{x}^{\prime} \subseteq U$ for an open neighbourhood $U \subseteq M$. Let $I \subseteq M$ be a subset such that ${ }^{19}$ each $\varphi$-orbit intersects at least one of the sets $M_{x}$. Then it is immediate from Lemma 3.7.2 that $\left\{s_{x}\left(M_{x}^{\prime}\right)\right\}_{x \in I}$ is a $\Phi$-covering of $P$. More generally, we have

Corollary 5.1. Let $(P, \pi, M, S)$ be a principal fibre bundle and $(G, \Phi)$ a Lie group of automorphisms of P. Denote by $\left(M_{\alpha}, s_{\alpha}\right)_{\alpha \in I}$ a family consisting of a collection of $\varphi$-patches $\left\{M_{\alpha}\right\}_{\alpha \in I}$ and smooth sections ${ }^{20} s_{\alpha}: M_{\alpha} \rightarrow P$. Then the sets $P_{\alpha}:=s_{\alpha}\left(M_{\alpha}\right)$ are $\Theta$-patches. They provide $a \Phi$-covering of $P$ iff each $\varphi$-orbit intersects at least one patch $M_{\alpha}$.

Proof. This is immediate from Lemma 3.7.2.
We now consider the case where $(G, \Phi)$ is a Lie group of gauge transformations of $P$, i.e., $\varphi_{g}=\mathrm{id}_{M}$ for all $g \in G$. Here we show that Theorem 4.9 can be seen as a generalization of the description of smooth connections by consistent families of local 1-forms on the base manifold $M$. For this let $\left\{U_{\alpha}\right\}_{\alpha \in I}$ be an open covering of $M$ and $\left\{s_{\alpha}\right\}_{\alpha \in I}$ a family of smooth sections $s_{\alpha}: U_{\alpha} \rightarrow P$. We define $U_{\alpha \beta}:=U_{\alpha} \cap U_{\beta}$ and consider the smooth maps $\delta_{\alpha \beta}: G \times$ $U_{\alpha \beta} \rightarrow S$ determined by $s_{\beta}(x)=\Phi\left(g, s_{\alpha}(x)\right) \cdot \delta_{\alpha \beta}(g, x) .{ }^{21}$ Finally, let $\mu_{\alpha \beta}\left(g, \vec{v}_{x}\right):=\mathrm{d} L_{\delta_{\alpha \beta}^{-1}(g, x)} \circ$ $\mathrm{d}_{x} \delta_{\alpha \beta}(g, \cdot)\left(\vec{v}_{x}\right)$ for $\vec{v}_{x} \in T_{x} U_{\alpha \beta}$ and $g \in G$. Then we have

Case 5.2 (Lie groups of gauge transformations). Let $(G, \Phi)$ be a Lie group of gauge transformations of the principal fibre bundle $(P, \pi, M, S)$. Then the $\Phi$-invariant connections on $P$ are in bijection with the families $\left\{\chi_{\alpha}\right\}_{\alpha \in I}$ of $\mathfrak{s}$-valued 1-forms $\chi_{\alpha}: U_{\alpha} \rightarrow \mathfrak{s}$ for which we have

$$
\begin{equation*}
\chi_{\beta}\left(\vec{v}_{x}\right)=\left(\operatorname{Ad}_{\delta_{\alpha \beta}(g, x)} \circ \chi_{\alpha}\right)\left(\vec{v}_{x}\right)+\mu_{\alpha \beta}\left(g, \vec{v}_{x}\right) \quad \forall \vec{v}_{x} \in T_{x} U_{\alpha \beta}, \quad \forall g \in G . \tag{5.1}
\end{equation*}
$$

Proof. By Corollary 5.1 $\left\{s_{\alpha}\left(U_{\alpha}\right)\right\}_{\alpha \in I}$ is a $\Phi$-covering of $P$. So, let $\left\{\psi_{\alpha}\right\}_{\alpha \in I}$ be a reduced connection w.r.t. this covering. We first show that condition $i$ ) from Corollary 4.2 implies

$$
\psi_{\beta}\left(\vec{g}, \overrightarrow{0}_{p}\right)=\mathrm{d}_{e} \phi_{p}(\vec{g}) \quad \forall \vec{g} \in \mathfrak{g}, \quad \forall p \in s_{\beta}(U)
$$

For this observe that condition $a$ ) from Remark 4.4 means that for all $\beta \in I, p \in s_{\beta}\left(U_{\beta}\right)$, $\vec{w}_{p} \in T_{p} s_{\beta}\left(U_{\beta}\right)$ and $\vec{g} \in \mathfrak{g}, \vec{s} \in \mathfrak{s}$ we have

$$
\mathrm{d}_{e} \Phi_{p}(\vec{g})+\vec{w}_{p}-\widetilde{s}(p)=0 \quad \Longrightarrow \quad \psi_{\beta}\left(\vec{g}, \vec{w}_{p}\right)-\vec{s}=0 .
$$

Now, $T_{p} s_{\beta}\left(U_{\beta}\right)$ is complementary to $T v_{p} P$ and $\operatorname{im}\left[\mathrm{d}_{e} \Phi_{p}\right] \subseteq \operatorname{ker}\left[\mathrm{d}_{p} \pi\right]$ so that $a$ ) is the same as

$$
\left.a^{\prime}\right) \mathrm{d}_{e} \Phi_{p}(\vec{g})=\widetilde{s}(p) \quad \Longrightarrow \quad \psi_{\beta}\left(\vec{g}, \overrightarrow{0}_{p}\right)=\vec{s} \quad \text { for } \vec{g} \in \mathfrak{g}, \vec{s} \in \mathfrak{s} \text { and all } p \in P_{\beta} .
$$

But, since $G_{x}=G$ for all $x \in M$, this just means ${ }^{22} \psi_{\beta}\left(\vec{g}, \overrightarrow{0}_{p}\right)=\mathrm{d}_{e} \phi_{p}(\vec{g})$ for all $\vec{g} \in \mathfrak{g}$ and is in line with condition $i i$ ) from Corollary 4.2 as $\phi_{p}$ is a Lie group homomorphism. Consequently, we can ignore this condition in the following. Now, we have $p_{\beta}=q \cdot p_{\alpha}$ for $q \in Q, p_{\alpha} \in P_{\alpha}$, $p_{\beta} \in P_{\beta}$ iff $\pi\left(p_{\alpha}\right)=\pi\left(p_{\beta}\right)=x \in U_{\alpha \beta}$ and $q=\left(g, \delta_{\alpha \beta}^{-1}(g, x)\right)$. Consequently, the left hand side of condition $i$ ) from Corollary 4.2 reads

$$
\widetilde{g}\left(s_{\beta}(x)\right)+\mathrm{d}_{x} s_{\beta}\left(\vec{v}_{\beta}\right)-\widetilde{s}\left(s_{\beta}(x)\right)=\left(\mathrm{d} L_{g} \circ \mathrm{~d} R_{\delta_{\alpha \beta}(g, x)} \circ \mathrm{d}_{x} s_{\alpha}\right)\left(\vec{v}_{\alpha}\right),
$$

[^9]where $\vec{v}_{\alpha}, \vec{v}_{\beta} \in T_{x} M$ and $g \in G$. This is true for $\vec{v}_{\alpha}=\vec{v}_{\beta}=\vec{v}_{x}, \vec{g}=0$ and $\vec{s}=\mu_{\alpha \beta}\left(g, \vec{v}_{x}\right)$, which follows from
\[

$$
\begin{aligned}
\mathrm{d}_{x} s_{\beta}\left(\vec{v}_{\beta}\right) & =\mathrm{d}_{x}\left[L_{g} \circ R_{\delta_{\alpha \beta}(g, \cdot)} \circ s_{\alpha}\right]\left(\vec{v}_{x}\right) \\
& =\mathrm{d} L_{g}\left[\mathrm{~d}_{s_{\alpha}(x)} R\left(\mathrm{~d}_{x} \delta_{\alpha \beta}(g, \cdot)\left(\vec{v}_{x}\right)\right)+\mathrm{d} R_{\delta_{\alpha \beta}(g, x)}\left(\mathrm{d}_{x} s_{\alpha}\left(\vec{v}_{x}\right)\right)\right] \\
\widetilde{s}\left(s_{\beta}(x)\right) & =\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} L_{g} \circ R_{\delta_{\alpha \beta}(g, x) \cdot \exp (t \vec{s})}\left(s_{\alpha}(x)\right) \\
& =\mathrm{d} L_{g}\left[\mathrm{~d}_{s_{\alpha}(x)} R\left(\mathrm{~d} L_{\delta_{\alpha \beta}(g, x)}(\vec{s})\right)\right]=\mathrm{d} L_{g}\left[\mathrm{~d}_{s_{\alpha}(x)} R\left(\mathrm{~d}_{x} \delta_{\alpha \beta}(g, \cdot)\left(\vec{v}_{x}\right)\right)\right] .
\end{aligned}
$$
\]

Consequently, by Corollary $4.2 . i$ ) we have ${ }^{23}$

$$
\begin{equation*}
\psi_{\beta}\left(\overrightarrow{0}_{\mathfrak{g}}, \mathrm{d}_{x} s_{\beta}\left(\vec{v}_{x}\right)\right)=\left(\operatorname{Ad}_{\delta_{\alpha \beta}(g, x)} \circ \psi_{\alpha} \circ \mathrm{d}_{x} s_{\alpha}\right)\left(\vec{v}_{x}\right)+\mu_{\alpha \beta}\left(g, \vec{v}_{x}\right) \tag{5.2}
\end{equation*}
$$

for all $g \in G$ and all $\vec{v}_{x} \in T_{x} U_{\alpha \beta}$. Due to part 2) in Remark 4.4 the condition $i$ ) from Corollary 4.2 now gives no further restrictions so that for $\chi_{\beta}:=\psi_{\beta} \circ \mathrm{d} s_{\beta}$ we have

$$
\psi_{\beta}\left(\vec{g}, \mathrm{~d}_{x} s_{\beta}\left(\vec{v}_{x}\right)\right)=\mathrm{d}_{e} \phi_{s_{\beta}(x)}(\vec{g})+\chi_{\beta}\left(\vec{v}_{x}\right) \quad \forall \vec{g} \in \mathfrak{g}, \quad \forall \vec{v}_{x} \in T_{x} M, \quad \forall x \in U_{\beta}
$$

Then $\psi_{\beta}$ is uniquely determined by $\chi_{\beta}$ for each $\beta \in I$ so that (5.2) yields the consistency condition (5.1) for the maps $\left\{\chi_{\alpha}\right\}_{\alpha \in I}$.

Example 5.3 (trivial action). If $G$ acts trivially, then for each $x \in U_{\alpha \beta}$ we have

$$
\delta_{\alpha \beta}(g, x)=\phi_{s_{\alpha}(x)}^{-1}(g) \cdot \delta_{\alpha \beta}(e, x)=\delta_{\alpha \beta}(e, x)
$$

so that $\delta_{\alpha \beta}$ is independent of $g \in G$. Here Case 5.2 just reproduces the description of smooth connections by means of consistent families of local 1-forms on the base manifold $M$.

## 5.2 (Almost-)fibre transitivity

In this subsection we discuss the situation where $M$ admits an element that is contained in the closure of each $\varphi$-orbit. For instance, this holds for all $x \in M$ if each $\varphi$-orbit is dense in $M$ and, in particular, is true for fibre transitive actions.

Case 5.4 (almost-fibre transitivity). Let $x \in M$ be contained in the closure of each $\varphi$-orbit and let $p \in \pi^{-1}(x)$. Then each $\Theta$-patch $P_{0} \subseteq P$ with $p \in P_{0}$ is a $\Phi$-covering of $P$. Hence, the $\Phi$-invariant connections on $P$ are in bijection with the smooth maps $\psi: \mathfrak{g} \times T P_{0} \rightarrow \mathfrak{s}$ for which $\left.\psi\right|_{\mathfrak{g} \times T_{p} P_{0}}$ is linear for all $p \in P_{0}$ and that fulfil the two conditions from Corollary 4.2.

Proof. It suffices to show that $\pi\left(P_{0}\right)$ intersects each $\varphi$-orbit [o]. Since $P_{0}$ is a $\Theta$-patch, there is an open neighbourhood $P^{\prime} \subseteq P_{0}$ of $p$ and a submanifold $Q^{\prime}$ of $Q$ through $\left(e_{G}, e_{S}\right)$ such that $\left.\Theta\right|_{Q^{\prime} \times P^{\prime}}$ is a diffeomorphism to an open subset $U \subseteq P$. Then $\pi(U)$ is an open neighbourhood of $\pi(p)$ and by assumption we have $[o] \cap \pi(U) \neq \varnothing$ for each $[o] \in M / G$. Consequently, for $[o] \in M / G$ we find $\widetilde{p} \in U$ with $\pi(\widetilde{p}) \in[o]$. Let $\widetilde{p}=\Theta\left(\left(g^{\prime}, s^{\prime}\right), p^{\prime}\right)$ for $\left(\left(g^{\prime}, s^{\prime}\right), p^{\prime}\right) \in Q^{\prime} \times P^{\prime}$. Then

$$
[o] \ni \pi(\widetilde{p})=\pi\left(\Phi\left(g^{\prime}, p^{\prime}\right) \cdot s^{\prime}\right)=\varphi\left(g^{\prime}, \pi\left(p^{\prime}\right)\right) \in\left[\pi\left(p^{\prime}\right)\right]
$$

shows that $[o]=\left[\pi\left(p^{\prime}\right)\right]$, hence $\pi\left(P_{0}\right) \cap[o] \neq \varnothing$.
The next example to Case 5.4 shows that evaluating the conditions $i$ ) and $i i$ ) from Corollary 4.2 at one single point can be sufficient to verify non-existence of invariant connections.

[^10]Example 5.5 (general linear group).

1. Let $P:=\mathrm{GL}(n, \mathbb{R})$ and $G=S=B \subseteq \mathrm{GL}(n, \mathbb{R})$ the subgroup of upper triangular matrices. Moreover, let $S_{n} \subseteq \mathrm{GL}(n, \mathbb{R})$ be the group of permutation matrices. Then $P$ is a principal fibre bundle with base manifold $M:=P / S$, structure group $S$ and projection map $\pi: P \rightarrow M, p \mapsto[p]$. Moreover, $G$ acts via automorphisms on $P$ by $\Phi(g, p):=g \cdot p$ and we have the Bruhat decomposition

$$
\operatorname{GL}(n, \mathbb{R})=\bigsqcup_{w \in S_{n}} B w B
$$

Then $M=\bigsqcup_{w \in S_{n}} G \cdot \pi(w), G \cdot \pi(e)=\pi(e)$ and $\pi(e) \in \overline{G \cdot \pi(w)}$ for all $w \in S_{n}$. Now, $\operatorname{im}\left[\mathrm{d}_{e} \Theta_{e}\right]=\mathfrak{g}$, since $\mathrm{d}_{e} \Theta_{e}(\vec{g})=\vec{g}$ for all $\vec{g} \in \mathfrak{g}$. Moreover, $\mathfrak{g}=\operatorname{span}_{\mathbb{R}}\left\{E_{i j} \mid 1 \leq i \leq j \leq n\right\}$ so that $V:=\operatorname{span}_{\mathbb{R}}\left\{E_{i j} \mid 1 \leq j<i \leq n\right\}$ is an algebraic complement of $\mathfrak{g}$ in $T_{e} P=\mathfrak{g l}(n, \mathbb{R})$. By Lemma 3.4.2 we find a patch $H \subseteq P$ through $e$ with $T_{e} H=V$ and due to Case 5.4 this is a $\Phi$-covering.
2. A closer look at the point $e \in P$ shows that there cannot exist any $\Phi$-invariant connection on $\operatorname{GL}(n, \mathbb{R})$. In fact, if $\psi: \mathfrak{g} \times T H \rightarrow \mathfrak{s}$ is a reduced connection w.r.t. $H$, then for $\vec{w}:=\overrightarrow{0}_{e}$ and $\vec{g}=\vec{s}$ we have

$$
\widetilde{g}(e)+\vec{w}-\widetilde{s}(e)=\vec{g}+\vec{w}-\vec{s}=0
$$

So, condition $i$ from Corollary 4.2 gives $\psi\left(\vec{g}, \overrightarrow{0}_{e}\right)-\vec{g}=0$, hence $\psi\left(\vec{g}, \overrightarrow{0}_{e}\right)=\vec{g}$ for all $\vec{g} \in \mathfrak{g}$. Now $q \cdot e=e$ iff $q=(b, b)$ for some $b \in B$. Let

$$
V \ni \vec{h}:=E_{n 1}, \quad B \ni b:=\mathbb{1}+E_{1 n}, \quad \mathfrak{g} \ni \vec{g}:=E_{11}-E_{1 n}-E_{n n} .
$$

Then $\widetilde{g}(e)+\vec{h}=\vec{g}+\vec{h}=b \vec{h} b^{-1}=\mathrm{d} L_{q} \vec{h}$ so that condition $i$ ) yields

$$
\psi(\vec{g}, \vec{h})=\rho(q) \circ \psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{h}\right)=\operatorname{Ad}_{b} \circ \psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{h}\right)
$$

hence $\vec{g}+\left[\mathrm{id}-\operatorname{Ad}_{b}\right] \circ \psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{h}\right)=0$. But $(\vec{g})_{11}=1$ and

$$
\left(\psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{h}\right)-\operatorname{Ad}_{b} \circ \psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{h}\right)\right)_{11}=\left(\psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{h}\right)\right)_{11}-\left(\psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{h}\right)\right)_{11}=0
$$

so that $\psi$ cannot exist.
Corollary 5.6. If $\Phi$ is fibre transitive, then $\{p\}$ is a $\Phi$-covering for all $p \in P$.
Proof. It suffices to show that $\{\pi(p)\}$ is a $\varphi$-patch, since then $\{p\}$ is a $\Theta$-patch by Corollary 5.1 and a $\Phi$-covering by Case 5.4. This, however, is clear from Remark 3.3.1. In fact, if $x:=\pi(p)$, then by general theory we know that $M$ is diffeomorphic to $G / G_{x}$ via $\vartheta:[g] \mapsto \varphi(g, x)$ and that for each $[g] \in G / G_{x}$ we find an open neighbourhood $U \subseteq G / G_{x}$ of $[g]$ and a smooth section $s: U \rightarrow G$. Then surjectivity of $\mathrm{d}_{e} \varphi_{x}$ is clear from surjectivity of $\mathrm{d}_{[e]} \vartheta$ and

$$
\mathrm{d}_{e} \varphi_{x} \circ \mathrm{~d}_{[e]} s=\mathrm{d}_{[e]}\left(\varphi_{x} \circ s\right)=\mathrm{d}_{[e]} \varphi(s(\cdot), x)=\mathrm{d}_{[e]} \vartheta
$$

showing that $T_{x} M=\mathrm{d}_{e} \varphi_{x}(\mathfrak{g})$.
Let $\phi$ be transitive and $p \in P$. Then $\{p\}$ is a $\Phi$-covering by Corollary 5.6 and $T_{p}\{p\}$ is the zero vector space. Moreover, we have $p_{\alpha}=q \cdot p_{\beta}$ iff $p_{\alpha}=p_{\beta}=p$ and $q \in Q_{p}$. It follows that a reduced connection w.r.t. $\{p\}$ can be seen as a linear map $\psi: \mathfrak{g} \rightarrow \mathfrak{s}$ that fulfils the following two conditions:

- $\mathrm{d}_{e} \Theta_{p}(\vec{g}, \vec{s})=0 \quad \Longrightarrow \quad \psi(\vec{g})=\vec{s} \quad$ for $\vec{g} \in \mathfrak{g}, \vec{s} \in \mathfrak{s}$,
- $\psi\left(\operatorname{Ad}_{q}(\vec{g})\right)=\rho(q) \circ \psi(\vec{g}) \quad \forall q \in Q_{p}, \quad \forall \vec{g} \in \mathfrak{g}$.

Since $\operatorname{ker}\left[d_{e} \Theta_{p}\right]=\mathfrak{q}_{p}$, we have shown

Case 5.7 (Hsien-Chung Wang, [8]). Let $(G, \Phi)$ be a fibre transitive Lie group of automorphisms of $P$. Then for each $p \in P$ there is a bijection between the $\Phi$-invariant connections on $P$ and the linear maps $\psi: \mathfrak{g} \rightarrow \mathfrak{s}$ that fulfil
a) $\psi(\vec{j})=\mathrm{d}_{e} \phi_{p}(\vec{j}) \quad \forall \vec{j} \in \mathfrak{g}_{\pi(p)}$,
b) $\psi \circ \operatorname{Ad}_{j}=\operatorname{Ad}_{\phi_{p}(j)} \circ \psi \quad \forall j \in G_{\pi(p)}$.

This bijection is explicitly given by $\omega \mapsto \Phi_{p}^{*} \omega$.

## Example 5.8.

1. Homogeneous connections. In the situation of Example 4.10 let $k=n, X=\mathbb{R}^{n}$. Then $\Phi$ is fibre transitive and for $p=(0, e)$ we have $G_{\pi(p)}=e$ and $\mathfrak{g}_{\pi(p)}=\{0\}$. Consequently, the reduced connections w.r.t. $\{p\}$ are just the linear maps $\psi: \mathbb{R}^{n} \rightarrow \mathfrak{s}$ and the corresponding homogeneous connections are given by

$$
\omega^{\psi}{ }_{(x, s)}\left(\vec{v}_{x}, \vec{\sigma}_{s}\right)=\operatorname{Ad}_{s^{-1}} \circ \psi\left(\vec{v}_{x}\right)+\mathrm{d} L_{s^{-1}}\left(\vec{\sigma}_{s}\right) \quad \forall\left(\vec{v}_{x}, \vec{\sigma}_{s}\right) \in T_{(x, s)} P
$$

2. Homogeneous isotropic connections. Let $P=\mathbb{R}^{3} \times \mathrm{SU}(2)$ and $\varrho: \mathrm{SU}(2) \rightarrow \mathrm{SO}(3)$ be the universal covering map. We consider the semi direct product $E:=\mathbb{R}^{3} \rtimes_{\varrho} \mathrm{SU}(2)$ whose multiplication is given by $(v, \sigma) \cdot \varrho\left(v^{\prime}, \sigma^{\prime}\right):=\left(v+\varrho(\sigma)\left(v^{\prime}\right), \sigma \sigma^{\prime}\right)$ for all $(v, \sigma),\left(v^{\prime}, \sigma\right) \in E$. Since $E$ equals $P$ as a set, we can define the action $\Phi$ of $E$ on $P$ just by $\varrho$. Then $E$ is a Lie group that resembles the euclidean one, and it follows from Wang's theorem that the $\Phi$-invariant connections are of the form (see e.g. Appendix A. 3 in [5])

$$
\omega_{(x, s)}^{c}\left(\vec{v}_{x}, \vec{\sigma}_{x}\right)=c \operatorname{Ad}_{s^{-1}}\left[\mu\left(\vec{v}_{x}\right)\right]+s^{-1} \vec{\sigma}_{s} \quad \forall\left(\vec{v}_{x}, \vec{\sigma}_{s}\right) \in T_{(x, s)} P
$$

Here $c$ runs over $\mathbb{R}$ and $\mu: \sum_{i=1}^{3} v^{i} \vec{e}_{i} \rightarrow \sum_{i=1}^{3} v^{i} \tau_{i}$ with matrices

$$
\tau_{1}:=\left(\begin{array}{cc}
0 & -\mathrm{i} \\
-\mathrm{i} & 0
\end{array}\right) \quad \tau_{2}:=\left(\begin{array}{cc}
0 & -1 \\
1 & 0
\end{array}\right) \quad \tau_{3}:=\left(\begin{array}{cc}
-\mathrm{i} & 0 \\
0 & \mathrm{i}
\end{array}\right)
$$

and $\left\{\vec{e}_{1}, \vec{e}_{2}, \vec{e}_{3}\right\}$ the standard basis in $\mathbb{R}^{3}$.
We close this section with a remark concerning the relations between sets of invariant connections that correspond to different lifts of the same Lie group action on the base manifold of a principal fibre bundle.

Remark 5.9. Let $P$ be a principal fibre bundle and $\Phi, \Phi^{\prime}: G \times P \rightarrow P$ be two Lie groups of automorphisms with $\varphi=\varphi^{\prime}$. Then the respective sets of invariant connections can differ significantly. In fact, in the situation of the second part of Example 5.8 let $\Phi^{\prime}((v, \sigma),(x, s)):=$ $(v+\varrho(\sigma)(x), s)$. Then $\varphi^{\prime}=\varphi$ and Appendix B. 1 shows that $\omega_{0}\left(\vec{v}_{x}, \vec{\sigma}_{s}\right):=s^{-1} \vec{\sigma}_{s}$ for $\left(\vec{v}_{x}, \vec{\sigma}_{s}\right) \in$ $T_{(x, s)} P$ is the only $\Phi^{\prime}$-invariant connection on $P$.

### 5.3 Trivial bundles - applications to LQG

In this section we determine the set of isotropic connections on $\mathbb{R}^{3} \times \mathrm{SU}(2)$ to be used for the description of isotropic gravitational systems in the framework of loop quantum gravity. To this end, we reformulate Theorem 4.9 for trivial bundles.

The spherically symmetric connections on $P=\mathbb{R}^{3} \times \mathrm{SU}(2)$ are such connections, invariant under the action $\Phi: \mathrm{SU}(2) \times P \rightarrow P,(\sigma,(x, s)) \mapsto(\sigma(x), \sigma s)$. Since $\Phi$ is not fibre transitive, we cannot use Case 5.7 for the necessary calculations. Moreover, it is not possible to apply the
results from [6] (see Example 4.6) because the $\varphi$-stabilizer of $x=0$ equals $\mathrm{SU}(2)$ whereas that of each $x \in \mathbb{R}^{3} \backslash\{0\}$ is given by the maximal torus $T_{x}:=\{\exp (t \mu(x) \mid t \in \mathbb{R})\} \subseteq \mathrm{SU}(2)$. Of course, we could ignore the origin and consider the bundle $\mathbb{R}^{3} \backslash\{0\} \times \mathrm{SU}(2)$ together with the $\Phi$-covering $\left\{\lambda \cdot \vec{e}_{1} \mid \lambda \in \mathbb{R}_{>0}\right\}$. This, however, is a different situation because an invariant connection on $\mathbb{R}^{3} \backslash\{0\} \times \mathrm{SU}(2)$ is not necessarily extendable to an invariant connection on $\mathbb{R}^{3} \times \mathrm{SU}(2)$ as the next example illustrates ${ }^{24}$.

## Example 5.10.

1. Let $S$ be a Lie group and $P=\mathbb{R}^{n} \times S$. We consider the action $\Phi: \mathbb{R}_{>0} \times P \rightarrow P$, $(\lambda,(x, s)) \mapsto(\lambda x, s)$ and claim that the only $\Phi$-invariant connection is given by

$$
\omega_{0}\left(\vec{v}_{x}, \vec{\sigma}_{s}\right):=\mathrm{d}_{s} L_{s^{-1}}\left(\vec{\sigma}_{s}\right) \quad \forall\left(\vec{v}_{x}, \vec{\sigma}_{s}\right) \in T_{(x, e)} P .
$$

In fact, $P_{\infty}:=\mathbb{R}^{n} \times\{e\}$ is a $\Phi$-covering of $P$ by Corollary 5.1 and it is straightforward to see ${ }^{25}$ that condition $i$ ) from Corollary 4.2 is equivalent to the conditions $a$ ) and $b$ ) from Remark 4.4. Let $\psi: \mathfrak{g} \times T P_{\infty}$ be a reduced connection w.r.t. $P_{\infty}$ and define $\psi_{x}:=\left.\psi\right|_{\mathfrak{g} \times T_{(x, e)}}$. Since the exponential map exp: $\mathfrak{g} \rightarrow \mathbb{R}_{>0}$ is just given by $\lambda \mapsto \mathrm{e}^{\lambda}$ for $\lambda \in \mathbb{R}=\mathfrak{g}$, we have $\widetilde{g}((x, e))=\vec{g} \cdot x \in T_{(x, e)} P_{\infty}$ for $\vec{g} \in \mathfrak{g}$. Then for $\vec{w}:=-\vec{g} \cdot x \in T_{(x, e)} P_{\infty}$ from $\left.a\right)$ we obtain

$$
\begin{equation*}
\psi_{x}(\vec{g}, \overrightarrow{0})=\psi_{x}\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{g} \cdot x\right) \quad \forall \vec{g} \in \mathfrak{g}, \quad \forall x \in \mathbb{R}^{n} \tag{5.3}
\end{equation*}
$$

In particular, $\psi_{0}(\vec{g}, \overrightarrow{0})=0$ and since $Q_{(0, e)}=\mathbb{R}_{>0} \times\{e\}$, for $q=(\lambda, e)$ condition $\left.b\right)$ yields

$$
\lambda \psi_{0}\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{w}\right)=\psi_{0}\left(\overrightarrow{0}_{\mathfrak{g}}, \lambda \vec{w}\right) \stackrel{b)}{=} \psi_{0}\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{w}\right) \quad \forall \lambda>0, \quad \forall \vec{w} \in T_{(0, e)} P_{\infty}
$$

hence $\psi_{0}=0$. Analogously, for $x \neq 0, \vec{w} \in T_{(\lambda x, e)} P_{\infty}, \lambda>0$ and $q=(\lambda, e)$, we obtain

$$
\lambda \psi_{\lambda x}\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{w}\right)=\psi_{\lambda x}\left(\overrightarrow{0}_{\mathfrak{g}}, \mathrm{d} L_{q}(\vec{w})\right) \stackrel{b)}{=} \rho(q) \circ \psi_{x}\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{w}\right)=\psi_{x}\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{w}\right)
$$

i.e., $\psi_{\lambda x}\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{w}\right)=\frac{1}{\lambda} \psi_{x}\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{w}\right)$. Here, in the second step, we have used the canonical identification of the linear spaces $T_{(x, e)} P_{\infty}$ and $T_{(\lambda x, e)} P_{\infty}$. Using the same identification, from continuity (smoothness) of $\psi$ and $\psi_{0}=0$ we obtain

$$
0=\lim _{\lambda \rightarrow 0} \psi_{\lambda x}\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{w}\right)=\lim _{\lambda \rightarrow 0} \frac{1}{\lambda} \psi_{x}\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{w}\right) \quad \forall x \in \mathbb{R}^{n}, \quad \forall \vec{w} \in T_{(x, e)} P_{\infty}
$$

so that $\psi_{x}\left(\overrightarrow{0}_{\mathfrak{g}}, \cdot\right)=0$ for all $x \in \mathbb{R}^{n}$, hence $\psi=0$ by (5.3). Finally, it is straightforward to see that $\left.\left(\Phi^{*} \omega_{0}\right)\right|_{\mathfrak{g} \times T P_{\infty}}=\psi=0$.
2. Let $P^{\prime}=\mathbb{R}^{n} \backslash\{0\} \times S$ and $\Phi$ be defined as above. Then $K \times\{e\}$, for the unit-sphere $K:=\left\{x \in \mathbb{R}^{n} \mid\|x\|=1\right\}$, is a $\Phi$-covering of $P^{\prime}$ with the properties from Example 4.6. Evaluating the corresponding conditions $i^{\prime \prime}$, $i i^{\prime \prime}$ ), $i i i^{\prime \prime}$ ) immediately shows that the set of $\Phi$-invariant connections on $P^{\prime}$ is in bijection with the smooth maps $\psi: \mathbb{R} \times T K \rightarrow \mathfrak{s}$ for which $\left.\psi\right|_{\mathbb{R} \times T_{k} K}$ is linear for all $k \in K$. The corresponding invariant connections are given by

$$
\omega_{(x, s)}^{\psi}\left(\vec{v}_{x}, \vec{\sigma}_{s}\right)=\psi\left(\frac{1}{\|x\|} \operatorname{pr}_{\|}\left(\vec{v}_{x}\right), \operatorname{pr}_{\perp}\left(\vec{v}_{x}\right)\right)+s^{-1} \vec{\sigma}_{s} \quad \forall\left(\vec{v}_{x}, \vec{\sigma}_{s}\right) \in T_{(x, s)} P^{\prime}
$$

Here $\mathrm{pr}_{\|}$denotes the projection onto the axis defined by $x \in \mathbb{R}^{n}$ and $\mathrm{pr}_{\perp}$ the projection onto the corresponding orthogonal complement in $\mathbb{R}^{n}$.

[^11]Also in the spherically symmetric case the $\varphi$-stabilizer of the origin has full dimension and it turns out to be convenient ${ }^{26}$ to use the $\Phi$-covering $\mathbb{R}^{3} \times\{e\}$ in this situation, too. Since the choice $P_{\infty}:=M \times\{e\}$ is always reasonable ${ }^{27}$ if there is a point in the base manifold $M$ (of the trivial bundle $M \times S$ ) whose stabilizer is the whole group, we now adapt Theorem 4.9 to this situation. For this, we identify $T_{x} M$ with $T_{(x, e)} P_{\infty}$ for each $x \in M$ in the sequel.

Case 5.11 (trivial principal fibre bundles). Let $(G, \Phi)$ be a Lie group of automorphisms of the trivial principal fibre bundle $P=M \times S$. Then the $\Phi$-invariant connections are in bijection with the smooth maps $\psi: \mathfrak{g} \times T M \rightarrow \mathfrak{s}$ for which $\left.\psi\right|_{\mathfrak{g} \times T_{x} M}$ is linear for all $x \in M$ and that fulfil the following properties. Let $\psi^{ \pm}\left(\vec{g}, \vec{v}_{y}, \vec{s}\right):=\psi\left(\vec{g}, \vec{v}_{y}\right) \pm \vec{s}$ for $\left((\vec{g}, \vec{s}), \vec{v}_{y}\right) \in \mathfrak{q} \times T_{y} M$. Then for $q \in Q$, $x \in M$ with $q \cdot(x, e)=(y, e) \in M \times\{e\}$ and all $\left((\vec{g}, \vec{s}), \vec{v}_{x}\right) \in \mathfrak{q} \times T_{x} M$ we have
i) $\widetilde{g}(x, e)+\vec{v}_{x}-\vec{s}=0 \quad \Longrightarrow \quad \psi^{-}\left(\vec{g}, \vec{v}_{x}, \vec{s}\right)=0$,
ii) $\psi^{+}\left(\mathrm{d} L_{q} \vec{v}_{x}\right)=\rho(q) \circ \psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{v}_{x}\right) \quad \forall \vec{v}_{x} \in T_{x} M$,
iii) $\psi\left(\operatorname{Ad}_{q}(\vec{g}), \overrightarrow{0}_{y}\right)=\rho(q) \circ \psi\left(\vec{g}, \overrightarrow{0}_{x}\right) \quad \forall \vec{g} \in \mathfrak{g}$.

Proof. The elementary proof can be found in Appendix A.
Example 5.12 (spherically symmetric systems in loop quantum gravity). Let $\varrho: \mathrm{SU}(2) \rightarrow$ $\mathrm{SO}(3)$ be the universal covering map and $\sigma(x):=\varrho(\sigma)(x)$ for $x \in \mathbb{R}^{3}$. Moreover, let $\mu: \mathbb{R} \rightarrow \mathfrak{s u}(2)$ be defined as in the second part of Example 5.8. We consider the action of $G=\mathrm{SU}(2)$ on $P=\mathbb{R}^{3} \times \mathrm{SU}(2)$ defined by $\Phi(\sigma,(x, s)):=(\varrho(\sigma)(x), \sigma s)$. It is shown in Appendix B. 2 that the corresponding invariant connections have the form

$$
\begin{align*}
\omega_{(x, s)}^{a b c}\left(\vec{v}_{x}, \vec{\sigma}_{s}\right):= & \operatorname{Ad}_{s^{-1}}\left[a(x) \mu\left(\vec{v}_{x}\right)+b(x)\left[\mu(x), \mu\left(\vec{v}_{x}\right)\right]\right.  \tag{5.4}\\
& \left.+c(x)\left[\mu(x),\left[\mu(x), \mu\left(\vec{v}_{x}\right)\right]\right]\right]+s^{-1} \vec{\sigma}_{s}
\end{align*}
$$

for $\left(\vec{v}_{x}, \vec{\sigma}_{s}\right) \in T_{(x, s)} P$ and with rotation invariant maps $a, b, c: \mathbb{R}^{3} \rightarrow \mathbb{R}$ for which the whole expression is a smooth connection. We claim that the functions $a, b, c$ can be assumed to be smooth as well. More precisely, we show that we can assume that

$$
a(x)=f\left(\|x\|^{2}\right), \quad b(x)=g\left(\|x\|^{2}\right), \quad c(x)=h\left(\|x\|^{2}\right)
$$

for smooth functions $f, g, h:(-\epsilon, \infty) \rightarrow \mathbb{R}$ with $\epsilon>0$. Then each pullback of such a spherically symmetric connection by the global section $x \mapsto(x, e)$ can be written in the form

$$
\widetilde{\omega}_{x}^{a b c}\left(\vec{v}_{x}\right)=f^{\prime}\left(\|x\|^{2}\right) \mu\left(\vec{v}_{x}\right)+g^{\prime}\left(\|x\|^{2}\right) \mu\left(x \times \vec{v}_{x}\right)+h^{\prime}\left(\|x\|^{2}\right) \mu\left(x \times\left(x \times \vec{v}_{x}\right)\right)
$$

for smooth functions $f^{\prime}, g^{\prime}, h^{\prime}:(-\epsilon, \infty) \rightarrow \mathbb{R}$ with $\epsilon>0$.
Proof of the claim. 1. Smoothness of $\omega^{a b c}$ implies smoothness of the real functions

$$
a_{\vec{n}}(\lambda):=a(\lambda \vec{n}), \quad b_{\vec{n}}(\lambda):=\lambda b(\lambda \vec{n}), \quad c_{\vec{n}}(\lambda):=\lambda^{2} c(\lambda \vec{n}) \quad \forall \lambda \in \mathbb{R}
$$

for each $\vec{n} \in \mathbb{R}^{3} \backslash\{0\}$. In fact, $a_{\vec{n}}(\lambda) \cdot \mu(\vec{n})=\omega_{(\lambda \vec{n}, e)}^{a b c}(\vec{n})$ is smooth, so that smoothness of $b_{\vec{n}}$ and $c_{\vec{n}}$ is immediate from smoothness of $\lambda \mapsto \omega_{\left(\lambda \vec{e}_{1}, e\right)}^{a b c}\left(\vec{e}_{2}\right)$.

2 . Let $\vec{n}$ be fixed. Then $a_{\vec{n}}$ is even so that $a_{\vec{n}}(\lambda)=f\left(\lambda^{2}\right)$ for a smooth function $f:\left(-\epsilon_{1}, \infty\right) \rightarrow \mathbb{R}$, see [10]. Moreover, $b_{\vec{n}}$ is smooth and odd so that $b_{\vec{n}}(\lambda)=\lambda g\left(\lambda^{2}\right)$ for some smooth function

[^12]$g:\left(-\epsilon_{2}, \infty\right) \rightarrow \mathbb{R}$, again by [10]. Similarly, $c_{\vec{n}}(\lambda)=l\left(\lambda^{2}\right)$ for a smooth function $l:\left(-\epsilon_{3}, \infty\right) \rightarrow \mathbb{R}$. Since $\lambda \mapsto l\left(\lambda^{2}\right)$ is even and $l(0)=0$, for $s \in \mathbb{N}_{>0}$ Taylor's formula yields
$$
l\left(x^{2}\right)=a_{1} x^{2}+\cdots+a_{s} x^{2 s}+x^{2(s+1)} \phi(x)=x^{2}\left(a_{1}+\cdots+a_{s} x^{2 s-2}+x^{2 s} \phi(x)\right)=x^{2} L(x)
$$
with remainder term $\phi(x):=\frac{1}{(2 s+1)!} \frac{1}{x^{2 s+2}} \int_{0}^{x}(x-t) f^{(2 s+2)}(t) \mathrm{d} t$ for $x \neq 0$ and $\phi(0):=l^{2 s+2}(0)$. Now, $\phi$ is continuous by Theorem 1 in [9] so that $L$ is continuous as well. But $x \mapsto x^{2} L(x)$ is smooth so that Corollary 1 in [9] shows that $L$ is smooth as well. Now, $L$ is even, hence $L(x)=h\left(x^{2}\right)$ for some smooth function $h:\left(-\epsilon_{4}, \infty\right) \rightarrow \mathbb{R}$. Then $c_{\vec{n}}(\lambda)=l\left(\lambda^{2}\right)=\lambda^{2} h\left(\lambda^{2}\right)$ and for $x \neq 0$ we get
\[

$$
\begin{aligned}
& b(x)=\|x\| b\left(\|x\| \frac{x}{\|x\|}\right) \frac{1}{\|x\|}=b \frac{x}{\|x\|}(\|x\|) \frac{1}{\|x\|}=g\left(\|x\|^{2}\right) \\
& c(x)=\|x\|^{2} c\left(\|x\| \frac{x}{\|x\|}\right) \frac{1}{\|x\|^{2}}=c \frac{x}{\|x\|}(\|x\|) \frac{1}{\|x\|^{2}}=h\left(\|x\|^{2}\right)
\end{aligned}
$$
\]

Moreover, for $x=0$ we have

$$
\begin{aligned}
& b(x)\left[\mu(x), \mu\left(\vec{v}_{x}\right)\right]=0=g\left(\|x\|^{2}\right)\left[\mu(x), \mu\left(\vec{v}_{x}\right)\right] \\
& c(x)\left[\mu(x),\left[\mu(x), \mu\left(\vec{v}_{x}\right)\right]\right]=0=h(x)\left[\mu(x),\left[\mu(x), \mu\left(\vec{v}_{x}\right)\right]\right]
\end{aligned}
$$

so that we can assume $a(x)=f\left(\|x\|^{2}\right), b(x)=g\left(\|x\|^{2}\right)$ and $c(x)=h\left(\|x\|^{2}\right)$ for the smooth functions $f, g, h:\left(-\min \left(\epsilon_{1}, \ldots, \epsilon_{4}\right), \infty\right) \rightarrow \mathbb{R}$.

In particular, there are spherically symmetric connections on $\mathbb{R}^{3} \backslash\{0\} \times \mathrm{SU}(2)$ which cannot be extended to those on $P$. For instance, if $b=c=0$ and $a(x):=1 /\|x\|$ for $x \in \mathbb{R}^{3} \backslash\{0\}$, then $\omega^{a b c}$ cannot be extended smoothly to an invariant connection on $\mathbb{R}^{3} \times \mathrm{SU}(2)$ since elsewise $a_{\vec{n}}$ could be extended to a continuous (smooth) function on $\mathbb{R}$.

## 6 Conclusions

We conclude with a short review of the particular cases that follow from Theorem 4.9. For this let $(G, \Phi)$ be a Lie group of automorphisms of the principal fibre bundle $(P, \pi, M, S)$ and $\varphi$ the induced action on $M$.

- If $P=M \times S$ is trivial, then $M \times\{e\}$ is a $\Phi$-covering of $P$. As we have demonstrated in the spherically symmetric and scale invariant case (cf. Examples 5.10 and 5.12 ), this choice can be useful for calculations if there is a point in $M$ whose $\varphi$-stabilizer is the whole group $G$.
- There is an element $x \in M$ which is contained in the closure of each $\varphi$-orbit. Then every $\Theta$-patch that contains some $p \in \pi^{-1}(x)$ is a $\Phi$-covering of $P$, see Example 5.5. If $\varphi$ acts transitively on $M$, then for each $p \in P$ the zero-dimensional submanifold $\{p\}$ is a $\Phi$-covering of $P$ leading to Wang's original theorem, see Case 5.7 and Example 5.8.
- Let $\Phi$ act via gauge transformations on $P$. In this case each open covering $\left\{U_{\alpha}\right\}_{\alpha \in I}$ of $M$ together with smooth sections $s_{\alpha}: U_{\alpha} \rightarrow P$ provides the $\Phi$-covering $\left\{s_{\alpha}\left(U_{\alpha}\right)\right\}_{\alpha \in I}$ of $P$. If $G$ acts trivially, this specializes to the usual description of smooth connections by means of consistent families of local 1-forms on the base manifold $M$.
- If we find a $\Theta$-patch $P_{0}$ such that $\pi\left(P_{0}\right)$ intersects each $\varphi$-orbit in a unique point, then $P_{0}$ is a $\Phi$-covering. If, in addition, the stabilizer $Q_{p}$ does not depend on $p \in P_{0}$, then we are in the situation of [6], see Example 4.6.
- Assume there is a collection of $\varphi$-orbits forming an open subset $U \subseteq M$. Then $O:=\pi^{-1}(U)$ is a principal fibre bundle and each $\Phi$-invariant connection on $P$ restricts to a $\Phi$-invariant connection on $O$. Conversely, if $U$ is in addition dense in $M$, then one can ask the question whether a $\Phi$-invariant connection on $O$ extends to a $\Phi$-invariant connection on $P$. Since such an extension is necessarily unique (continuity), $\varphi$-orbits not contained in $U$ can be seen as sources of obstructions for the extendability of invariant connections on $O$ to $P$. Indeed, the examples in Subsection 5.3 show that smoothness of these extension gives crucial restrictions. Moreover, by Example 5.5 taking one additional orbit into account can shrink the number of invariant connections to zero. Of particular interest, in this context, is the case where $G$ is compact, as then the orbits of principal type always form a dense and open subset of $M$. Moreover, due to compactness the situation of [6] always holds locally [7], giving rise to a canonical $\Phi$-covering (of $P$ and $O$ ) consisting of convenient patches. So, using the present characterization theorem there is a realistic chance to get some general classification results in the compact case ${ }^{28}$.

In the general situation one can always construct $\Phi$-coverings of $P$ from families of $\varphi$-patches in $M$ as Corollary 5.1 shows. In particular, the first three cases arise in this way.

## Appendix

## A A technical proof

Proof of Case 5.11. The only patch is $M \times\{e\}$ so that a reduced connection is a smooth $\operatorname{map} \psi: \mathfrak{g} \times T M \rightarrow \mathfrak{s}$ with the claimed linearity property and that fulfils the two conditions from Corollary 4.2. Obviously, $i i$ ) and iii) are equivalent. Moreover, i) follows from $i$ for $p_{\alpha}=p_{\beta}=(x, e), q=(e, e), \vec{w}_{p_{\beta}}=\vec{v}_{x}$ and $\vec{w}_{p_{\alpha}}=\overrightarrow{0}_{(x, e)}$, see also $\left.a\right)$ in Remark 4.4. To obtain ii) let $\vec{v}_{x} \in T_{x} M, q \in Q$ and $q \cdot(x, e)=(y, e)$. Then $\mathrm{d} L_{q} \vec{v}_{x}=\left(\vec{v}_{y},-\vec{s}\right)$ for elements $\vec{v}_{y} \in T_{y} M$ and $\vec{s} \in \mathfrak{s}$ so that

$$
\psi^{+}\left(\mathrm{d} L_{q} \vec{v}_{x}\right)=\psi^{+}\left(\vec{v}_{y},-\vec{s}\right)=\psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{v}_{y}\right)-\vec{s} \stackrel{i)}{=} \rho(q) \circ \psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{v}_{x}\right)
$$

It remains to show that i) and ii) imply $i$. To this end, let $(y, e)=q \cdot(x, e)$ for $x, y \in M$ and $q \in Q$. Then $i$ ) reads

$$
\widetilde{g}(y, e)+\vec{v}_{y}-\vec{s}=\mathrm{d} L_{q} \vec{v}_{x} \quad \Longrightarrow \quad \psi^{-}\left(\vec{g}, \vec{v}_{y}, \vec{s}\right)=\rho(q) \circ \psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{v}_{x}\right)
$$

where $\vec{v}_{x} \in T_{x} M, \vec{v}_{y} \in T_{y} M, \vec{s} \in \mathfrak{s}$ and $\vec{g} \in \mathfrak{g}$. Let $\mathrm{d} L_{q} \vec{v}_{x}=\left(\vec{v}_{y},-\vec{s}\right)$ be as above. If ii) is true, then it is clear from

$$
\psi^{-}\left(\vec{v}_{y}, \vec{s}\right)=\psi^{+}\left(\mathrm{d} L_{q} \vec{v}_{x}\right) \stackrel{\mathrm{i})}{=} \rho(q) \circ \psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{v}_{x}\right)
$$

that $i)$ is true for $\left(\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{s}\right), \vec{v}_{y}\right)$, i.e.,

$$
\overrightarrow{0}_{\mathfrak{g}}+\vec{v}_{y}-\vec{s}=\mathrm{d} L_{q} \vec{v}_{x} \quad \Longrightarrow \quad \psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{v}_{y}\right)-\vec{s}=\rho(q) \circ \psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{v}_{x}\right)
$$

Due to i) and the linearity properties of $\psi$, the condition $i$ ) then is also true for each other element $\left(\left(\vec{g}^{\prime}, \vec{s}^{\prime}\right), \vec{v}_{y}^{\prime}\right) \in \mathfrak{q} \times T_{y} M$ with $\widetilde{g}^{\prime}(y, e)+\vec{v}_{y}^{\prime}-\vec{s}^{\prime}=\mathrm{d} L_{q} \vec{v}_{x}$.

[^13]
## B Technical calculations

Let $P=\mathbb{R}^{3} \times \mathrm{SU}(2), \varrho: \mathrm{SU}(2) \rightarrow \mathrm{SO}(3)$ the universal covering map, $E=\mathbb{R}^{3} \rtimes_{\varrho} \mathrm{SU}(2)$ and $\mu: \mathbb{R}^{3} \rightarrow \mathfrak{s u}(2)$ be defined as in the second part of Example 5.8. Then $\varrho(\sigma)=\mu^{-1} \circ \mathrm{Ad}_{\sigma} \circ \mu$ and each $\sigma \in \mathrm{SU}(2)$ can be written as

$$
\sigma=\cos (\alpha / 2) \mathbb{1}+\sin (\alpha / 2) \mu(\vec{n})=\exp (\alpha / 2 \cdot \mu(\vec{n}))
$$

for some $|\vec{n}|=1$ and $\alpha \in[0,2 \pi]$. In this case $\varrho(\sigma)$ rotates a point $x$ by the angle $\alpha$ w.r.t. the axis $\vec{n}$. For simplicity, if $\sigma \in \mathrm{SU}(2)$ and $x \in \mathbb{R}^{3}$, we write $\sigma(x)$ instead of $\varrho(\sigma)(x)$ in the following.

## B. 1 A result used in the end of Section 5

We consider the fibre transitive action $\Phi^{\prime}: E \times P \rightarrow P$ that is defined by $\Phi^{\prime}((v, \sigma),(x, s)):=$ $(v+\sigma(x), s)$ and claim that the connection

$$
\omega_{0}\left(\vec{v}_{x}, \vec{\sigma}_{s}\right)=s^{-1} \vec{\sigma}_{s} \quad \forall\left(\vec{v}_{x}, \vec{\sigma}_{s}\right) \in T_{(x, s)} P
$$

is the only $\Phi^{\prime}$-invariant one. For this observe that the stabilizer of $x=0$ w.r.t. $\varphi^{\prime}$ is given by $\mathrm{SU}(2)$ and $\phi_{(0, e)}^{\prime}(\sigma)=e$ for all $\sigma \in \mathrm{SU}(2)$. We apply Wang's theorem to $p=(0, e)$. Then condition a) yields $\psi(\vec{s})=0$ for all $\vec{s} \in \mathfrak{s u}(2)$ and b) now reads $\psi \circ \mathrm{Ad}_{\sigma}=\psi$ for all $\sigma \in \mathrm{SU}(2)$. Consequently, for $\vec{v} \in \mathbb{R}^{3} \subseteq \mathfrak{e}=\mathbb{R}^{3} \times \mathfrak{s u}(2)$ we obtain

$$
0=\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} \psi(\vec{v})=\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} \psi \circ \operatorname{Ad}_{\exp (t \vec{s})}(\vec{v})=\psi\left(\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} \varrho(\exp (t \vec{s}))(\vec{v})\right)=\psi \circ \mu^{-1}([\vec{s}, \mu(\vec{v})])
$$

for all $\vec{s} \in \mathfrak{s u}(2)$ just by linearity of $\psi$. This gives

$$
0=\psi\left(\mu^{-1}\left(\left[\tau_{i}, \mu\left(\vec{e}_{j}\right)\right]\right)\right)=\psi\left(\mu^{-1}\left(\left[\tau_{i}, \tau_{j}\right]\right)\right)=2 \epsilon_{i j k} \psi\left(\vec{e}_{k}\right)
$$

hence $\psi=0=\Phi_{p}^{*} \omega_{0}$.

## B. 2 Spherically symmetric connections

We consider the action $\Phi$ of $\mathrm{SU}(2)$ on $P$ defined by $\Phi(\sigma,(x, s)):=(\sigma(x), \sigma s)$ and show that the corresponding invariant connections are given by (see (5.4) in Example 5.12)

$$
\omega_{(x, s)}^{a b c}\left(\vec{v}_{x}, \vec{\sigma}_{s}\right):=\operatorname{Ad}_{s^{-1}}\left[a(x) \mu\left(\vec{v}_{x}\right)+b(x)\left[\mu(x), \mu\left(\vec{v}_{x}\right)\right]+c(x)\left[\mu(x),\left[\mu(x), \mu\left(\vec{v}_{x}\right)\right]\right]\right]+s^{-1} \vec{\sigma}_{s}
$$

with rotation invariant maps $a, b, c: \mathbb{R}^{3} \rightarrow \mathbb{R}$ for which the whole expression is a smooth connection. Now, a straightforward calculation shows that each $\omega^{a b c}$ is $\Phi$-invariant so that it remains to verify that each $\Phi$-invariant connection is of the upper form. To this end, we reduce the connections $\omega^{a b c}$ w.r.t. $P_{\infty}=\mathbb{R}^{3} \times\{e\}$ and show that each map $\psi$ as in Case 5.11 can be obtained in this way. For this, let $\vec{g} \in \mathfrak{g}, p=(x, e) \in P_{\infty}$ and $\gamma_{x}:(-\epsilon, \epsilon) \rightarrow M$ be a smooth curve with $\dot{\gamma}_{x}(0)=\vec{v}_{x} \in T_{x} M \subseteq T_{p} P_{\infty}$. Then

$$
\begin{align*}
\mathrm{d}_{(e, p)} \Phi\left(\vec{g}, \vec{v}_{x}\right) & =\left(\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} \mu^{-1}\left(\exp (t \vec{g}) \mu\left(\gamma_{x}(t)\right) \exp (t \vec{g})^{-1}\right), \exp (t \vec{g})\right)  \tag{B.1}\\
& =\left(\mu^{-1}([\vec{g}, \mu(x)])+\vec{v}_{x}, \vec{g}\right)
\end{align*}
$$

This equals $\vec{s}$ iff $\vec{g}=\vec{s}$ and $\vec{v}_{x}=\mu^{-1}([\mu(x), \vec{g}])$. Consequently, for the reduced connection $\psi^{a b c}$ that corresponds to $\omega^{a b c}$ we obtain

$$
\begin{aligned}
\psi^{a b c}\left(\vec{g}, \vec{v}_{x}\right)= & \left(\Phi^{*} \omega^{a b c}\right)_{(e, p)}\left(\vec{g}, \vec{v}_{x}\right)=\omega_{p}^{a b c}\left(\mu^{-1}\left([\vec{g}, \mu(x)]+\mu\left(\vec{v}_{x}\right)\right), \vec{g}\right) \\
= & a(x)\left[[\vec{g}, \mu(x)]+\mu\left(\vec{v}_{x}\right)\right]+b(x)\left[[\mu(x),[\vec{g}, \mu(x)]]+\left[\mu(x), \mu\left(\vec{v}_{x}\right)\right]\right] \\
& +c(x)\left[[\mu(x),[\mu(x),[\vec{g}, \mu(x)]]]+\left[\mu(x),\left[\mu(x), \mu\left(\vec{v}_{x}\right)\right]\right]\right]+\vec{g}
\end{aligned}
$$

Now, assume that $\psi$ is as in Case 5.11. Then for $q \in Q$ and $p \in P_{\infty}$ we have $q \cdot p \in P_{\infty}$ iff $q=(\sigma, \sigma)$ for some $\sigma \in \mathrm{SU}(2)$ and $p=(x, e)$ for some $x \in M$. Consequently, $q \cdot p=(\sigma(x), e)$ as well as $\mathrm{d} L_{q}\left(\vec{v}_{x}\right)=\sigma\left(\vec{v}_{x}\right)$ for all $\vec{v}_{x} \in T_{x} M$ so that ii) gives

$$
\psi\left(\overrightarrow{0}_{\mathfrak{g}}, \sigma\left(\vec{v}_{x}\right)\right)=\psi^{+}\left(\mathrm{d} L_{q}\left(\vec{v}_{x}\right)\right)=\operatorname{Ad}_{\sigma} \circ \psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{v}_{x}\right),
$$

hence

$$
\begin{equation*}
\psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{v}_{x}\right)=\operatorname{Ad}_{\sigma^{-1}} \circ \psi\left(\overrightarrow{0}_{\mathfrak{g}}, \sigma\left(\vec{v}_{x}\right)\right) \quad \forall \vec{v}_{x} \in T_{x} M . \tag{B.2}
\end{equation*}
$$

If $x \neq 0$, then for $\sigma_{t}:=\exp (t \mu(x))$ we have $\sigma_{t}(x)=x$ and $\sigma_{t}\left(\vec{v}_{x}\right) \in T_{(x, e)} P_{\infty}$ for all $t \in \mathbb{R}$. Then linearity of $\psi_{x}:=\left.\psi\right|_{\mathfrak{g} \times T_{(x, e)} P_{\infty}}$ yields

$$
\begin{aligned}
0 & =\left.\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} \psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{v}_{x}\right) \stackrel{(\mathrm{B} .2)}{=} \frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} \mathrm{Ad}_{\sigma_{t}^{-1}} \circ \psi\left(\overrightarrow{0}_{\mathfrak{g}}, \sigma_{t}\left(\vec{v}_{x}\right)\right) \\
& =\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} \sigma_{t}^{-1}\left(\psi \circ \mu^{-1}\right)\left(\sigma_{t} \mu\left(\vec{v}_{x}\right) \sigma_{t}^{-1}\right) \sigma_{t} \\
& \text { lin. } \\
= & \mu(x) \psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{v}_{x}\right)+\left(\psi_{x} \circ \mu^{-1}\right)\left[\mu(x) \mu\left(\vec{v}_{x}\right)-\mu\left(\vec{v}_{x}\right) \mu(x)\right]+\psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{v}_{x}\right) \mu(x),
\end{aligned}
$$

hence $\left[\mu(x), \psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{v}_{x}\right)\right]=\left(\psi_{x} \circ \mu^{-1}\right)\left(\left[\mu(x), \mu\left(\vec{v}_{x}\right)\right]\right)$. For $x=\lambda \vec{e}_{1} \neq 0$ and $\kappa_{j}:=\psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{v}_{x}\right)$ with $\vec{v}_{x}=\vec{e}_{j}$ this reads

$$
\left[\tau_{1}, \kappa_{j}\right]=\left(\psi_{x} \circ \mu^{-1}\right)\left(\left[\tau_{1}, \tau_{j}\right]\right)=\left(\psi_{x} \circ \mu^{-1}\right)\left(2 \epsilon_{1 j k} \tau_{k}\right)=2 \epsilon_{1 j k} \psi_{x}\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{e}_{k}\right)=2 \epsilon_{1 j k} \kappa_{k} .
$$

From these relations it follows that

$$
\kappa_{1}=r(\lambda) \tau_{1}, \quad \kappa_{2}=s(\lambda) \tau_{2}+t(\lambda) \tau_{3}, \quad \kappa_{3}=s(\lambda) \tau_{3}-t(\lambda) \tau_{2}
$$

for real constants $r(\lambda), s(\lambda), t(\lambda)$ depending on $\lambda \in \mathbb{R} \backslash\{0\}$. Then for $x=\lambda \vec{e}_{1}$ and

$$
a\left(\lambda \vec{e}_{1}\right):=r(\lambda), \quad b\left(\lambda \vec{e}_{1}\right):=\frac{t(\lambda)}{2 \lambda}, \quad c\left(\lambda \vec{e}_{1}\right):=\frac{r(\lambda)-s(\lambda)}{4 \lambda^{2}}
$$

linearity of $\psi_{x}$ yields that

$$
\psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{v}_{x}\right)=a(x) \mu\left(\vec{v}_{x}\right)+b(x)\left[\mu(x), \mu\left(\vec{v}_{x}\right)\right]+c(x)\left[\mu(x),\left[\mu(x), \mu\left(\vec{v}_{x}\right)\right]\right] .
$$

Now, if $x \neq 0$ is arbitrary, then $x=\sigma\left(\lambda \vec{e}_{1}\right)$ for some $\sigma \in \mathrm{SU}(2)$ and $\lambda>0$. So, $(\sigma, \sigma) \cdot\left(\lambda \vec{e}_{1}, e\right)=$ $(x, e)$ and if we consider $\sigma^{-1}\left(\vec{v}_{x}\right)$ as an element of $T_{\left(\lambda \overrightarrow{e_{1}}, e\right)} P_{\infty}$, then ii) yields

$$
\begin{aligned}
\psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{v}_{x}\right) & =\psi^{+}\left(\vec{v}_{x}\right)=\psi^{+}\left(\mathrm{d} L_{(\sigma, \sigma)}\left(\sigma^{-1}\left(\vec{v}_{x}\right)\right)\right) \\
& \stackrel{\mathrm{ii})}{=} \operatorname{Ad}_{\sigma} \circ \psi^{+}\left(\sigma^{-1}\left(\vec{v}_{x}\right)\right)=\operatorname{Ad}_{\sigma} \circ \psi\left(\overrightarrow{0}_{\mathfrak{g}}, \sigma^{-1}\left(\vec{v}_{x}\right)\right) \\
& =a\left(\lambda \vec{e}_{1}\right) \mu\left(\vec{v}_{x}\right)+b\left(\lambda \vec{e}_{1}\right)\left[\mu(x), \mu\left(\vec{v}_{x}\right)\right]+c\left(\lambda \vec{e}_{1}\right)\left[\mu(x),\left[\mu(x), \mu\left(\vec{v}_{x}\right)\right]\right] .
\end{aligned}
$$

For $x=0$ we have $\sigma(x)=x$ for all $\sigma \in \mathrm{SU}(2)$ and analogous to the case $x \neq 0$ but now for $\sigma_{t}:=\exp (t \vec{g})$ with $\vec{g} \in \mathfrak{g}$ we obtain from (B.2) that

$$
\left[\vec{g}, \psi_{0}\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{v}_{0}\right)\right]=\left(\psi_{0} \circ \mu^{-1}\right)\left(\left[\vec{g}, \mu\left(\vec{v}_{0}\right)\right]\right) \quad \forall \vec{g} \in \mathfrak{s u}(2), \quad \forall \vec{v}_{0} \in T_{0} M
$$

This gives $\left[\tau_{i}, \psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{e}_{j}\right)\right]=2 \epsilon_{i j k} \psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{e}_{k}\right)$ and forces $\psi\left(\vec{v}_{0}\right)=a(0) \mu\left(\vec{v}_{0}\right)$ for all $\vec{v}_{0} \in T_{(0, e)} P_{\infty}$ where $a(0) \in \mathbb{R}$ is some constant. Together, this shows

$$
\psi\left(\overrightarrow{0}_{\mathfrak{g}}, \vec{v}_{x}\right)=a(x) \mu\left(\vec{v}_{x}\right)+b(x)\left[\mu(x), \mu\left(\vec{v}_{x}\right)\right]+c(x)\left[\mu(x),\left[\mu(x), \mu\left(\vec{v}_{x}\right)\right]\right]
$$

with functions $a, b, c$ that depend on $\|x\|$ in such a way that the whole expression is smooth. Finally, to determine $\psi\left(\vec{g}, \overrightarrow{0}_{x}\right)$ for $\vec{g} \in \mathfrak{s u}(2)=\mathfrak{g}$, we consider $\mu^{-1}([\mu(x), \vec{g}])$ as an element of $T_{(x, e)} P_{\infty}$. Then by (B.1) we obtain from i) that $\psi\left(\vec{g}, \mu^{-1}([\mu(x), \vec{g}])\right)-\vec{g}=0$, hence

$$
\begin{aligned}
\psi\left(\vec{g}, \overrightarrow{0}_{x}\right) & =\vec{g}-\psi\left(\overrightarrow{0}_{\mathfrak{g}}, \mu^{-1}([\mu(x), \vec{g}])\right) \\
& =\vec{g}-a(x)[\mu(x), \vec{g}]-b(x)[\mu(x),[\mu(x), \vec{g}]]-c(x)[\mu(x),[\mu(x),[\mu(x), \vec{g}]]] \\
& =a(x)[\vec{g}, \mu(x)]+b(x)[\mu(x),[\vec{g}, \mu(x)]]+c(x)[\mu(x),[\mu(x),[\mu(x), \vec{g}]]]+\vec{g} .
\end{aligned}
$$

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[^0]:    ${ }^{1}$ Each Lie group of automorphisms of a bundle induces a smooth action on the base manifold.
    ${ }^{2}$ Amongst others, they assume the $\varphi$-stabilizer of $\pi\left(p_{0}\right)$ to be the same for all $p_{0} \in P_{0}$.
    ${ }^{3}$ At the moment assume that $P_{\alpha} \subseteq P$ is a subset which, at the same time, is a manifold such that the inclusion $\operatorname{map} \iota_{\alpha}: P_{\alpha} \rightarrow P$ is an immersion. Here we tacitly identify $T_{p_{\alpha}} P_{\alpha}$ with $\operatorname{im}\left[\mathrm{d}_{p_{\alpha}} \iota_{\alpha}\right]$. Note that we do not require $P_{\alpha}$ to be a topological submanifold of $P$. For details see Convention 3.1.

[^1]:    ${ }^{4}$ Here pointwise means to consider such elements $q \in G \times S$ that are contained in the $\Theta$-stabilizer of some fixed $p_{\alpha} \in P_{\alpha}$ for $\alpha \in I$.
    ${ }^{5}$ See, e.g., calculations in Appendix B.2.

[^2]:    ${ }^{6}$ Here $\mathrm{pr}_{2}$ denotes the projection onto the second factor.

[^3]:    ${ }^{7}$ In literature sometimes the latter condition is used to define $\Phi$-invariance of connections.

[^4]:    ${ }^{8}$ The sum is not necessarily direct.
    ${ }^{9}$ In fact, let $V \subseteq \mathrm{~d}_{e} \Psi_{x}(\mathfrak{g})$ be an algebraic complement of $T_{x} N$ in $T_{x} M$ and $V^{\prime} \subseteq \mathfrak{g}$ a linear subspace with $\operatorname{dim}\left[V^{\prime}\right]=\operatorname{dim}[V]$ and $\mathrm{d}_{e} \Psi_{x}\left(V^{\prime}\right)=V$. Then we find a submanifold $H$ of $G$ through $e$ with $T_{e} H=V^{\prime}$ so that $\mathrm{d}_{(e, x)} \Psi: T_{e} H \times T_{x} N \rightarrow T_{x} M$ is bijective.
    ${ }^{10}$ Recall that $\left.\mathrm{d}_{(e, x)} \Psi\right|_{T_{e} H^{\prime} \times T_{e} N^{\prime}}:\left(\vec{h}, \vec{v}_{x}\right) \mapsto \mathrm{d}_{(e, x)} \Psi\left(\mathrm{d}_{e} \tau_{H}(\vec{h}), \mathrm{d}_{x} \tau_{N}\left(\vec{v}_{x}\right)\right)$.

[^5]:    ${ }^{11}$ More precisely, $\tau_{P_{\beta}}\left(p_{\beta}\right)=q \cdot \tau_{P_{\alpha}}\left(p_{\alpha}\right)$ by Convention 3.1.
    ${ }^{12}$ See end of Subsection 2.2.

[^6]:    ${ }^{13}$ Observe that due to surjectivity of $d_{\left(e, p_{\beta}\right)} \Phi$ such elements always exist.
    ${ }^{14}$ Recall equation (4.1).

[^7]:    ${ }^{15}$ Cf. Lemma 3.7.1.
    ${ }^{16}$ In the sense that $\left.\psi\right|_{\mathfrak{g} \times T_{p} P_{0}}$ is linear for all $p \in P_{0}$.

[^8]:    ${ }^{17} \operatorname{Then} \operatorname{im}\left[\mathrm{~d}_{0} s_{0}\right]+\operatorname{im}\left[\mathrm{d}_{e} \Phi_{s_{0}(x)}\right]+T v_{s_{0}(0)} P=\operatorname{span}_{\mathbb{R}}\left(\vec{e}_{1}\right) \oplus T v_{(0, e)} P \neq \mathbb{R}^{2} \oplus T v_{(0, e)} P=T_{(0, e)} P$ so that $\left(U_{0}, s_{0}\right)$ is not a $\Theta$-patch as it fails the transversality condition (3.2) from Remark 3.6.2.
    ${ }^{18}$ These are trivial in this case because $H=\{e\}$.

[^9]:    ${ }^{19}$ It is always possible to choose $I=M$.
    ${ }^{20}$ This is that $\pi \circ s_{\alpha}=\mathrm{id}_{M_{\alpha}}$.
    ${ }^{21}$ Observe that $\delta_{\alpha \beta}(g, x)=\phi_{s_{\alpha}(x)}^{-1}(g) \cdot \delta_{\alpha \beta}(e, x)$.
    ${ }^{22} \mathrm{~d}_{e} \Phi_{p}(\vec{g})-\widetilde{s}(p)=0$ iff $(\vec{g}, \vec{s}) \in \mathfrak{q}_{p}$ iff $\vec{s}=\mathrm{d}_{e} \phi_{p}(\vec{g})$.

[^10]:    ${ }^{23}\left(\psi_{\alpha} \circ \mathrm{d}_{x} s_{\alpha}\right)\left(\vec{v}_{x}\right):=\psi_{\alpha}\left(\overrightarrow{0}_{\mathfrak{g}}, \mathrm{d}_{x} s_{\alpha}\left(\vec{v}_{x}\right)\right)$ for $\vec{v}_{x} \in T_{x} U_{\alpha \beta}$.

[^11]:    ${ }^{24}$ See also remarks following Example 5.12.
    ${ }^{25}$ Cf. Remark 4.4.3.

[^12]:    ${ }^{26}$ Cf. Appendix B.2.
    ${ }^{27}$ Cf. Lemma 3.4.1.

[^13]:    ${ }^{28} \mathrm{To}$ be used, e.g., to extend the framework of the foundational LQG reduction paper [2].

