

EXPANSIONS OF THE KUREPA FUNCTION

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Dedicated to the memory of Professor Đ. Kurepa

Abstract. The Taylor series expansions of the Kurepa function $K(a+z)$, $a \geq 0$, and numerical determination of their coefficients $b_\nu(a)$ for $a=0$ and $a=1$ are given. An asymptotic behaviour of $b_\nu(a)$ as well as that $|b_\nu(a)/b_{\nu+1}(a)| \sim a+1$, when $\nu \rightarrow \infty$, are shown. Using this fact, a transformation of series with much faster convergence is done. Numerical values of coefficients in such a transformed series for $a=0$ and $a=1$ are given with 30 decimal digits. Also, the Chebyshev expansions of $K(1+z)$ and $1/K(1+z)$ are obtained.

1. Introduction

In 1971 Professor Đ. Kurepa (see [8–9]) defined so-called left factorial as

$$!n = 0! + 1! + \cdots + (n-1)!$$

and extended it to the complex plane

$$K(z) = \int_0^\infty \frac{t^z - 1}{t - 1} e^{-t} dt \quad (\operatorname{Re} z > 0). \quad (1.1)$$

This function can be extended analytically to the whole complex plane by

$$K(z) = K(z+1) - \Gamma(z+1), \quad (1.2)$$

where $\Gamma(z)$ is the gamma function defined by

$$\Gamma(z) = \int_0^\infty e^{-t} t^{z-1} dt \quad (\operatorname{Re} z > 0) \quad \text{and} \quad z\Gamma(z) = \Gamma(z+1).$$

Kurepa [9] proved that $K(z)$ is a meromorphic function with simple poles at the points $z_k = -k$ ($k \in \mathbb{N} \setminus \{2\}$). Slavić [17] found the representation

$$K(z) = -\frac{\pi}{e} \cot \pi z + \frac{1}{e} \left(\sum_{n=1}^{\infty} \frac{1}{n!n} + \gamma \right) + \sum_{n=0}^{\infty} \Gamma(z-n),$$

where γ is Euler's constant. These formulas were mentioned also in the book [11]. A number of problems and hypotheses, especially in number theory, were posed by Kurepa and then considered by several mathematicians. For details and a complete list of references see a recent survey written by Ivić and Mijajlović [6].

Evaluation of $K(z)$ for some specific z in $(0, 1)$, using quadrature formulas with relatively small accuracy, was done by Slavić and the author of this paper (see [9]).

In this paper we give power series expansions of the Kurepa function $K(a+z)$, $a \geq 0$, and determine numerical values of their coefficients $b_\nu(a)$ for $a = 0$ and $a = 1$. Using an asymptotic behaviour of $b_\nu(a)$, when $\nu \rightarrow \infty$, we give a transformation of series with much faster convergence. Also, we obtain the Chebyshev expansions for $K(1+z)$ and $1/K(1+z)$. For similar expansions of the gamma function see e.g. Davis [3], Luke [10], Fransén and Wrigge [4], and Bohman and Fröberg [1].

2. Power series expansions

There are many investigations about the gamma function $\Gamma(z)$, including numerical calculations. Some old references can be found in Luke [10]. In 1980 Fransén and Wrigge [4] determined numerical values of the coefficients in the Taylor series expansion

$$\Gamma(a+x)^m = \sum_{k=0}^{\infty} g_k(m, a) x^k$$

for certain values of m and a and used these values to calculate $\Gamma(p/q)$ ($p, q = 1(1)10$; $p < q$) with a high precision. Recently, Bohman and Fröberg [1] have given power series expansions of the form

$$\Gamma(n+1+z) = n!(1 + d_1 z + d_2 z^2 + \dots) \quad (n = 2, 3, 4, 10)$$

and

$$(-1)^n n! \Gamma(-n+z) = \frac{n}{1-z} - \frac{1}{(n+1)(1+z)} + \frac{1}{z} (1 + f_1 z + f_2 z^2 + \dots)$$

for $n = 0, 1, 2, 10$.

In this section we investigate the Kurepa function in a similar way.

Let $a \geq 0$. Differentiating (1.1) we find

$$K^{(\nu)}(a) = \int_0^{\infty} \frac{t^a \log^\nu t}{t-1} e^{-t} dt \quad (\nu \geq 1), \quad (2.1)$$

so that the Taylor series expansion is given by

$$K(a+z) = b_0(a) + b_1(a)z + b_2(a)z^2 + b_3(a)z^3 + \dots, \quad (2.2)$$

where

$$b_0(a) = K(a), \quad b_\nu(a) = \frac{1}{\nu!} K^{(\nu)}(a) \quad (\nu \geq 1). \quad (2.3)$$

Taking a to be an integer, we have that $b_0(a) = !a$. For example, $b_0(0) = 0$, $b_0(1) = 1$, $b_0(2) = 2$, $b_0(3) = 4$, etc.

The following theorem gives an asymptotic behaviour of $b_\nu(a)$.

THEOREM 2.1. *For every $a > 0$, $b_\nu(a) \rightarrow 0$, when $\nu \rightarrow \infty$. If $a = 0$ we have*

$$\lim_{\nu \rightarrow \infty} b_{2\nu}(0) = -1, \quad \lim_{\nu \rightarrow \infty} b_{2\nu+1}(0) = 1.$$

Precisely,

$$b_\nu(a) \sim \frac{(-1)^{\nu+1}}{(a+1)^{\nu+1}}, \quad \text{when } \nu \rightarrow \infty.$$

Proof. We start with the identity

$$K(a) = R_1(a) + R_2(a), \quad (2.4)$$

where

$$R_1(a) = \int_0^1 \frac{t^a - 1}{t - 1} e^{-t} dt \quad \text{and} \quad \int_1^\infty \frac{t^a - 1}{t - 1} e^{-t} dt.$$

Differentiating (2.4) ν times with respect to a , we obtain

$$b_\nu(a) = \frac{K^{(\nu)}(a)}{\nu!} = -\frac{1}{\nu!} \int_0^1 \frac{t^a \log^\nu t}{1-t} \sum_{k=0}^\infty \frac{(-1)^k t^k}{k!} dt + \frac{1}{\nu!} R_2^{(\nu)}(a), \quad (2.5)$$

i.e.,

$$b_\nu(a) = -\frac{1}{\nu!} \sum_{k=0}^\infty \frac{(-1)^k}{k!} \int_0^1 \frac{t^{k+a} \log^\nu t}{1-t} dt + \frac{1}{\nu!} R_2^{(\nu)}(a). \quad (2.6)$$

Using the Karamata inequality [7] (cf. [14, p. 272]),

$$\frac{\log t}{t-1} \leq \frac{1}{\sqrt{t}} \quad (t \geq 1),$$

for the last term on the right in (2.6), we have

$$\begin{aligned} 0 \leq \frac{1}{\nu!} R_2^{(\nu)}(a) &= \frac{1}{\nu!} \int_1^\infty \frac{t^a \log^\nu t}{t-1} e^{-t} dt \\ &\leq \frac{1}{\nu!} \int_1^\infty \frac{(t-1)^{\nu-1}}{t^{\nu/2-a}} e^{-t} dt \\ &\leq \frac{1}{\nu! e} \int_0^\infty \frac{z^{\nu-1}}{(z+1)^{\nu/2-a}} e^{-z} dz, \end{aligned}$$

i.e.,

$$\frac{1}{\nu!} R_2^{(\nu)}(a) < \frac{1}{\nu!e} \int_0^\infty z^{\nu/2+a-1} e^{-z} dz = \frac{1}{\nu!e} \Gamma\left(\frac{\nu}{2} + a\right) \rightarrow 0,$$

when $\nu \rightarrow \infty$.

On the other side, we have that (cf. [16, p. 491])

$$\int_0^1 \frac{t^{k+a} \log^\nu t}{1-t} dt = (-1)^\nu \nu! \zeta(\nu+1, k+a+1),$$

where $\zeta(z, \alpha) = \sum_{m=0}^\infty (\alpha+m)^{-z}$ is the generalized Riemman function (the Hurwitz Zeta function).

Therefore,

$$b_r(a) = (-1)^{\nu+1} \sum_{k=0}^\infty \frac{(-1)^k}{k!} \zeta(\nu+1, k+a+1) + \frac{1}{\nu!} R_2^{(\nu)}(a). \quad (2.7)$$

When $\nu \rightarrow \infty$, we see that

$$\zeta(\nu+1, q+1) = \frac{1}{(q+1)^{\nu+1}} + \sum_{m=1}^\infty \frac{1}{(q+1+m)^{\nu+1}} \rightarrow \begin{cases} 0, & \text{if } q > 0, \\ 1, & \text{if } q = 0. \end{cases}$$

Thus, when $\nu \rightarrow \infty$ all terms in the right-hand side of (2.7) approaches zero except the term $(-1)^{\nu+1}$ in the case when $a = 0$ ($k = 0$, $q = k + a = 0$). The last statement of the theorem follows immediately from (2.7). \square

From this theorem we obtain:

COROLLARY 2.2. *For the coefficients $b_\nu(0)$ in the series (2.2) we have that $b_\nu(0) = (-1)^{\nu+1}(1 + \varepsilon_\nu)$, where $\lim_{\nu \rightarrow \infty} \varepsilon_\nu = 0$.*

Remark 2.1. We note that the radius of convergence of the series (2.2) increases as a increases, because of $|b_\nu(a)/b_{\nu+1}(a)| \sim a + 1$, when $\nu \rightarrow \infty$.

The following theorem gives an expansion with a faster convergence.

THEOREM 2.3. *For $|z| < 1$ we have*

$$K(z) = \frac{1}{1+z} \sum_{\nu=1}^\infty \alpha_\nu z^\nu, \quad (2.8)$$

where $\alpha_\nu = b_\nu(0) + b_{\nu-1}(0) = (-1)^{\nu-1} \Delta \varepsilon_{\nu-1}$, $\nu \geq 1$. The coefficients $b_\nu(0)$ and ε_ν are given in Corollary 2.2.

Proof. Applying the Euler-Abel transformation (cf. Milovanović [12, pp. 49–50]) to (2.2), with $a = 0$, we obtain (2.8). \square

Remark 2.2. The series (2.8) converges also for $z = 1$. Notice that $\sum_{\nu=1}^\infty \alpha_\nu = 2$.

Similar to Theorem 2.3 we can state and prove the following more general result:

THEOREM 2.4. For $|z| < a + 1$ we have

$$K(a + z) = \frac{1}{a + 1 + z} \sum_{\nu=0}^{\infty} \beta_{\nu}(a) z^{\nu}, \quad (2.9)$$

where

$$\beta_0(a) = (a + 1)b_0(a), \quad \beta_{\nu}(a) = (a + 1)b_{\nu+1}(a) + b_{\nu}(a), \quad \nu \geq 1. \quad (2.10)$$

Proof. Put $K(a + z) = b_0(a) + (z/(a + 1))g(z)$, where

$$g(z) = \sum_{\nu=1}^{\infty} (a + 1)^{\nu} b_{\nu}(a) \left(\frac{z}{a + 1}\right)^{\nu-1}.$$

Since

$$\left(1 + \frac{z}{a + 1}\right)g(z) = \sum_{\nu=0}^{\infty} \{(a + 1)b_{\nu+1}(a) + b_{\nu}(a)\} z^{\nu} - b_0(a),$$

we obtain (2.9), where the coefficients $\beta_{\nu}(a)$ are given by (2.10). This is, in fact, the Euler-Abel transformation of the series (2.2). \square

Remark 2.3. Suppose that $a \geq 1$. For $z = 1$ and $z = -1$, (2.9) gives

$$\sum_{\nu=0}^{\infty} \beta_{\nu}(a) = (a + 2)K(a + 1) \quad \text{and} \quad \sum_{\nu=0}^{\infty} (-1)^{\nu} \beta_{\nu}(a) = aK(a - 1),$$

respectively. If $a = 1$, we have that

$$\sum_{\nu=0}^{\infty} \beta_{\nu}(1) = 6 \quad \text{and} \quad \sum_{\nu=0}^{\infty} (-1)^{\nu} \beta_{\nu}(1) = 0.$$

Numerical calculation of $b_{\nu}(a)$ for $\nu \geq 1$ is not so easy, because of singular integrand in (2.1). Some Gaussian type of quadrature formulas, e.g. Gauss-Laguerre formulas or, after changing variable, Gauss-Einstein formulas on $(0, \infty)$ (cf. Gautschi and Milovanović [5]) show a slowly convergence. Using the fact that for an integral of an analytic function over $(-\infty, \infty)$ the trapezoidal rule with an equal mesh size is asymptotically optimal among formulas with the same density of sampling points (see Takahasi and Mori [18]), we will transform our integrals to ones over $(-\infty, \infty)$, taking $t = \exp(\pi \sinh x)$, and then apply the trapezoidal quadrature formula

$$I = \int_{-\infty}^{\infty} g(x) dx \approx I_h = h \sum_{n=-\infty}^{\infty} g(nh) \quad (2.11)$$

For a such kind of quadratures obtained by variable transformation, especially for the DE-rule, see a survey written by Mori [15].

TABLE 2.1: The coefficients in the series (2.8) and (2.9) when $a = 1$

ν	$\beta_\nu(0)$						$\beta_\nu(1)$					
0	0.00000	00000	00000	00000	00000	00000	2.00000	00000	00000	00000	00000	00000
1	1.43220	57346	53224	41481	10310	06215	2.70998	01395	03383	10840	90378	32265
2	0.46888	45450	62632	23498	89857	40662	0.90645	96812	26452	30535	12196	88028
3	0.06386	42753	51364	60958	84436	58189	0.26514	75834	59521	37636	12078	98719
4	0.02916	62163	10236	51587	58769	27935	0.08712	40652	66430	32799	74653	52816
5	0.00384	29068	21880	78393	78013	20684	0.02344	30113	03591	62453	71921	03613
6	0.00183	20459	86444	14916	92261	41329	0.00610	50968	58694	42506	90720	59033
7	0.00009	73017	96964	77704	35972	03248	0.00137	03067	97499	58760	40919	95658
8	0.00010	88168	64034	75461	36273	53524	0.00029	98495	87774	13386	75427	15149
9	-0.00000	83124	80989	14474	80756	83316	0.00005	71279	77327	82762	62249	42045
10	0.00000	76093	36747	52906	90149	05470	0.00001	09812	77997	28649	23721	64909
11	-0.00000	16730	50119	56080	52965	01901	0.00000	17875	37686	00232	08394	11611
12	0.00000	06860	95605	33952	08525	32143	0.00000	03177	80645	24735	74359	66125
13	-0.00000	02068	17935	17460	02890	90873	0.00000	00427	02999	64944	38064	00163
14	0.00000	00715	53267	36586	97321	77786	0.00000	00078	08364	44191	12066	23204
15	-0.00000	00233	56503	79198	66091	37138	0.00000	00007	17143	23915	78700	00764
16	0.00000	00078	18902	04675	37681	86066	0.00000	00001	87654	45234	13299	33159
17	-0.00000	00025	94278	60137	72193	76381	0.00000	00000	01586	36642	25452	36344
18	0.00000	00008	64332	38788	78858	80196	0.00000	00000	05714	87184	82397	51814
19	-0.00000	00002	87697	66083	26943	17995	-0.00000	00000	00680	75693	40826	49802
20	0.00000	00000	95832	46989	44643	89904	0.00000	00000	00256	93394	02012	50700
21	-0.00000	00000	31923	54355	14763	58582	-0.00000	00000	00054	25128	23079	80583
22	0.00000	00000	10636	48853	89104	25473	0.00000	00000	00014	61950	44184	99559
23	-0.00000	00000	03544	28133	05353	85378	-0.00000	00000	00003	53643	92780	66646
24	0.00000	00000	01181	12896	81039	46306	0.00000	00000	00000	89493	51613	20739
25	-0.00000	00000	00393	63482	92231	22875	-0.00000	00000	00000	22239	88173	56536
26	0.00000	00000	00131	19298	21915	64240	0.00000	00000	00000	05567	76999	51411
27	-0.00000	00000	00043	72633	98654	72079	-0.00000	00000	00000	01390	20407	68723
28	0.00000	00000	00014	57428	48201	29110	0.00000	00000	00000	00347	51114	79725
29	-0.00000	00000	00004	85780	47285	87821	-0.00000	00000	00000	00086	84201	06978
30	0.00000	00000	00001	61919	57473	07582	0.00000	00000	00000	00021	70582	21355
31	-0.00000	00000	00000	53971	38024	88922	-0.00000	00000	00000	00005	42530	90710
32	0.00000	00000	00000	17990	00746	96841	0.00000	00000	00000	00001	35611	61567
33	-0.00000	00000	00000	05996	55604	63516	-0.00000	00000	00000	00000	33898	54027
34	0.00000	00000	00000	01998	82374	65548	0.00000	00000	00000	00000	08473	77501
35	-0.00000	00000	00000	00666	26751	66870	-0.00000	00000	00000	00000	02118	27096
36	0.00000	00000	00000	00222	08740	61997	0.00000	00000	00000	00000	00529	53328
37	-0.00000	00000	00000	00074	02869	39615	-0.00000	00000	00000	00000	00132	37643
38	0.00000	00000	00000	00024	67612	09747	0.00000	00000	00000	00000	00033	09273
39	-0.00000	00000	00000	00008	22534	60745	-0.00000	00000	00000	00000	00008	27291

40	0.00000	00000	00000	00002	74177	51295	0.00000	00000	00000	00000	00002	06817
41	-0.00000	00000	00000	00000	91392	33194	-0.00000	00000	00000	00000	00000	51703
42	0.00000	00000	00000	00000	30464	06756	0.00000	00000	00000	00000	00000	12926
43	-0.00000	00000	00000	00000	10154	67841	-0.00000	00000	00000	00000	00000	03231
44	0.00000	00000	00000	00000	03384	89011	0.00000	00000	00000	00000	00000	00808
45	-0.00000	00000	00000	00000	01128	29603	-0.00000	00000	00000	00000	00000	00202

Thus, coefficients given by (2.3) and (2.1) reduce to

$$b_\nu(a) = \frac{\pi^{\nu+1}}{2\nu!} \int_{-\infty}^{\infty} \frac{\sinh^\nu x \cosh x \exp\left(\left(a + \frac{1}{2}\right)\pi \sinh x - e^{\pi \sinh x}\right)}{\sinh\left(\frac{\pi}{2} \sinh x\right)} dx \quad (\nu \geq 1).$$

Similarly,

$$\alpha_\nu = \frac{\pi^{\nu+1}}{2(\nu+1)!} \int_{-\infty}^{\infty} \frac{\sinh^\nu x \cosh x \exp\left(\frac{1}{2}\pi \sinh x - e^{\pi \sinh x}\right)}{\sinh\left(\frac{\pi}{2} \sinh x\right)} [\pi \sinh x + \nu + 1] dx$$

for $\nu \geq 1$.

Using (2.11) in Q-arithmetic on the MICROVAX 3400 computer (machine precision $\approx 1.93 \times 10^{-34}$), we computed the coefficients $b_\nu(a)$ and $\beta_\nu(a)$ ($\nu \leq 60$) for $a = 0$ and $a = 1$. Here, we give in Table 2.1 the coefficients $\beta_\nu(0)$ and $\beta_\nu(1)$ only for $\nu = 0(1)45$ with 30D.

3. Chebyshev expansions

Let $T_k^*(z)$ be the shifted Chebyshev polynomial of the first kind of degree k given by $T_k^*(z) = T_k(2z - 1)$, $T_k(x) = \cos k\theta$ ($x = \cos \theta$). Numerical values of the coefficients in the Chebyshev expansions for the gamma function

$$\Gamma(z+1) = \sum_{k=0}^{\infty} \gamma_k T_k^*(z) \quad \text{and} \quad \frac{1}{\Gamma(z+1)} = \sum_{k=0}^{\infty} \delta_k T_k^*(z) \quad (0 \leq z \leq 1),$$

were given by Clenshaw [2] to 20D (see also Luke [10, p. 13]).

In this section we obtain similar expansions for the Kurepa function. Namely, we determine the coefficients c_k and d_k in the Chebyshev expansions

$$K(z+1) = \sum_{k=0}^{\infty} c_k T_k^*(z) \quad \text{and} \quad \frac{1}{K(z+1)} = \sum_{k=0}^{\infty} d_k T_k^*(z) \quad (0 \leq z \leq 1). \quad (3.1)$$

Taking $x = 2z - 1$, we can reduce our problem to the interval $(-1, 1)$ and then use the least square approximation, with respect to the Chebyshev weight $w(x) = (1 - x^2)^{-1/2}$, in order to find the coefficients c_k and d_k . Thus,

$$c_k = \frac{(K((x+3)/2), T_k)}{\|T_k\|^2}, \quad d_k = \frac{(1/K((x+3)/2), T_k)}{\|T_k\|^2} \quad (k \geq 0),$$

where the inner product and the norm are given by

$$(f, g) = \int_{-1}^1 f(x)\overline{g(x)}w(x) dx, \quad \|f\|^2 = (f, f),$$

respectively.

TABLE 3.1: The coefficients c_k and d_k in series (3.1)

k	c_k						d_k					
0	1.47404	93837	23786	91003	52009	68753	0.71933	88033	78564	80199	26920	51918
1	0.49590	32661	94158	22700	15501	19905	-0.24486	53611	27056	06741	80872	17436
2	0.02579	30869	03070	45788	53165	66164	0.02976	94905	09647	98622	36085	35374
3	0.00407	08812	31791	61386	34001	49176	-0.00498	33867	56921	33739	51852	46148
4	0.00015	73577	18751	30746	82134	28808	0.00086	57497	58509	74532	38298	56153
5	0.00002	57037	05826	92036	26292	93518	-0.00014	67971	82657	38722	42886	97797
6	0.00000	01775	04779	26512	29793	59549	0.00002	51924	09034	18335	36671	90502
7	0.00000	01478	12089	72264	27761	10704	-0.00000	43238	54800	40900	43392	52848
8	-0.00000	00057	65834	77355	97860	46608	0.00000	07414	54647	70219	98837	74381
9	0.00000	00010	46615	98880	68659	86996	-0.00000	01272	20052	30389	19464	41242
10	-0.00000	00000	83632	76388	29507	65862	0.00000	00218	27577	03559	46944	44963
11	0.00000	00000	09422	94336	75632	53320	-0.00000	00037	44928	13250	51616	37824
12	-0.00000	00000	00910	90916	80081	86592	0.00000	00006	42530	56969	56371	69126
13	0.00000	00000	00093	62274	78249	23554	-0.00000	00001	10240	75381	57048	29851
14	-0.00000	00000	00009	39687	44816	46570	0.00000	00000	18914	30308	43243	17659
15	0.00000	00000	00000	95140	34153	77730	-0.00000	00000	03245	18210	36461	92610
16	-0.00000	00000	00000	09603	61612	70033	0.00000	00000	00556	78518	95627	89374
17	0.00000	00000	00000	00970	39299	20617	-0.00000	00000	00095	52923	18136	84543
18	-0.00000	00000	00000	00098	02159	05061	0.00000	00000	00016	39022	51658	48302
19	0.00000	00000	00000	00009	90238	99413	-0.00000	00000	00002	81211	80450	47223
20	-0.00000	00000	00000	00001	00033	64229	0.00000	00000	00000	48248	31777	95406
21	0.00000	00000	00000	00000	10105	45983	-0.00000	00000	00000	08278	10261	20992
22	-0.00000	00000	00000	00000	01020	85778	0.00000	00000	00000	01420	29786	64079
23	0.00000	00000	00000	00000	00103	12756	-0.00000	00000	00000	00243	68458	86448
24	-0.00000	00000	00000	00000	00010	41800	0.00000	00000	00000	00041	80966	55298
25	0.00000	00000	00000	00000	00001	05243	-0.00000	00000	00000	00007	17340	45282
26	-0.00000	00000	00000	00000	00000	10632	0.00000	00000	00000	00001	23076	16403
27	0.00000	00000	00000	00000	00000	01074	-0.00000	00000	00000	00000	21116	53134
28	-0.00000	00000	00000	00000	00000	00108	0.00000	00000	00000	00000	03623	02400
29	0.00000	00000	00000	00000	00000	00011	-0.00000	00000	00000	00000	00621	61264
30	-0.00000	00000	00000	00000	00000	00001	0.00000	00000	00000	00000	00106	65187

Since $\|T_0\|^2 = \pi$ and $\|T_k\|^2 = \pi/2$ ($k \geq 1$), we have

$$c_0 = \frac{1}{\pi} \int_{-1}^1 \frac{1}{\sqrt{1-x^2}} K\left(\frac{x+3}{2}\right) dx, \quad c_k = \frac{2}{\pi} \int_{-1}^1 \frac{1}{\sqrt{1-x^2}} K\left(\frac{x+3}{2}\right) T_k(x) dx,$$

where $k \geq 1$. Similarly,

$$d_0 = \frac{1}{\pi} \int_{-1}^1 \frac{1}{\sqrt{1-x^2}} \frac{1}{K\left(\frac{x+3}{2}\right)} dx, \quad d_k = \frac{2}{\pi} \int_{-1}^1 \frac{1}{\sqrt{1-x^2}} \frac{T_k(x)}{K\left(\frac{x+3}{2}\right)} dx \quad (k \geq 1).$$

For numerical determination of these coefficients, we apply the N -point Gauss-Chebyshev quadrature formula (see [13])

$$\int_{-1}^1 \frac{1}{\sqrt{1-x^2}} f(x) dx = \frac{\pi}{N} \sum_{k=1}^N f(x_k^{(N)}) + R_N(f),$$

where $x_k^{(N)}$, $k = 1, \dots, N$, are the zeros of N -th Chebyshev polynomial $T_N(x)$, i.e., $x_k^{(N)} = \cos\left((2k-1)\frac{\pi}{2N}\right)$. This formula is exact for all algebraic polynomials of degree at most $2N-1$, i.e., $R_N(\mathcal{P}_{2N-1}) \equiv 0$. To compute values of the Kurepa function we use the power series expansion from Section 2. In order to obtain numerical values to 30D of coefficients c_k and d_k for $k \leq 30$, the Gaussian formula (3.1) needs $N = 40$ and $N = 35$ points, respectively. Numerical values are displayed in Table 3.1.

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