Research Article

Homotopy Perturbation Method and Variational Iteration Method for Harmonic Waves Propagation in Nonlinear Magneto-Thermoelasticity with Rotation

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The homotopy perturbation method and variational iteration method are applied to obtain the approximate solution of the harmonic waves propagation in a nonlinear magneto-thermoelasticity under influence of rotation. The problem is solved in one-dimensional elastic half-space model subjected initially to a prescribed harmonic displacement and the temperature of the medium. The displacement and temperature are calculated for the methods with the variations of the magnetic field and the rotation. The results obtained are displayed graphically to show the influences of the new parameters and the difference between the methods' technique. It is obvious that the homotopy perturbation method is more effective and powerful than the variational iteration method.

1. Introduction

In the past recent years, much attentions have been devoted to simulate some real-life problems which can be described by nonlinear coupled differential equations using reliable and more efficient methods. The nonlinear coupled system of partial differential equations often appear in the study of circled fuel reactor, high-temperature hydrodynamics, and thermoelasticity problems, see [1–4]. From the analytical point of view, lots of work have been done for such systems. With the rapid development of nanotechnology, there appears an everincreasing interest of scientists and researchers in this field of science. Nanomaterials, because of their exceptional mechanical, physical, and chemical properties, have been the main topic of research in many scientific publications. Wave generation in nonlinear thermoelasticity problems has gained a considerable interest for its utilitarian aspects in understanding the nature of interaction between the elastic and thermal fields as well as for its applications. A lot of applications was paid on existence, uniqueness, and stability of the solution of the problem, see [5–7].

Much attention has been devoted to numerical methods, which do not require discretization of space-time variables or linearization of the nonlinear equations, among which the variational iteration method (VIM) suggested in [8–20] shows its remarkable merits over others. The method was successfully applied to a nonlinear one dimensional coupled equations in thermoelasticity [21], revealing that the method is very convenient, efficient, and accurate. The basic idea of variational iteration method is to construct a correction functional with a general Lagrange multiplier which can be identified optimally via variational theory.

The homotopy perturbation method [8, 22] has the merits of simplicity and easy execution. Unlike the traditional numerical methods, the HPM does not need discretization and linearization. Most perturbation methods assume that a small parameter exists, but most nonlinear problems have no small parameter at all. Many new methods have been proposed to eliminate the small parameter. Recently, the applications of homotopy theory among scientists appeared, and the homotopy theory becomes a powerful mathematical tool, when it is successfully coupled with perturbation theory. Sweilam and Khader [1] investigated variational iteration method for one dimensional nonlinear thermoelasticity. Applying He's variational iteration method for solving differential-difference equation is discussed by Yildirim [23]. Noor and Mohyud-Din [24], Mohyud-Din et al. [25-27] used He's polynomials or Padé approximants to solve solving higher-order nonlinear boundary value problems, second-order singular problems, and nonlinear boundary value problems. Mohyud-Din et al. [28] applied the modified variational iteration method for free-convective boundary-layer equation using Padé approximation. Mohyud-Din and Noor [29, 30] used Homotopy perturbation method for solving some new boundary value problems. Mohyud-Din et al. [31] investigated some relatively new techniques for nonlinear problems.

In this paper, the homotopy perturbation method and variational iteration method are used to solve the coupled harmonic waves nonlinear magneto-thermoelasticity equations under influence of rotation. The Maple and Mathematica software packages are used to obtain the approximate solutions in one-dimensional half-space. The displacement and temperature which obtained have been calculated numerically and presented graphically.

2. Basic Idea of He's Homotopy Perturbation Method

We illustrate the following nonlinear differential equation [8, 22]:

$$A(u) - f(r) = 0, \quad r \in \Lambda, \tag{2.1}$$

with the boundary conditions:

$$B\left(u,\frac{\partial u}{\partial n}\right) = 0, \quad r \in \Gamma,$$
(2.2)

where *A* is a general differential operator, *B* is a boundary operator, f(r) is an analytic function, and Γ is the boundary of the domain Λ . Generally speaking, the operator *A* can be divided into two parts which are *L* and *N*, where *L* is linear operator but *N* is nonlinear operator. Equation (2.1) can therefore be rewritten as follows:

$$L(u) + N(u) - f(r) = 0.$$
(2.3)

By the homotopy technique, we construct a homotopy $V(r, p): \Lambda \times [0, 1] \rightarrow R$ which satisfies

$$H(V,p) = (1-p)[L(V) - L(u_0)] + p[A(V) - f(r)] = 0, \quad r \in \Lambda,$$
(2.4)

or

$$H(V,p) = L(V) - L(u_0) + pL(u_0) + p(N(V) - f(r)) = 0, \quad r \in \Lambda,$$
(2.5)

where $p \in [0, 1]$ is an embedding parameter and u_0 is an initial approximation of (2.1) which satisfies the boundary conditions (2.2). Obviously, from (2.4) and (2.5) we have

$$H(V,0) = L(V) - L(u_0) = 0,$$

$$H(V,1) = A(V) - f(r) = 0.$$
(2.6)

The changing process of p from zero to unity is just that of V(r, p) from $u_0(r)$ to u(r). In topology, this is called deformation, and $L(V) - L(u_0)$ and A(V) - f(r) are called homotopy. According to the homotopy perturbation method, we can first use the embedding parameter "p" as a small parameter and assume that the solution of (2.4) and (2.5) can be written as a power series in "p" as follows:

$$V = V_0 + pV_1 + p^2 V_2 + \cdots . (2.7)$$

On setting p = 1 results in the approximate solution of (2.3), we have

$$u = \lim_{p \to 1} V = V_0 + V_1 + V_2 + \cdots .$$
(2.8)

The combination of the perturbation method and the homotopy method is called the homotopy perturbation method, which has eliminated the limitations of the traditional perturbation methods. On the other hand, this technique can have full advantage of the traditional perturbation techniques. The series (2.8) is convergent to most cases. However, the convergent rate depends on the nonlinear operator A(V).

- (1) The second derivative of N(V) with respect to V must be small because the parameter may be relatively large, that is, $p \rightarrow 1$.
- (2) The norm of $L^{-1}(\partial N/\partial V)$ must be smaller than one so that the series converges.

3. Application of Homotopy Perturbation Method on the Nonlinear Magneto-Thermoelastic with Rotation Equations

In this section, we use the homotopy perturbation method to calculate the approximate solutions of the following nonlinear magneto-thermoelastic with rotation equations:

$$(1 + \sigma_1)u_{tt} + \Omega u_t - u_{xx} \left(1 - \sigma_2 + 2\gamma u_x + 3\delta u_x^2 \right) - \beta_1 \theta_x - \beta_2 (\theta u_x)_x = 0,$$

$$\left(\theta - au_x - \frac{1}{2}bu_x^2 \right)_t - \left[(1 + \alpha u_x)\theta_x \right]_x = 0,$$
(3.1)

where γ , β_1 , β_2 , a, b, α are arbitrary constants, σ_1 , σ_2 are the sensitive parts of the magnetic field, and Ω is the rotation parameter, with the initial conditions

$$u(x,0) = \theta(x,0) = A(1 - \cos(x)), \qquad u_t(x,0) = \theta_t(x,0) = 0, \tag{3.2}$$

where A is an arbitrary constant and the boundary conditions

$$u(0,t) = \theta(0,t) = 0, \qquad u_t(0,t) = \theta_t(0,t) = 0.$$
(3.3)

To investigate the traveling wave solution of (3.1), we first construct a homotopy perturbation method as follows:

$$(1-p)[(1+\sigma_{1})(V_{tt}-V_{0tt})] + p[(1+\sigma_{1})V_{tt} + \Omega V_{t} - V_{xx}(1-\sigma_{2}+2\gamma V_{x}+3\delta V_{x}^{2}) -\beta_{1}\Theta_{x} - \beta_{2}(\Theta V_{x})_{x}] = 0,$$
(3.4)
$$(1-p)(\Theta_{t}-\Theta_{0t}) + p[\Theta_{t}-aV_{xt}-bV_{x}V_{xt}-\Theta_{xx}-aV_{xx}\Theta_{x}-aV_{x}\Theta_{xx}] = 0,$$

where the initial approximations take the following form:

$$V_0(x,t) = u_0(x,t) = u(x,0) = A(1 - \cos(x)),$$

$$\Theta_0(x,t) = \theta_0(x,t) = \theta(x,0) = A(1 - \cos(x)).$$
(3.5)

According to the homotopy perturbation method, we can first use the embedding parameter "p" as a small parameter and assume that the solution of (3.4) can be written as a power series in "p" as the following:

$$V = V_0(x,t) + pV_1(x,t) + p^2V_2(x,t) + p^3V_3(x,t) + \cdots,$$

$$\Theta(x,t) = \Theta_0(x,t) + p\Theta_1(x,t) + p^2\Theta_2(x,t) + p^3\Theta_3(x,t) + \cdots,$$
(3.6)

where V_j and Θ_j , j = 1, 2, 3, ... are functions to be determined.

Substituting from (3.6) into (3.4) and arranging the coefficients of "p" powers, we have

$$\begin{aligned} (1+\sigma_{1})V_{0,tt} + \left(\sigma_{1}V_{1,tt} - 3\delta V_{0,x}^{2}V_{0,xx} - \beta_{2}V_{0,xx}\Theta_{0} - 2\gamma V_{0,x}V_{0,xx} + \sigma_{2}V_{0,xx} - \beta_{1}\Theta_{0,x} \right. \\ \left. -V_{0,xx} - \beta_{2}V_{0,x}\Theta_{0,x} + \Omega V_{0,t} + V_{1,tt}\right)p \\ + \left(\sigma_{1}V_{2,tt} - 2\gamma V_{1,x}V_{0,xx} + V_{2,tt} - 2\gamma V_{0,x}V_{1,xx} - 6\delta V_{0,x}V_{1,x}V_{0,xx} - V_{1,xx} - \beta_{2}V_{1,xx}\Theta_{0} \right. \\ \left. -\beta_{1}\Theta_{1,x} - \beta_{2}V_{0,x}\Theta_{1,x} - \beta_{2}V_{1,x}\Theta_{0,x} - \beta_{2}V_{0,xx}\Theta_{1} + \Omega V_{1,t} - 3\delta V_{0,x}^{2}V_{1,xx} + \sigma_{2}V_{1,xx}\right)p^{2} \\ + \left(V_{3,tt} - \beta_{2}V_{2,x}\Theta_{0,x} - \beta_{2}V_{2,xx}\Theta_{0} - 2\gamma V_{2,x}V_{0,xx} - 2\gamma V_{1,x}V_{1,xx} + \sigma_{1}V_{3,tt} - \beta_{1}\Theta_{2,x} \right. \\ \left. - 6\delta V_{0,x}V_{2,x}V_{0,xx} - 3\delta V_{1,x}^{2}V_{0,xx} + \sigma_{2}V_{2,xx} + \Omega V_{2,t} - V_{2,xx} - 3\delta V_{0,x}^{2}V_{2,xx} - 2\gamma V_{0,x}V_{2,xx} \right. \\ \left. -\beta_{2}V_{1,xx}\Theta_{1} - \beta_{2}V_{0,xx}\Theta_{2} - \beta_{2}V_{0,x}\Theta_{2,x} - 6\delta V_{0,x}V_{1,x}V_{1,xx} - \beta_{2}V_{1,x}\Theta_{1,x}\right)p^{3} + \dots = 0, \\ \Theta_{0,t} + \left(-aV_{0,xt} - aV_{0,x}\Theta_{0,xx} + \Theta_{1,t} - aV_{0,xx}\Theta_{0,x} - 2bV_{0,x}V_{0,xt} - \Theta_{0,xx}\right)p \\ + \left(\Theta_{2,t} - \Theta_{1,xx} - aV_{1,xt} - aV_{1,xx}\Theta_{0,x} - 2bV_{1,x}V_{0,xt} - aV_{0,x}\Theta_{1,xx} - aV_{1,x}\Theta_{0,xx} \right. \\ \left. -2bV_{0,x}V_{1,xt} - aV_{0,xx}\Theta_{1,x}\right)p^{2} \\ \times \left(-aV_{1,xx}\Theta_{1,x} - aV_{2,xx}\Theta_{0,x} - 2bV_{2,x}V_{0,xt} - aV_{0,xx}\Theta_{2,x} - 2bV_{1,x}V_{1,xt} - aV_{2,xt} \right. \\ \left. +\Theta_{3,t} - 2bV_{0,x}V_{2,xt} - aV_{0,x}\Theta_{2,xx} - aV_{1,x}\Theta_{1,xx} - aV_{2,x}\Theta_{0,xx} - \Theta_{2,xx}\right)p^{3} + \dots = 0. \end{aligned}$$

$$(3.7)$$

In order to obtain the unknowns of V_j and Θ_j , (j = 1, 2, 3, ...), we construct and solve the following system considering the initial conditions (3.2):

$$\begin{split} &\sigma_{1}V_{1,tt} - 3\delta V_{0,x}^{2}V_{0,xx} - \beta_{2}V_{0,xx}\Theta_{0} - 2\gamma V_{0,x}V_{0,xx} + \sigma_{2}V_{0,xx} - \beta_{1}\Theta_{0,x} - V_{0,xx} \\ &- \beta_{2}V_{0,x}\Theta_{0,x} + \Omega V_{0,t} + V_{1,tt} = 0, \\ &\sigma_{1}V_{2,tt} - 2\gamma V_{1,x}V_{0,xx} + V_{2,tt} - 2\gamma V_{0,x}V_{1,xx} - 6\delta V_{0,x}V_{1,x}V_{0,xx} - V_{1,xx} - \beta_{2}V_{1,xx}\Theta_{0} \\ &- \beta_{1}\Theta_{1,x} - \beta_{2}V_{0,x}\Theta_{1,x} - \beta_{2}V_{1,x}\Theta_{0,x} - \beta_{2}V_{0,xx}\Theta_{1} + \Omega V_{1,t} - 3\delta V_{0,x}^{2}V_{1,xx} + \sigma_{2}V_{1,xx} = 0, \\ &V_{3,tt} - \beta_{2}V_{2,x}\Theta_{0,x} - \beta_{2}V_{2,xx}\Theta_{0} - 2\gamma V_{2,x}V_{0,xx} - 2\gamma V_{1,x}V_{1,xx} + \sigma_{1}V_{3,tt} - \beta_{1}\Theta_{2,x} \\ &- 6\delta V_{0,x}V_{2,x}V_{0,xx} - 3\delta V_{1,x}^{2}V_{0,xx} + \sigma_{2}V_{2,xx} + \Omega V_{2,t} - V_{2,xx} - 3\delta V_{0,x}^{2}V_{2,xx} - 2\gamma V_{0,x}V_{2,xx} \\ &- \beta_{2}V_{1,xx}\Theta_{1} - \beta_{2}V_{0,x}\Theta_{2} - \beta_{2}V_{0,x}\Theta_{2,x} - 6\delta V_{0,x}V_{1,x}V_{1,xx} - \beta_{2}V_{1,x}\Theta_{1,x} = 0, \\ &- aV_{0,xt} - aV_{0,x}\Theta_{0,xx} + \Theta_{1,t} - aV_{0,xx}\Theta_{0,x} - 2bV_{0,x}V_{0,xt} - \Theta_{0,xx} = 0, \\ &\Theta_{2,t} - \Theta_{1,xx} - aV_{1,xt} - aV_{1,xx}\Theta_{0,x} - 2bV_{1,x}V_{0,xt} - aV_{0,x}\Theta_{1,xx} - aV_{1,x}\Theta_{0,xx} \\ &- 2bV_{0,x}V_{1,xt} - aV_{0,x}\Theta_{1,x} = 0, \\ &- aV_{1,xx}\Theta_{1,x} - aV_{2,xx}\Theta_{0,x} - 2bV_{2,x}V_{0,xt} - aV_{0,x}\Theta_{2,x} - 2bV_{1,x}V_{1,xt} - aV_{2,xt} + \Theta_{3,t} \\ &- 2bV_{0,x}V_{2,xt} - aV_{0,x}\Theta_{2,xx} - aV_{1,x}\Theta_{1,xx} - aV_{2,x}\Theta_{0,xx} - \Theta_{2,xx} = 0. \end{split}$$

(3.8)

Consequently, we deduce after some calculations the following results:

$$u = \lim_{p \to 1} V = V_0 + V_1 + V_2 + \cdots,$$

$$\theta = \lim_{p \to 1} \Theta = \Theta_0 + \Theta_1 + \Theta_2 + \cdots,$$
(3.9)

where

$$\begin{split} V_0 &= A(1-\cos x), \\ V_1 &= \frac{\hbar^2}{2(4+4\sigma_1)} \Big[4\gamma A^2 \sin 2x + 3\delta A^3 \cos x - 3\delta A^3 \cos 3x + 4\beta_2 A^2 \cos x \\ &\quad -4\beta_2 A^2 \cos 2x 4\sigma_2 A \cos x + 4\beta_1 A \sin x + 4A \cos x \Big], \\ V_2 &= \frac{\hbar^4}{32(\sigma_1+1)^2} \Big\{ \Big(-\frac{8}{3}\beta_2^2 A^3 - \frac{8}{3}\beta_2 A^2 - \frac{8}{3}\gamma^2 A^3 - 4\delta A^3 + 4A^3\sigma_2 \delta + \frac{8}{3}\sigma_2 A - \frac{9}{2}\delta^2 A^5 - \frac{4}{3}A\sigma_2^2 \\ &\quad -\frac{4}{3}A + \frac{8}{3}A^2\sigma_2\beta_2 - 4\beta_2 A^4 \delta \Big) \cos x \\ &\quad - \Big(\frac{4}{3}\beta_1 A + \delta A^3\beta_1 + \frac{4}{3}\gamma A^3\beta_2 - \frac{4}{3}A\sigma_2\beta_1 + \frac{4}{3}\beta_2 A^2\beta_1 \Big) \sin x \\ &\quad + \Big(-8\gamma A^2 + 8A^2\sigma_2\gamma - 8\beta_2 A^3\gamma + \frac{4}{3}\beta_2 A^2\beta_1 - 16\delta A^4\gamma \Big) \sin 2x \\ &\quad + \Big(12\gamma A^3\beta_2 + 3\delta A^3\beta_1 \Big) \sin 3x + 20\delta A^4\gamma \sin 4x \\ &\quad + \Big(-\frac{20}{3}A^2\sigma_2\beta_2 + \frac{20}{3}\beta_2^2 A^3 + \frac{8}{3}\gamma A^2\beta_1 + \frac{20}{3}\beta_2 A^2 + 6\delta A^4\beta_2 \Big) \cos 2x \\ &\quad + \Big(\frac{63}{4}\delta^2 A^5 - 4\beta_2^2 A^3 + 12\delta A^3 + 12\beta_2 A^4 \delta + 8\gamma^2 A^3 - 12A^3\sigma_2 \delta \Big) \cos 3x \\ &\quad -14\beta_2 A^4 \delta \cos 4x - \frac{45}{4}\delta^2 A^5 \cos 5x \Big\} \\ &\quad + \frac{\hbar^3}{32(\sigma_1 + 1)^2} \Big\{ \Big(-\frac{16}{3}\Omega A\beta_1 - \frac{16}{3}\beta_1 A\sigma_1 - \frac{8}{3}\beta_2 a\sigma_1 A^3 - \frac{16}{3}\beta_1 A - \frac{8}{3}\beta_2 A^3 a \Big) \sin x \\ &\quad - \frac{16}{3}\Omega A^2\gamma \sin 2x + \Big(8\beta_2 A^3 a\sigma_1 + 8\beta_2 A^3 a \Big) \sin 3x \\ &\quad + \Big(\frac{16}{3}\Omega A\sigma_2 - 4\Omega A^3 \delta - \frac{16}{3}\Omega A^2\beta_2 - \frac{16}{3}\Omega A \Big) \cos(x) \\ &\quad + \Big(\frac{16}{3}\beta_2 A^2 + \frac{32}{3}\beta_1 A^2 a\sigma_1 + \frac{16}{3}\Omega A^2\beta_2 + \frac{16}{3}\beta_2 A^2\sigma_1 + \frac{32}{3}\beta_1 A^2a \Big) \cos 2x \\ &\quad + 4\Omega A^3 \delta \cos 3x \Big\}, \end{split}$$

$$\begin{split} \Theta_{0} &= A(1 - \cos(x)), \\ \Theta_{1} &= A^{2}ta\sin(2x) + A\cos(x)t, \\ \Theta_{2} &= \frac{A}{2(1 + \sigma_{1})} \left(\frac{t^{3}}{3} \left(12\alpha A^{2}\gamma\cos^{3}(x) + 2\alpha A\cos^{2}(x)\beta_{1} + 36\alpha A^{3}\delta\sin(x)\cos^{3}(x) \right. \\ &+ 12\alpha A^{2}\beta_{2}\sin(x)\cos^{2}(x) + 2\alpha A\sigma_{2}\sin(x)\cos(x) - 24\alpha A^{3}\delta\sin(x)\cos(x) \\ &- 2\alpha A^{2}\beta_{2}\sin(x)\cos(x) - 2\alpha A\sin(x)\cos(x) - \alpha A\beta_{1} - 10\alpha\gamma A^{2}\cos(x) \\ &- 4\alpha\beta_{2}A^{2}\sin(x) \right) \\ &+ \frac{t^{2}}{2} \left(-2\cos(x) + 18\alpha\delta A^{2}\sin(x)\cos^{2}(x) + 16bA^{2}\gamma\sin(x)\cos^{2}(x) \\ &+ 8\alpha\beta_{2}A\sin(x)\cos(x) + 2\alpha\sigma_{2}\sin(x) + 2\alpha\beta_{1}\cos(x) + 24\alpha^{2}A^{2}\sigma_{1}\cos^{3}(x) \\ &- 36bA^{3}\delta\cos^{4}(x) - 16bA^{2}\beta_{2}\cos^{3}(x) + 48bA^{3}\delta\cos^{2}(x) + 4bA^{2}\beta_{2}\cos^{2}(x) \\ &+ 8\cos^{2}(x)a\gamma A - 4\cos(x)^{2}bA\sigma_{2} - 20a^{2}A^{2}\cos(x) + 24\cos(x)^{3}a^{2}A^{2} \\ &+ 4bA\cos^{2}(x) - 4bA - 6\alpha\delta A^{2}\sin(x) - 20\alpha A\sin(x)\cos(x) \\ &- 2\alpha\beta_{2}A\sin(x) - 4bA^{2}\beta_{2} - 12bA^{3}\delta + 4bA\sigma_{2} - 8bA^{2}\sin(x)\gamma \\ &+ 16bA^{2}\beta_{2}\cos(x) - 20a^{2}A^{2}\cos(x) - 20\alpha A\sin(x)\cos(x)\sigma_{1} \right) \right). \end{split}$$

$$(3.10)$$

Now we make calculations for the results obtained by the homotopy perturbation method using the Maple software package with the following arbitrary constants:

 $a = 0.5, \quad A = 0.001, \quad b = 0.5, \quad \alpha = 1, \quad \beta_1 = \beta_2 = 0.05, \quad \gamma = 1, \quad \delta = 0.8.$ (3.11)

The results obtained in (3.9) are displayed graphically in Figures 1-4.

3.1. Special Cases

- (1) If we take into our consideration the first iteration (i.e., $u = V_0 + V_1$ and $\theta = \Theta_0 + \Theta_1$). See Figures 5, 6, 7, and 8.
- (2) If the magnetic field and rotation are neglected, the components of the displacement u and temperature θ take the following forms. See Figures 9 and 10.



Figure 1: Variations of the displacement *u* and temperature θ for various values of the axis *x* and time *t* when $\Omega = 0.1$, $\sigma_1 = 0.2$, $\sigma_2 = 0.1$.



Figure 2: Variations of the displacement *u* and temperature θ for various values of the axis *x* and rotation Ω when t = 0.1, $\sigma_1 = 0.2$, $\sigma_2 = 0.1$.



Figure 3: Variations of the displacement *u* and temperature θ for various values of the axis *x* and magnetic field σ_1 when t = 0.1, $\Omega = 0.1$, $\sigma_2 = 0.1$.



Figure 4: Variations of the displacement *u* and temperature θ for various values of the axis *x* and magnetic field σ_2 when t = 0.1, $\Omega = 0.1$, $\sigma_1 = 0.1$.



Figure 5: Variations of the displacement *u* and temperature θ for various values of the axis *x* and time *t* when $\Omega = 0.1$, $\sigma_1 = 0.2$, $\sigma_2 = 0.1$.



Figure 6: Variations of the displacement *u* and temperature θ for various values of the axis *x* and rotation Ω when t = 0.1, $\sigma_1 = 0.2$, $\sigma_2 = 0.1$.



Figure 7: Variations of the displacement *u* and temperature θ for various values of the axis *x* and magnetic field σ_1 when t = 0.1, $\Omega = 0.1$, $\sigma_2 = 0.1$.



Figure 8: Variations of the displacement *u* and temperature θ for various values of the axis *x* and magnetic field σ_2 when t = 0.1, $\Omega = 0.1$, $\sigma_1 = 0.1$.

4. Basic Idea of Variational Iteration Method

Consider the following nonhomogeneous nonlinear system of partial differential equations:

$$L_1 u(x,t) + N_1 (u(x,t), \theta(x,t)) = f(x,t),$$
(4.1)

$$L_2\theta(x,t) + N_2(u(x,t),\theta(x,t)) = g(x,t),$$
(4.2)



Figure 9: Variations of the displacement *u* and temperature θ for various values of the axis *x* and time *t* $(u = V_0 + V_1 + V_2 \text{ and } \theta = \Theta_0 + \Theta_1 + \Theta_2)$ when $\Omega = \sigma_1 = \sigma_2 = 0$.



Figure 10: Variations of the displacement *u* and temperature Θ for various values of the axis *x* and time *t* ($u = V_0 + V_1$ and $\theta = \Theta_0 + \Theta_1$) when $\Omega = \sigma_1 = \sigma_2 = 0$.

where L_1 , L_2 are linear differential operators with respect to time, N_1 , N_2 are nonlinear operators, and f(x,t), g(x,t) are given functions.

According to the variational iteration method, we can construct correct functionals as follows:

$$u_{n+1}(x,t) = u_n(x,t) + \int_0^t \lambda_1(\tau) \Big[L_1 u_n(x,\tau) + N_1 \Big(\tilde{u}_n(x,\tau), \tilde{\theta}_n(x,\tau) \Big) - f(x,\tau) \Big] d\tau,$$
(4.3)

$$\theta_{n+1}(x,t) = \theta_n(x,t) + \int_0^t \lambda_2(\tau) \Big[L_2 \theta_n(x,\tau) + N_2 \Big(\widetilde{u}_n(x,\tau), \widetilde{\theta}_n(x,\tau) \Big) - g(x,\tau) \Big] d\tau, \qquad (4.4)$$

where λ_1 and λ_2 are general Lagrange multipliers, which can be identified optimally via variational theory [8–20]. The second term on the right-hand side in (4.3) and (4.4) is called the corrections, and the subscript *n* denotes the *n*th order approximation, \tilde{u}_n and $\tilde{\theta}_n$ are restricted variations. We can assume that the above correctional functionals are stationary (i.e., $\delta u_{n+1} = 0$ and $\delta \theta_{n+1} = 0$), then the Lagrange multipliers can be identified. Now we can start with the given initial approximation and by the previous iteration formulas we can obtain the approximate solutions.

5. Application of the Variational Iteration Method on the Nonlinear Magneto-Thermoelastic with Rotation Equations

According to the variational iteration method and after some manipulation of (4.3) and (4.4), the correct functionals are as follows:

$$u_{n+1}(x,t) = u_n(x,t) + \int_0^t \lambda_1(\tau) \Big[(1+\sigma_1) u_{n,tt}(x,\tau) + \Omega \widetilde{u}_{n,t}(x,\tau) \\ - \widetilde{u}_{n,xx} \Big(1-\sigma_2 + 2\gamma \widetilde{u}_{n,x}(x,\tau) + 3\delta \widetilde{u}_{n,x}^2(x,\tau) \Big) \\ - \beta_1 \widetilde{\theta}_{n,x}(x,\tau) - \beta_2 \Big(\widetilde{\theta}_n(x,\tau) \widetilde{u}_{n,x}(x,\tau) \Big)_x \Big] d\tau, \\ \theta_{n+1}(x,t) = \theta_n(x,t) + \int_0^t \lambda_2(\tau) \Big[\theta_{n,t}(x,\tau) - a \widetilde{u}_{n,xt}(x,\tau) - b \widetilde{u}_{n,x}(x,\tau) \widetilde{u}_{n,xt}(x,\tau) \\ - \widetilde{\theta}_{n,xx}(x,\tau) - \alpha \widetilde{u}_{n,xx}(x,\tau) \widetilde{\theta}_{n,x}(x,\tau) - \alpha \widetilde{u}_{n,xx}(x,\tau) \Big] d\tau,$$

$$(5.1)$$

where \tilde{u}_n and $\tilde{\theta}_n$ are considered as a restricted variation, that is, $\delta \tilde{u}_{n+1} = 0$ and $\delta \tilde{\theta}_{n+1} = 0$. Consequently, the general Lagrange multipliers λ_1 and λ_2 take the following form:

$$\lambda_1(\tau) = \frac{\tau - t}{1 + \sigma_1}, \qquad \lambda_2(\tau) = -1.$$
 (5.2)

By the substitution of the identified Lagrange multipliers (5.2) into (5.1), we have the following iteration relations:

$$u_{n+1}(x,t) = u_n(x,t) + \int_0^t \frac{\tau - t}{1 + \sigma_1} \Big[(1 + \sigma_1) u_{n,tt}(x,\tau) + \Omega u_{n,t}(x,\tau) \\ - u_{n,xx} \Big(1 - \sigma_2 + 2\gamma u_{n,x}(x,\tau) + 3\delta u_{n,x}^2(x,\tau) \Big) \\ - \beta_1 \theta_{n,x}(x,\tau) - \beta_2 (\theta_n(x,\tau) u_{n,x}(x,\tau))_x \Big] d\tau,$$

$$\theta_{n+1}(x,t) = \theta_n(x,t) - \int_0^t \Big[\theta_n(x,\tau) - au_{n,xt}(x,\tau) - bu_{n,x}(x,\tau)u_{n,xt}(x,\tau) - \theta_{n,xx}(x,\tau) \\ -\alpha u_{n,xx}(x,\tau)\theta_{n,x}(x,\tau) - \alpha u_{n,x}(x,\tau)\theta_{n,xx}(x,\tau) \Big] d\tau, \quad n \ge 0.$$
(5.3)

With help of Maple or Mathematica, we get the following results:

$$\begin{split} u_0 &= \theta_0 = A(1 - \cos x), \\ u_1 &= -\frac{t^2}{2(1 + \sigma_1)} \Big[-\beta_1 \sin x - 2A\gamma \cos x \sin(x) + 2\beta_2 A \cos^2 x - \cos x - 3\delta A^2 \cos x \\ &-\beta_2 A + \sigma_2 \cos x - \beta_2 A \cos x + 3\delta A^2 \cos^3 x \Big] + A(1 - \cos x), \\ \theta_1 &= A \big[\cos x + 2\alpha A \cos x \sin x \big] t + A(1 - \cos x), \end{split}$$

$$u_{2} = \left\{ -\frac{A\cos^{6}(x)}{6720(1+\sigma_{1})^{4}} \left(-19440A^{5}\delta^{2}\beta_{1}\gamma + 5760A^{5}\delta\beta_{2}^{3} - 247860A^{7}\delta^{3}\beta_{2} - 19440A^{6}\delta^{2}\beta_{2}^{2} \right. \right. \\ \left. -19440A^{5}\delta^{2}\beta_{2} - 17280\delta\beta_{2}\gamma^{2}A^{5} + 19440A^{5}\delta^{2}\beta_{2}\sigma_{2} \right) \right\}$$

$$+ \left[-\frac{A\cos^{2}(x)}{6720(1+\sigma_{1})^{4}} \left(-6480A^{5}\delta\beta_{2}^{2}\gamma + 2025A^{4}\delta\beta_{1}\beta_{2}^{2} + 270A^{3}\delta\beta_{1}\sigma_{2}\beta_{2} + 3240A^{4}\delta^{2}\beta_{1}\sigma_{2} \right. \\ \left. - 38880A^{6}\delta^{2}\beta_{2}\gamma - 8505A^{6}\delta^{3}\beta_{1} - 135A^{2}\delta\beta_{1}\sigma_{2}^{2} - 6480A^{4}\gamma\beta_{2}\delta_{2}\delta_{1} \right. \\ \left. - 135A^{2}\delta\beta_{1} + 6480A^{4}\delta\beta_{2}\gamma\sigma_{2} - 3240A^{5}\delta^{2}\beta_{1}\beta_{2} + 270A^{2}\delta\beta_{1}\sigma_{2} \right. \\ \left. - 270A^{3}\delta\beta_{1}\beta_{2} - 3240A^{4}\delta^{2}\beta_{1} - 2160A^{4}\delta\beta_{1}\gamma^{2} + 45A^{2}\delta\beta_{1}^{3} \right) \right] \\ \left. - \frac{A\cos^{4}(x)}{6720(1+\sigma_{1})^{4}} \left(7200A^{4}\gamma\beta_{2}\delta + 118800A^{6}\delta^{2}\beta_{2}\gamma - 4050A^{4}\delta^{2}\beta_{1}\sigma_{2} + 4050A^{4}\delta^{2}\beta_{1} \right. \\ \left. + 4050A^{5}\delta^{2}\beta_{1}\beta_{2} + 3600A^{4}\delta\beta_{1}\gamma^{2} - 3600A^{4}\delta\beta_{1}\beta_{2}^{2} \right. \\ \left. + 7200A^{5}\delta\beta_{2}^{2}\gamma + 30375A^{6}\delta^{3}\beta_{1} - 7200A^{4}\delta\beta_{2}\gamma\sigma_{2} \right) \right] \\ \left. - \frac{A\cos^{3}(x)}{6720(1+\sigma_{1})^{4}} \left(17280A^{5}\delta^{2}\beta_{1}\beta_{2} - 1440A^{4}\gamma\beta_{2}\delta - 93960A^{7}\delta^{3}\gamma + 16560A^{5}\delta\beta_{2}^{2}\gamma \right. \\ \left. + 1440A^{4}\delta\beta_{2}\gamma\sigma_{2} - 23760A^{6}\delta^{2}\beta_{2}\gamma + 1440A^{4}\delta\beta_{1}\beta_{2}^{2} - 720A^{3}\gamma\delta \right. \\ \left. + 1440A^{3}\delta\beta_{1}\beta_{2} + 23760A^{5}\delta^{2}\gamma\sigma_{2} - 720A^{3}\delta\beta_{1}^{2}\gamma + 1440A^{3}\gamma\delta\sigma_{2} \right) \right] \\ \left. - \frac{A\cos^{5}(x)}{6720(1+\sigma_{1})^{4}} \left(-19440A^{5}\delta^{2}\beta_{1}\beta_{2} + 189540A^{7}\delta^{3}\gamma + 19440A^{6}\delta^{2}\beta_{2}\gamma - 17280A^{5}\delta\beta_{2}^{2}\gamma \right) \right] \right]$$

$$\begin{split} &-\frac{A\cos(x)}{6720(1+\sigma_1)^4} \left(-1080A^3\gamma \delta \sigma_2 + 540A^3\gamma \delta + 540A^3\delta \gamma \sigma_2^2 - 180A^3\delta \beta_1^2 \gamma \right. \\ &-1080A^4 \delta \beta_2 \gamma \sigma_2 - 720A^4 \delta \beta_1 \beta_2^2 - 2160A^5 \delta^2 \beta_1 \beta_2 + 1080A^4 \gamma \beta_2 \delta \\ &+720A^3 \delta \beta_1 \sigma_2 \beta_2 + 6480A^6 \delta^2 \beta_2 \gamma - 2340A^5 \delta \beta_2^2 \gamma - 6480A^5 \delta^2 \gamma \sigma_2 \\ &-720A^3 \delta \beta_1 \beta_2 + 14580A^7 \delta^3 \gamma + 1440A^5 \delta \gamma^3 + 6480A^5 \gamma \delta^2 \right) e^{899} \\ &+\frac{243A^8 \delta^3 \gamma}{141(1+\sigma_1)^4} \cos^2(x) - \frac{A\cos^6(x)}{6720(1+\sigma_1)^4} \left(-90720A^6 \delta^2 \beta_2 \gamma - 25515A^6 \delta^3 \beta_1\right) \\ &-\frac{A}{6720(1+\sigma_1)^4} \left(720A^5 \delta \beta_2^2 \gamma + 45A^2 \delta \beta_1 \sigma_2^2 + 405A^6 \delta^3 \beta_1 + 720A^4 \gamma \beta_2 \delta \\ &-720A^4 \delta \beta_2 \gamma \sigma_2 + 90A^3 \delta \beta_1 \beta_2 + 45A^4 \delta \beta_1 \beta_2^2 + 2160A^6 \delta^2 \beta_2 \gamma \\ &+ 270A^5 \delta^2 \beta_1 \beta_2 + 45A^2 \delta \beta_1 - 90A^3 \delta \beta_1 \sigma_2 \beta_2 + 270A^4 \delta^2 \beta_1 \\ &- 270A^4 \delta^2 \beta_1 \sigma_2 - 90A^2 \delta \beta_1 \sigma_2 + 180A^4 \delta \beta_1 \gamma^2 \right) \right] \sin(x) \\ &-\frac{A\cos^5(x)}{6720(1+\sigma_1)^4} \left(204120A^8 \delta^4 - 7200A^4 \delta \beta_1 \gamma \beta_2 + 52245A^6 \delta^3 + 2025\delta^2 A^4 \\ &- 52245A^6 \delta^3 \sigma_2 + 3600A^4 \gamma^2 \delta - 80055A^6 \delta^2 \beta_2^2 - 3600A^5 \delta \beta_2^2 \\ &+ 3600\delta \beta_2 \gamma^2 A^5 + 52245A^7 \delta^3 \beta_2 + 82080\delta^2 \gamma^2 A^6 + 2025A^4 \delta^2 \sigma_2^2 \\ &+ 3600A^4 \delta \beta_2^2 \sigma_2 - 4050A^5 \delta^2 \beta_2 \sigma_2 - 4050A^4 \delta^2 \sigma_2 - 2025A^4 \delta^2 \beta_1^2 \\ &+ 4050A^5 \delta^2 \beta_2 - 3600A^4 \delta \beta_2^2 - 3600A^4 \delta \gamma^2 \sigma_2 \right) - \frac{243A^8 \delta^3 \beta_2 \cos^3(x)}{141(1+\sigma_1)^4} \\ &-\frac{A\cos^4(x)}{6720(1+\sigma_1)^4} \left(25920\delta \beta_2 \gamma^2 A^5 + 1440A^3 \gamma \beta_1 \delta - 720A^3 \delta \beta_1^2 \beta_2 + 33480A^5 \delta^2 \beta_2 \\ &- 1440A^4 \delta \beta_2^2 \sigma_2 - 1440A^3 \delta \beta_2 \sigma_2 + 720A^3 \delta \beta_1^2 \beta_2 + 33480A^6 \delta^2 \beta_2^2 \\ &- 1440A^3 \delta \beta_1 \gamma \sigma_2 - 7920A^5 \delta \beta_2^3 + 1440A^4 \delta \beta_1 \gamma \beta_2 + 2700A^3 \delta \beta_1^2 \beta_2 \\ &- 3240\delta^2 A^4 + 7920A^4 \delta \beta_1 \gamma \beta_2 - 70470A^8 \delta^4 - 32805A^6 \delta^3 \\ &- 3240\delta^2 A^4 + 7920A^4 \delta \beta_1 \gamma \beta_2 - 70470A^8 \delta^4 - 32805A^6 \delta^3 \\ &- 4905A^4 \delta \beta_2^2 \sigma_2 - 135A^2 \delta \sigma_2^2 + 135A^2 \delta \beta_1^2 + 270A^3 \delta \beta_2 \sigma_2 + 45A^2 \delta \sigma_3^2 - 45A^2 \delta \sigma_3 \\ &- 32805A^7 \delta^3 \beta_2 - 45900\delta^2 \gamma^2 A^6 + 32805A^6 \delta^3 \sigma_2 + 37800A^6 \delta^2 \beta_2^2 \\ &- 135A^3 \delta \beta_1 \delta_2 - 25900\delta^2 \gamma^2 A^6 + 32805A^6 \delta^3 \sigma_2 + 37800A^6 \delta^2 \beta_2^2 \\ &- 32805A^7 \delta^3 \beta_2 - 45900\delta^2 \gamma^2 A^6 + 32805A^6 \delta^3 \sigma_2 + 37800A^6 \delta^2 \beta_2^2 \\ &- 32805A^7 \delta^3 \beta_2 - 45900\delta^2 \gamma^2 A^6 +$$

$$\begin{split} &-1356 \rho_2 A^3 + 4995 A^5 \delta \beta_2^3 - 135 A^2 \delta \beta_1^2 \sigma_2 - 6480 A^5 \delta^2 \rho_2 - 5040 \delta \rho_2 \eta^2 A^5 \\ &+ 4905 A^4 \delta \beta_2^2 - 5040 A^4 \gamma^2 \delta - 3240 A^4 \delta^2 \sigma_2^2 \Big) - \frac{6561 A^9 \delta^4 \cos^9(x)}{4481(1 + \sigma_1)^4} \\ &- \frac{A\cos^2(x)}{6720(1 + \sigma_1)^4} \Big(540 A^3 \delta \beta_1^2 \beta_2 + 1980 A^5 \delta \beta_2^3 - 9180 A^5 \delta^2 \beta_1 \gamma - 9006 \beta_2 A^3 \\ &+ 1800 A^4 \delta \beta_2^2 \sigma_2 - 15120 A^5 \delta^2 \beta_2 - 10080 \delta \beta_2 \eta^2 A^5 - 15120 A^6 \delta^2 \beta_2^2 \\ &+ 1440 A^3 \delta \beta_1 \gamma \sigma_2 - 1440 A^3 \gamma \beta_1 \delta - 37260 A^7 \delta^3 \beta_2 - 1440 A^4 \delta \beta_1 \gamma \beta_2 \\ &+ 15120 A^5 \delta^2 \beta_2 \sigma_2 + 1800 A^3 \delta \beta_2 \sigma_2 - 900 A^3 \delta \beta_2 \sigma_2^2 - 1800 A^4 \delta \beta_2^2 \Big) \\ &- \frac{A}{6720(1 + \sigma_1)^4} \Big(180 A^4 \delta \beta_1 \gamma \beta_2 + 360 A^4 \delta \beta_2^2 + 180 \delta \beta_2 A^3 - 360 A^4 \delta \beta_2^2 \sigma_2 \\ &+ 1080 A^6 \delta^2 \beta_2^2 + 1080 A^5 \delta^2 \beta_2 + 1620 A^7 \delta^3 \beta_2 - 180 A^3 \delta \beta_1 \gamma \sigma_2 \\ &+ 180 A^3 \gamma \beta_1 \delta + 180 A^3 \delta \beta_2 \sigma_2^2 + 720 \delta \beta_2 \Lambda^2 A^5 \Big) \\ &- \frac{A\cos^2(x)}{6720(1 + \sigma_1)^4} \Big(45360 A^6 \delta^2 \beta_2^2 + 25515 A^6 \delta^3 \sigma_2 - 25515 A^6 \delta^3 - 45360 \delta^2 \gamma^2 A^6 \\ &- 240570 A^8 \delta^4 - 25515 A^7 \delta^3 \beta_2 \Big) \\ &- \frac{A\cos(x)}{6720(1 + \sigma_1)^4} \Big(-270 A^4 \delta^2 \beta_1^2 - 2430 A^4 \delta^2 \sigma_2 - 1620 A^4 \delta \gamma^2 \sigma_2 - 90 A^3 \delta \beta_1^2 \beta_2 + 1215 \delta^2 A^4 \\ &- 1440 A^4 \delta \beta_1 \gamma \beta_2 + 8505 A^8 \delta^4 + 6075 A^6 \delta^3 + 1305 A^4 \delta \beta_2^2 \sigma_2 \\ &+ 90 A^2 \delta \beta_1^2 - 270 A^3 \delta \beta_2 \sigma_2 - 2430 A^5 \delta^2 \beta_2 \sigma_2 + 135 A^3 \delta \beta_2 \sigma_2^2 \\ &+ 135 A^2 \delta \sigma_2^2 - 135 A^2 \sigma_2 \delta - 45 A^2 \delta \sigma_2^3 + 45 A^2 \delta + 6075 A^7 \delta^3 \beta_2 \\ &+ 8100 \delta^2 \gamma^2 A^6 - 6075 A^6 \delta^3 \sigma_2 - 3105 A^6 \delta^2 \beta_2^2 + 135 \delta \beta_2 A^3 \\ &- 1395 A^4 \delta \beta_2^2 + 1620 A^4 \gamma^2 \delta + 1215 A^4 \delta^2 \sigma_2^2 \Big) \Big) t^8 \\ + \left\{ -\frac{A\cos^5(x)}{6720(1 + \sigma_1)^4} \Big(-2876 A^5 \delta^2 \beta_2 - 99792 A^5 \gamma \delta^2) \\ &- \frac{A\cos^5(x)}{(6720(1 + \sigma_1)^4} \Big(1008 A^3 \delta \beta_1 \beta_2 + 1344 A^2 \gamma^2 \beta_1 - 1344 A^2 \gamma \beta_2 \sigma_2 - 1008 A^2 \delta \beta_1 \sigma_2 \\ - \frac{A\cos^2(x)}{(1 + \sigma_1)^4} \Big(1008 A^3 \delta \beta_1 \beta_2 + 1344 A^2 \gamma^2 \beta_1 - 1344 A^2 \gamma \beta_2 \sigma_2 - 1008 A^2 \delta \beta_1 \sigma_2 \\ - \frac{A\cos^2(x)}{(6720(1 + \sigma_1)^4} \Big(1008 A^3 \delta \beta_1 \beta_2 + 1344 A^2 \gamma^2 \beta_1 - 1344 A^2 \gamma \beta_2 \sigma_1 - 1008 A^2 \delta \beta_1 \sigma_2 \\ - \frac{A\cos^2(x)}{(6720(1 + \sigma_1)^4} \Big) \Big(1008 A^3 \delta \beta_1 \beta_2 + 1344 A^2 \gamma^2 \beta_1 - 1344 A^2 \gamma \beta_2 \sigma_1 - 100$$

$$+ 1344A^{3}\gamma\beta_{2}^{2} + 1344A^{2}\gamma\beta_{2}\sigma_{1} + 1344A^{3}\gamma\beta_{2}^{2}\sigma_{1} + 1344A^{2}\gamma\beta_{2}$$

$$- 1008A^{2}\delta\beta_{1}\sigma_{1}\sigma_{2} + 1344A^{2}\gamma^{2}\beta_{1}\sigma_{1} + 12096A^{4}\delta^{2}\beta_{1}$$

$$+ 1008A^{3}\delta\beta_{1}\sigma_{1}\beta_{2} + 12096A^{4}\delta^{2}\beta_{1}\sigma_{1} - 1344A^{2}\gamma\beta_{2}\sigma_{2}$$

$$+ 40320A^{4}\gamma\beta_{2}\delta\sigma_{1} + 40320A^{4}\gamma\beta_{2}\delta + 1008A^{2}\delta\beta_{1}$$

$$+ 1008A^{2}\delta\beta_{1}\sigma_{1})$$

$$- \frac{A}{6720(1+\sigma_{1})^{4}} \Big(-448A^{2}\gamma\beta_{2} + 448A^{2}\gamma\beta_{2}\sigma_{2}\sigma_{1} - 448A^{2}\gamma\beta_{2}\sigma_{1} - 448A^{3}\gamma\beta_{2}^{2}\sigma_{1}$$

$$+ 448A^{2}\gamma\beta_{2}\sigma_{2} + 336A^{2}\delta\beta_{1}\sigma_{2} - 1008A^{4}\delta^{2}\beta_{1} - 336A^{3}\delta\beta_{1}\beta_{2} - 336A^{2}\delta\beta_{1} - 336A^{3}\delta\beta_{1}\sigma_{1}\beta_{2} + 336A^{2}\delta\beta_{1}\sigma_{1}\sigma_{2} - 336A^{2}\delta\beta_{1}\sigma_{1} - 4032A^{4}\gamma\beta_{2}\delta - 224A^{2}\gamma^{2}\beta_{1} - 448A^{3}\gamma\beta_{2}^{2} - 4032A^{4}\gamma\beta_{2}\delta\sigma_{1} - 224A^{2}\gamma^{2}\beta_{1}\sigma_{1} - 1008A^{4}\delta^{2}\beta_{1}\sigma_{1} \Big)$$

$$-\frac{A\cos(x)}{6720(1+\sigma_{1})^{4}} \Big(2688A^{3}\delta\beta_{1}\beta_{2} + 1680A^{3}\gamma\beta_{2}^{2}\sigma_{1} - 1792A^{3}\gamma^{3} + 2688A^{3}\delta\beta_{1}\sigma_{1}\beta_{2} + 6720A^{3}\gamma\delta\sigma_{2}\sigma_{1} - 6720A^{3}\gamma\delta\sigma_{1} + 6720A^{3}\gamma\delta\sigma_{2} - 31248A^{5}\gamma\delta^{2} - 224A^{2}\gamma\beta_{2}\sigma_{1} + 224A^{2}\gamma\beta_{2}\sigma_{2}\sigma_{1} + 1680A^{3}\gamma\beta_{2}^{2} + 224A^{2}\gamma\beta_{2}\sigma_{2} - 1792A^{3}\gamma^{3}\sigma_{1} - 6720A^{3}\gamma\delta - 224A^{2}\gamma\beta_{2} - 112A\gamma + 112A\gamma\beta_{1}^{2} + 112A\gamma\beta_{1}^{2}\sigma_{1} - 112A\gamma\sigma_{2}^{2}\sigma_{1} + 224A\gamma\sigma_{2}\sigma_{1} - 112A\gamma\sigma_{2}^{2} + 224A\gamma\sigma_{2} - 112A\gamma\sigma_{1} - 6720A^{4}\gamma\beta_{2}\delta\sigma_{1} - 31248A^{5}\gamma\delta^{2}\sigma_{1} - 6720A^{4}\gamma\beta_{2}\delta\Big) \Big)$$

$$-\frac{A\cos^{3}(x)}{6720/(1+\sigma_{1})^{4}} \Big(-3584A^{3}\gamma\beta_{2}^{2}\sigma_{1}+9408A^{4}\gamma\beta_{2}\delta\sigma_{1}+118944A^{5}\gamma\delta^{2}+9408A^{4}\gamma\beta_{2}\delta\sigma_{1}$$

$$+3584A^{3}\gamma^{3}\sigma_{1}+9408A^{3}\gamma\delta+9408A^{3}\gamma\delta\sigma_{1}+118944A^{5}\gamma\delta^{2}\sigma_{1}$$

$$-3584A^{3}\gamma\beta_{2}^{2}-5376A^{3}\delta\beta_{1}\beta_{2}+3584A^{3}\gamma^{3}-9408A^{3}\gamma\delta\sigma_{2}$$

$$-5376A^{3}\delta\beta_{1}\sigma_{1}\beta_{2}-9408A^{3}\gamma\delta\sigma_{2}\sigma_{1}\Big)$$

$$-\frac{A\cos^{4}(x)}{6720(1+\sigma_{1})^{4}} \Big(-15120A^{4}\delta^{2}\beta_{1}-47040A^{4}\gamma\beta_{2}\delta-15120A^{4}\delta^{2}\beta_{1}\sigma_{1}$$

$$-47040A^{4}\gamma\beta_{2}\delta\sigma_{1}\Big)\Big]\sin(x)$$

$$-\frac{A\cos^{7}(x)}{6720(1+\sigma_{1})^{4}} \Big(95256A^{6}\delta^{3}\sigma_{1}+95256A^{6}\delta^{3}\Big)$$

$$\begin{split} &-\frac{A\cos^3(x)}{6720(1+\sigma_1)^4}\Big(-24192A^4\delta^2\sigma_2-24192A^4\delta^2\sigma_1\sigma_2+24192\delta^2A^4-1344A^2\gamma\beta_1\beta_2\sigma_1\\ &+122472A^6\delta^3+1008A^3\beta_2\delta\sigma_1-1344A^2\gamma\beta_1\beta_2-18312A^4\delta\beta_2^2\sigma_1\\ &+1344A^3\gamma^2\beta_2\sigma_1-504A^2\delta\beta_1^2+1344A^2\gamma^2\sigma_1-1008A^3\delta\beta_2\sigma_2\\ &+24192A^5\delta^2\beta_2\sigma_1+504A^2\delta\sigma_1+24192A^4\delta^2\sigma_1+54024A^4\gamma^2\delta\sigma_1\\ &-1344A^2\gamma^2\sigma_2+504A^2\delta\sigma_2-1008A^2\sigma_2\delta+1008A^2\sigma_2\delta\sigma_1\\ &+504A^2\delta-1008A^3\delta\beta_2\sigma_1\sigma_2-1344A^2\gamma^2\sigma_2\sigma_1-122472A^6\delta^3\sigma_1\\ &+1008\delta\beta_2A^3+1344A^2\gamma^2+504A^2\delta\sigma_1\sigma_2^2-504A^2\delta\beta_1^2\sigma_1\\ &+24192A^5\delta^2\beta_2-18312A^4\delta\beta_2^2+45024A^4\gamma^2\delta+1344A^3\gamma^2\beta_2\Big) \end{split}$$

$$-\frac{A\cos^5(x)}{6720(1+\sigma_1)^4}\Big(-195048A^6\delta^3+15120A^4\delta^2\sigma_1\sigma_2+15120A^4\delta^2\sigma_2-195048A^6\delta^3\sigma_1\\ &-15120A^5\delta^2\beta_2+13440A^4\delta\beta_2^2-33600A^4\gamma^2\delta\sigma_1\\ &-33600A^4\gamma^2\delta-15120\delta^2A^4-15120A^4\delta^2\sigma_1-15120A^5\delta^2\beta_2\sigma_1\\ &+13440A^4\delta\beta_2^2\sigma_1\Big) \cr$$

$$-\frac{A}{6720(1+\sigma_1)^4}\Big(-112A\gamma\beta_1\sigma_1-1344\delta\beta_2A^3-112A^2\gamma\beta_1\beta_2-1008A^3\gamma\beta_1\delta\sigma_1\\ &+112A\gamma\beta_1\sigma_2\sigma_1+1344A^3\delta\beta_2\sigma_1\sigma_2-896A^3\gamma^2\beta_2-112A^2\gamma\beta_1\beta_2\sigma_1\\ &-1344A^3\beta_2\delta\sigma_1-1344A^4\delta\beta_2^2\sigma_1-1008A^3\gamma\beta_1\delta-112A\gamma\beta_1\\ &+112A\gamma\beta_1\sigma_2-896A^3\gamma^2\beta_2\sigma_1-4032A^5\delta^2\beta_2-1344A^4\delta\beta_2^2\\ &+1344A^3\delta\beta_2\sigma_2-4032A^5\delta^2\beta_2\sigma_1\Big) \cr$$

$$-\frac{A\cos^4(x)}{6720(1+\sigma_1)^4}\Big(-7168A^3\gamma^2\beta_2-5376A^4\delta\beta_2\sigma_1-5376\delta\beta_2A^3-124992A^5\delta^2\beta_2\sigma_1\\ &-5376A^4\delta\beta_2^2+5376A^3\delta\beta_2\delta\sigma_1+896A^2\gamma\beta_1\beta_2+4892A^4\delta\beta_2\sigma_1\\ &-5376A^4\delta\beta_2^2+5376A^3\delta\beta_2\sigma_1+896A^2\gamma\beta_1\beta_2+4892A^4\delta\beta_2\sigma_1\\ &-22680A^6\delta^3-1008A^3\beta_2\delta\sigma_1+896A^2\gamma\beta_1\beta_2+4872A^4\delta\beta_2^2\sigma_1\\ &-1120A^3\gamma^2\beta_2\sigma_1-504A^2\delta\sigma_1-9072A^4\delta^2\sigma_1-1344A^4\gamma^2\sigma_1\\ &-112A^3\gamma_1\beta_2\sigma_1-5376A^3\beta_2\delta\sigma_1+896A^2\gamma\beta_1\beta_2+4872A^4\delta\beta_2\sigma_1\\ &-1120A^3\gamma_1\beta_2\sigma_1+536A^2\delta\beta_1^2-1120A^2\gamma^2\sigma_1+1008A^3\delta\beta_2\sigma_2\\ &-9072A^5\delta^2\beta_2\sigma_1-504A^2\delta\sigma_1-9072A^4\delta^2\sigma_1-9072A^4\delta^2\sigma_1-1344A^4\gamma^2\delta\sigma_1\\ \cr$$

$$\begin{split} &+ 1120A^{2}p^{2}\sigma_{2} - 504A^{2}\delta\sigma_{2}^{2} + 1008A^{2}\sigma_{2}\delta + 1008A^{2}\sigma_{2}\delta\sigma_{1} \\ &- 504A^{2}\delta + 1008A^{3}\delta\beta_{2}\sigma_{1}\sigma_{2} + 1120A^{2}\gamma^{2}\sigma_{2}\sigma_{1} - 22680A^{6}\delta^{3}\sigma_{1} \\ &- 9072A^{5}\delta^{2}\beta_{2} + 4872A^{4}\delta\beta_{2}^{2} - 13440A^{4}\gamma^{2}\delta - 1120A^{3}\gamma^{2}\beta_{2} \Big) \\ &- \frac{A\cos^{2}(x)}{6720(1+\sigma_{1})^{4}} \Big(- 224A\gamma\beta_{1}\sigma_{2}\sigma_{1} + 224A\gamma\beta_{1}\sigma_{1} + 224A^{2}\gamma\beta_{1}\beta_{2} + 56448A^{5}\delta^{2}\beta_{2} \\ &+ 6720A^{4}\delta\beta_{2}^{2} + 9072A^{3}\gamma\beta_{1}\delta\sigma_{1} + 224A\gamma\gamma_{1} + 6720\delta\beta_{2}A^{3} \\ &+ 9072A^{3}\gamma\beta_{1}\delta + 224A^{2}\gamma\beta_{1}\beta_{2}\sigma_{1} + 7168A^{3}\gamma^{2}\beta_{2} + 56448A^{5}\delta^{2}\beta_{2}\sigma_{1} \\ &+ 6720A^{3}\beta_{2}\delta\sigma_{1} + 7168A^{3}\gamma^{2}\beta_{2}\sigma_{1} - 6720A^{3}\delta\beta_{2}\sigma_{1}\sigma_{2} \\ &- 224A\gamma\beta_{1}\sigma_{2} + 6720A^{4}\delta\beta_{2}^{2}\sigma_{1} - 6720A^{3}\delta\beta_{2}\sigma_{2} \Big) \bigg\} t^{6} \\ + \Bigg\{ -\frac{A\cos^{3}(x)}{6720(1+\sigma_{1})^{4}} \Big(-1008A^{2}\beta_{2}\beta_{1}a\sigma_{1}^{2} - 4032A^{2}\beta_{2}^{2}\sigma_{1} - 1008A^{2}\beta_{2}\beta_{1}a - 2016A^{2}\beta_{2}\beta_{1}a\sigma_{1} \\ &- 2016A^{2}\beta_{2}^{2}-2016A^{2}\beta_{2}\sigma_{1}^{2} \Big) \\ - \frac{A\cos(x)}{6720(1+\sigma_{1})^{4}} \Big(1344A^{2}\beta_{2}^{2}\sigma_{1}^{2} + 672A^{2}a\beta_{2}\beta_{1}a + 1344A^{2}\beta_{2}^{2} + 2688A^{2}\beta_{2}^{2}\sigma_{1} \\ &+ 1344A^{2}\beta_{2}\beta_{1}a\sigma_{1} + 672A^{2}\beta_{2}\beta_{1}a\sigma_{1}^{2} \Big) \\ - \frac{A}{6720(1+\sigma_{1})^{4}} \Big(-504\delta\beta_{2}A^{3} - 336A^{2}\beta_{2}^{2}\sigma_{1} - 168A\beta_{2}\sigma_{1}^{2} + 336A\beta_{2}\sigma_{2}\sigma_{1} \\ &+ 168A\beta_{2}\sigma_{2} - 1008A^{3}\beta_{2}\delta\sigma_{1} - 168\beta_{2}A - 504A^{3}\beta_{2}\delta\sigma_{1}^{2} \\ &- 168A^{2}\beta_{2} - 672A^{3}\beta_{2}\gamma\alpha_{1}^{2} - 1344A^{3}\beta_{2}\gamma\alpha_{1} \Big) \\ - \frac{A\cos^{2}(x)}{6720(1+\sigma_{1})^{4}} \Big(5544A^{3}\beta_{2}\delta\sigma_{1}^{2} + 5376A^{3}\beta_{2}\gamma\alpha_{1} - 168A^{2}\beta_{2}^{2}\sigma_{1} - 336A\beta_{2}\sigma_{2} \\ &+ 336\beta_{2}A + 672A^{2}\beta_{2}^{2}\sigma_{1} + 336A^{2}\beta_{2}^{2} + 5376A^{3}\beta_{2}\gamma\alpha_{1} \\ &+ 672A\beta_{2}\sigma_{1} - 336A\beta_{2}\sigma_{2}^{2} + 5376A^{3}\beta_{2}\gamma\alpha_{1} \\ &+ 336A^{2}\beta_{2}\sigma_{1}^{2} + 5544\delta\beta_{2}A^{3} + 11088A^{3}\beta_{2}\delta\sigma_{1} \Big) \\ + \Big[-\frac{A\cos(x)}{(6720(1+\sigma_{1})^{4}} \Big(672\beta_{2}A\beta_{1}\sigma_{1} + 336\beta_{2}A\beta_{1}\sigma_{1}^{2} + 2688A^{3}\beta_{2}^{2}\alpha_{1} \\ &+ 336A^{2}\beta_{2}^{2}\alpha_{1} + 336\beta_{2}A\beta_{1}\sigma_{1}^{2} + 2688A^{3}\beta_{2}^{2}\alpha_{1} \Big) \\ \end{cases}$$

$$\begin{split} &-\frac{A\cos^2(x)}{6720(1+\sigma_1)^4} \Big(2016A^2\gamma\beta_2 + 24192A^4\beta_2\delta a\sigma_1 + 2016A^2\gamma\beta_2\sigma_1^2 + 2016A^2\beta_2\sigma_1a \\&+ 1008A^3\beta_2^2a + 12096A^4\beta_2\delta a\sigma_1^2 + 1008A^3\beta_2^2a\sigma_1^2 + 4032A^2\gamma\beta_2\sigma_1 \\&- 1008A^2\beta_2\sigma_2a\sigma_1^2 + 2016A^3\beta_2^2a\sigma_1 - 2016A^2\beta_2\sigma_2a\sigma_1 \\&+ 12096A^4\beta_2\delta a + 1008A^2\beta_2\sigma_1^2a - 1008A^2\beta_2\sigma_2a - 1008A^2\beta_2\sigma_2a \Big) \\&- \frac{A}{6720(1+\sigma_1)^4} \Big(- 336A^2\gamma\beta_2 - 336A^2\beta_2a - 2016A^4\beta_2\delta a\sigma_1 - 1008A^4\beta_2\delta a\sigma_1^2 \\&- 1008A^4\beta_2\delta a - 336A^2\beta_2\sigma_1^2a - 672A^2\gamma\beta_2\sigma_1 - 336A^2\beta_2\sigma_2a \\&- 336A^3\beta_2a\sigma_1^2 + 672A^2\beta_2\sigma_2a\sigma_1 + 336A^2\beta_2\sigma_2a\sigma_1^2 - 336A^3\beta_2a \\&- 672A^3\beta_2a\sigma_1 - 336A^2\gamma\beta_2\sigma_1^2 - 672A^2\beta_2\sigma_1a \Big) \\&- \frac{A\cos^4(x)}{6720(1+\sigma_1)^4} \Big(-30240A^4\beta_2\delta a\sigma_1 - 15120A^4\beta_2\delta a - 15120A^4\beta_2\delta a\sigma_1^2 \Big) \\&- \frac{A\cos^4(x)}{6720(1+\sigma_1)^4} \Big(-5376A^3\beta_2a\sigma_1^2 - 5376A^3\beta_2a - 10752A^3\beta_2a\sigma_1 \Big) \Big] \sin(x) \\&- \frac{A\cos^4(x)}{6720(1+\sigma_1)^4} \Big(-5376A^3\beta_2\gamma a - 12096A^3\beta_2\delta \sigma_1 - 6048\delta\beta_2A^3 - 6048A^3\beta_2\delta \sigma_1^2 \\&- 10752A^3\beta_2\gamma a\sigma_1 - 5376A^3\beta_2\gamma a\sigma_1^2 \Big) \Big\} t^5 \\&+ \Big\{ -\frac{\cos^2(x)}{6720(1+\sigma_1)^4} \Big(5600A\beta_2\sigma_2\sigma_1 - 26040A^3\beta_2\delta \sigma_1^2 - 5600A\beta_2\sigma_1 - 2800A^2\beta_2^2 \\&- 2800A^2\beta_2\sigma_1^2 - 2800\beta_2A - 52080A^3\beta_2\delta \sigma_1 + 2800A\beta_2\sigma_2\sigma_1^2 \\&- 2800A\beta_2\sigma_1^2 - 1120A\gamma\beta_1\sigma_1^2 + 2800A\beta_2\sigma_2 \Big) \\&- \frac{A\cos(x)}{6720(1+\sigma_1)^4} \Big(280 + 560\sigma_1 - 560\sigma_2 + 22680\delta^2A^4 - 560\sigma_2\sigma_1^2 + 560\beta_2A \\&+ 1120A\beta_2\sigma_1 - 560A\beta_2\sigma_2 + 560A\beta_2\sigma_1^2 - 1960A^2\beta_2\sigma_1^2 + 560A\beta_2\sigma_1^2 \\&+ 16800A^3\beta_2\delta\sigma_1 + 5600A^2\gamma_2\sigma_1^2 + 560\sigma_2\sigma_1 + 2800A^2\beta_2\sigma_1^2 \\&+ 280\sigma_2^2 + 11200A^2\gamma_2\sigma_1 + 16800A^2\delta_2\sigma_1^2 - 1960A^2\beta_2\sigma_1 + 2800A^2\delta\sigma_1^2 \\&+ 8400A^2\sigma_2\delta - 8400A^2\sigma_2\delta\sigma_1^2 - 16800A^2\sigma_2\delta\sigma_1 + 2800A^2\sigma_1^2 + 2800A^2\sigma_1^2 \\&+ 8400A^2\sigma_2\delta - 8400A^2\sigma_2\sigma_1 - 16800A^2\sigma_2\delta_2\sigma_1 - 1960A^2\gamma_2^2 \\&- 8400A^2\sigma_2\delta - 8400A^2\sigma_2\sigma_1^2 - 1960A^2\gamma_2^2 \\&- 1960A^2\beta_2^2 - 1120A\beta_2\sigma_2\sigma_1 - 3920A^2\beta_2\sigma_1 \Big) \Big\}$$

$$\begin{split} &-\frac{A\cos^4(x)}{6720(1+\sigma_1)^4} \Big(23520A^3\beta_2\delta\sigma_1^2 + 47040A^3\beta_2\delta\sigma_1 + 23520\delta\beta_2A^3 \Big) \\ &-\frac{A\cos^{3}(x)}{6720(1+\sigma_1)^4} \Big(10080A^2\sigma_2\delta - 6720A^2\gamma^2 - 13440A^2\gamma^2\sigma_1 + 10080A^2\sigma_2\delta\sigma_1^2 \\ &+ 3360A^2\beta_2^2 - 10080A^2\delta\sigma_1^2 - 60480\delta^2A^4 - 10080A^2\delta \\ &- 6720A^2\gamma^2\sigma_1^2 - 120960A^4\delta^2\sigma_1 - 20160A^3\beta_2\delta\sigma_1 - 10080\delta\beta_2A^3 \\ &- 20160A^2\delta\sigma_1 + 20160A^2\sigma_2\delta\sigma_1 + 6720A^2\beta_2^2\sigma_1 + 3360A^2\beta_2\sigma_1^2 \\ &- 60480A^4\delta^2\sigma_1^2 - 10080A^3\beta_2\delta\sigma_1^2 \Big) \\ &+ \Big[-\frac{A\cos(x)}{6720(1+\sigma_1)^4} \Big(-560\beta_2A\beta_1 + 23520A^3\gamma\delta\sigma_1^2 + 23520A^3\gamma\delta - 1120\beta_2A\beta_1\sigma_1 \\ &+ 3360A\gamma\sigma_1^2 - 3360A\gamma\sigma_2\sigma_1^2 + 6720A\gamma\sigma_1 - 560\beta_2A\beta_1\sigma_1^2 \\ &+ 47040A^3\gamma\delta\sigma_1 + 3360A^2\gamma\beta_2\sigma_1^2 - 3360A\gamma\sigma_2 + 3360A^2\gamma\beta_2 \\ &+ 6720A^2\gamma\beta_2\sigma_1 - 6720A\gamma\sigma_2\sigma_1 + 3360A\gamma \Big) \\ &- \frac{A}{6720(1+\sigma_1)^4} \Big(560\beta_1\sigma_1 - 280\sigma_2\beta_1 - 280\sigma_2\beta_1\sigma_1^2 - 560\sigma_2\beta_1\sigma_1 + 280\beta_1 \\ &+ 1680A^2\delta\beta_1\sigma_1 + 280\beta_2A\beta_1\sigma_1^2 + 560A^2\gamma\beta_2\sigma_1 + 840A^2\delta\beta_1\sigma_1^2 \\ &+ 280\beta_1\sigma_1^2 + 2800A^2\gamma\beta_2\sigma_1^2 + 5600A^2\gamma\beta_2\sigma_1 + 840A^2\delta\beta_1\sigma_1^2 \\ &+ 280\beta_2A\beta_1 + 840A^2\delta\beta_1 \Big) \\ &- \frac{A\cos^2(x)}{6720(1+\sigma_1)^4} \Big(-10080A^2\gamma\beta_2 - 20160A^2\gamma\beta_2\sigma_1 - 5040A^2\delta\beta_1\sigma_1 - 2520A^2\delta\beta_1\sigma_1^2 \\ &- 10080A^2\gamma\beta_2\sigma_1^2 - 2520A^2\delta\beta_1 \Big) \Big] \sin(x) \\ &- \frac{A\cos^5(x)}{6720(1+\sigma_1)^4} \Big(37800\delta^2A^4 + 37800A^4\delta^2\sigma_1^2 + 75600A^4\delta^2\sigma_1 \Big) \\ &- \frac{A\cos^5(x)}{6720(1+\sigma_1)^4} \Big(4200A^3\beta_2\delta\sigma_1^2 + 2800A^2\beta_2\sigma_1 - 1400A\beta_2\sigma_2\sigma_1^2 + 1400A\beta_2\sigma_1^2 \\ &+ 2800A\beta_2\sigma_1 + 1120A\gamma\beta_1\sigma_1 + 560A\gamma\beta_1 \Big) \Big\} t^4 \end{split}$$

$$\begin{split} &+ \left\{ \left[-\frac{A\cos(x)}{6720(1+\sigma_1)^4} \left(2240\Omega A_{\rm Y} + 4480\Omega A_{\rm Y}\sigma_1 + 2240\Omega A_{\rm Y}\sigma_1^2 \right) \right. \\ &- \frac{A\cos^2(x)}{6720(1+\sigma_1)^4} \left(-6720A^2\beta_2\sigma_1^2\alpha - 20160A^2\beta_2\sigma_1^2\alpha - 6720A^2\beta_2\alpha - 20160A^2\beta_2\sigma_1\alpha \right) \right. \\ &- \frac{A}{6720(1+\sigma_1)^4} \left(6720A^2\beta_2\sigma_1\alpha + 1120\Omega\beta_1\sigma_1^2 + 1120\beta_1\sigma_1^3 + 1120\beta_1 \right. \\ &+ 3360\beta_1\sigma_1^2 + 2240\Omega\beta_1\sigma_1 + 3360\beta_1\sigma_1 + 1120\Omega\beta_1 \\ &+ 6720A^2\beta_2\sigma_1^2\alpha + 2240A^2\beta_2\sigma_1^3\alpha + 2240A^2\beta_2\alpha \right) \right] \sin(x) \\ &- \frac{A\cos^2(x)}{6720(1+\sigma_1)^4} \left(-4480\Omega A\beta_2\sigma_1 - 2240\Omega A\beta_2\sigma_1^2 - 2240\beta_2A - 6720A\beta_2\sigma_1^2 \right. \\ &- 13440A\beta_1\alpha\sigma_1^2 - 2240A\beta_2\sigma_1^3 - 4480A\beta_1\alpha - 2240\Omega A\beta_2\sigma_1^2 \right. \\ &- 13440A\beta_1\alpha\sigma_1 - 6720A\beta_2\sigma_1^2 - 4480A\beta_1\alpha\sigma_1^3 \right) \\ &- \frac{A}{6720(1+\sigma_1)^4} \left(1120\beta_2A + 1120\Omega A\beta_2\sigma_1^2 + 1120A\beta_2\sigma_1^3 + 2240A\beta_1\alpha \right. \\ &+ 2240A\beta_1\alpha\sigma_1^3 + 1120\Omega A\beta_2\sigma_1 + 6720A\beta_1\alpha\sigma_1^2 \right) \\ &- \frac{A\cos(x)}{6720(1+\sigma_1)^4} \left(3360\Omega A^2\delta + 1120\Omega A\beta_2\sigma_1 + 6720A\beta_1\alpha\sigma_1^2 \right) \\ &- \frac{A\cos(x)}{6720(1+\sigma_1)^4} \left(3360\Omega A^2\delta\sigma_1 - 3360\Omega A^2\delta\sigma_1^2 - 3360\Omega A^2\delta\sigma_1 \right) \right\} t^3 \\ &+ \left\{ -\frac{A\cos^3(x)}{6720(1+\sigma_1)^4} \left(-6720\Omega A^2\delta\sigma_1 - 3360\Omega A^2\delta\sigma_1^2 - 3360\Omega A^2\delta\sigma_1 \right) \right\} t^3 \\ &+ \left\{ -\frac{A\cos^3(x)}{6720(1+\sigma_1)^4} \left(-10080A\beta_2\sigma_1^2 - 3360\beta_2A - 3360A\beta_2\sigma_1^3 - 10080A\beta_2\sigma_1 \right) \right\} t^3 \\ &+ \left\{ -\frac{A\cos(x)}{6720(1+\sigma_1)^4} \left(-10080A\beta_2\sigma_1^2 - 3360\beta_1 - 3360A\beta_2\sigma_1^3 - 10080A\beta_2\sigma_1 \right) \right\} t^3 \\ &+ \left\{ -\frac{A\cos(x)}{6720(1+\sigma_1)^4} \left(-20160A\gamma\sigma_1 - 6720A\gamma\sigma_1^3 - 6720A\gamma - 20160A\gamma\sigma_1^2 \right) \right\} \sin(x) \right\} \right\} t^3 \\ &+ \left\{ -\frac{A\cos^2(x)}{6720(1+\sigma_1)^4} \left(20160A\beta_2\sigma_1 + 20160A\beta_2\sigma_1^2 + 6720A\beta_2\sigma_1^3 + 6720\beta_2A \right) \right\} t^3 \\ &+ \left\{ -\frac{A\cos^2(x)}{6720(1+\sigma_1)^4} \left(20160A\beta_2\sigma_1 - 20160A\beta_2\sigma_1^2 + 6720A\beta_2\sigma_1^3 + 6720\beta_2A \right) \right\} t^3 \\ &+ \left\{ -\frac{A\cos^2(x)}{6720(1+\sigma_1)^4} \left(20160A\beta_2\sigma_1 - 20160A\beta_2\sigma_1^2 + 6720A\beta_2\sigma_1^3 + 6720\beta_2A \right) \right\} t^3 \\ + \left\{ -\frac{A\cos^2(x)}{6720(1+\sigma_1)^4} \left(20160A\beta_2\sigma_1 + 20160A\beta_2\sigma_1^2 + 6720A\beta_2\sigma_1^3 + 6720\beta_2A \right) \right\} t^3 \\ + \left\{ -\frac{A\cos^2(x)}{6720(1+\sigma_1)^4} \left(20160A\beta_2\sigma_1 + 20160A\beta_2\sigma_1^2 + 6720A\beta_2\sigma_1^3 + 6720\beta_2A \right) \right\} t^3 \\ t^3 + t^3 +$$

$$-\frac{A/\cos(x)}{6720(1+\sigma_{1})^{4}} \Big(-30240A^{2}\delta\sigma_{1}^{2} - 10080A\beta_{2}\sigma_{1} + 3360\sigma_{2}\sigma_{1}^{3} - 10080A^{2}\delta \\ + 10080\sigma_{2}\sigma_{1} - 3360A\beta_{2}\sigma_{1}^{3} - 3360 - 30240A^{2}\delta\sigma_{1} - 3360\beta_{2}A \\ + 10080\sigma_{2}\sigma_{1}^{2} + 3360\sigma_{2} - 10080\sigma_{1}^{2} - 10080A^{2}\delta\sigma_{1}^{3} \\ - 10080A\beta_{2}\sigma_{1}^{2} - 3360\sigma_{1}^{3} - 10080\sigma_{1}\Big) \Big\}t^{2} \\ - \frac{A\cos(x)}{6720(1+\sigma_{1})^{4}} \Big(6720 + 6720\sigma_{1}^{4} + 40320\sigma_{1}^{2} + 26880\sigma_{1}^{3} + 26880\sigma_{1}\Big) \\ - \frac{A}{6720(1+\sigma_{1})^{4}} \Big(-6720 - 26880\sigma_{1} - 26880\sigma_{1}^{3} - 6720\sigma_{1}^{4} - 40320\sigma_{1}^{2}\Big),$$
(5.4)

$$\begin{split} \theta_{2} &= \left\{ -\frac{A\cos\left(x\right)^{5}}{24(1+\sigma_{1})^{2}} \left(-540A^{4}\delta\alpha^{2} - 540A^{4}\delta\alpha^{2}\sigma_{1} + 216A^{4}b\delta\beta_{2} \right) \\ &- \frac{A\cos\left(x\right)^{3}}{24(1+\sigma_{1})^{2}} \left(-24A^{3}b\beta_{2}^{2} + 36A^{3}\beta_{2}\alpha^{2} + 24A^{2}b\sigma_{2}\beta_{2} + 36A^{3}\beta_{2}\alpha^{2}\sigma_{1} + 36A^{2}\alpha^{2}\sigma_{1} \\ &+ 36A^{2}\alpha\gamma\sigma_{1} + 702A^{4}\delta\alpha^{2}\sigma_{1} - 36A^{2}\sigma_{2}\alpha^{2}\sigma_{1} - 36A^{2}\sigma_{2}\alpha^{2} + 36A^{2}\alpha^{2} \\ &- 24A^{2}b\beta_{1}\gamma - 288A^{4}b\delta\beta_{2} - 24A^{2}b\beta_{2} + 36A^{2}\alpha\gamma + 702A^{4}\delta\alpha^{2} \right) \\ &+ \left[-\frac{A\cos^{2}(x)}{24(1+\sigma_{1})^{2}} \left(-24A^{2}b\gamma\sigma_{2} + 36A^{2}\beta_{1}\alpha^{2}\sigma_{1} - 24A^{2}b\beta_{1}\beta_{2} + 36A^{2}\beta_{1}\alpha^{2} + 36A^{2}\beta_{2}\alpha \\ &+ 24A^{3}b\gamma\beta_{2} + 24A^{2}b\gamma + 180A^{4}b\gamma\delta + 36A^{2}\beta_{2}\sigma_{1}\alpha \right) + \frac{9A^{5}b\gamma\delta\cos^{4}(x)}{(1+\sigma_{1})^{2}} \\ &- \frac{A\cos(x)}{24(1+\sigma_{1})^{2}} \left(-6\alpha A\sigma_{1} + 6bA\beta_{1} + 18A^{3}b\beta_{1}\delta - 72\alpha A^{3}\delta - 96A^{3}\gamma\alpha^{2}\sigma_{1} + 48A^{3}b\gamma\beta_{2} \\ &- 6A^{2}\beta_{2}\sigma_{1}\alpha - 6A^{2}\beta_{2}\alpha - 96A^{3}\gamma\alpha^{2} + 6\alpha A\sigma_{2} - 6\alpha A - 6Ab\beta_{1}\sigma_{2} \\ &- 72\alpha A^{3}\delta\sigma_{1} + 6\alpha A\sigma_{2}\sigma_{1} + 6A^{2}b\beta_{1}\beta_{2} \right) \\ &- \frac{A\cos^{3}(x)}{24(1+\sigma_{1})^{2}} \left(192A^{3}\gamma\alpha^{2}\sigma_{1} + 192A^{3}\gamma\alpha^{2} + 108\alpha A^{3}\delta - 96A^{3}b\gamma\beta_{2} \\ &+ 108\alpha A^{3}\delta\sigma_{1} - 54A^{3}b\beta_{1}\delta \right) \\ &- \frac{A}{24(1+\sigma_{1})^{2}} \left(-12A^{2}b\gamma - 12A^{2}\beta_{2}\alpha - 12A^{3}b\gamma\beta_{2} - 36A^{4}b\gamma\delta - 6A^{2}\beta_{1}\alpha^{2} - 12A^{2}\beta_{2}\sigma_{1}\alpha \\ &+ 12A^{2}b\gamma\sigma_{2} - 6A^{2}\beta_{1}\alpha^{2}\sigma_{1} \right) \right] \sin(x) \end{split}$$

$$\begin{split} &-\frac{A\cos^2(x)}{24(1+\sigma_1)^2} \Big(3Ab + 72A^4b\delta\beta_2 + 48A^3b\gamma^2 + 3Ab\sigma_2^2 + 189A^5b\delta^2 + 6A\beta_1a\sigma_1 \\ &- 6A^2b\sigma_2\beta_2 + 6A\beta_1a - 3Ab\beta_1^2 + 192A^3\beta_2a^2 + 192A^3\beta_2a^2\sigma_1 \\ &+ 72A^3b\delta + 6A^2b\beta_2 - 45A^3b\beta_2^2 - 72A^3b\delta\sigma_2 - 6Ab\sigma_2 \Big) \\ &-\frac{A\cos(x)}{24(1+\sigma_1)^2} \Big(72A^4b\delta\beta_2 - 198A^4\delta a^2\sigma_1 - 30A^3\beta_2a^2 - 30A^2a^2\sigma_1 + 24A^3b\beta_2^2 \\ &- 30A^3\beta_2a^2\sigma_1 + 30A^2\sigma_2a^2 - 24A^2b\sigma_2\beta_2 + 12A^2b\beta_1\gamma + 30A^2\sigma_2a^2\sigma_1 \\ &- 198A^4\delta a^2 - 30A^2a\gamma\sigma_1 + 24A^2b\beta_2 - 30A^2a\gamma - 30A^2a^2 \Big) \\ &-\frac{A\cos^4(x)}{24(1+\sigma_1)^2} \Big(-405A^5b\delta^2 - 192A^3\beta_2a^2 + 54A^3b\delta\sigma_2 - 54A^4b\delta\beta_2 - 48A^3b\gamma^2 \\ &+ 48A^3b\beta_2^2 - 54A^3b\delta - 192A^3\beta_2a^2\sigma_1 \Big) - \frac{81A^6b\delta^2\cos^6(x)}{8(1+\sigma_1)^2} \\ &-\frac{A}{24(1+\sigma_1)^2} \Big(-12A^3b\gamma^2 - 18A^4b\delta\beta_2 - 3Ab - 3Ab\sigma_2^2 - 3A^3b\beta_2^2 \\ &+ 6A^2b\sigma_2\beta_2 - 3A\beta_1a - 27A^5b\delta^2 - 3A\beta_1a\sigma_1 - 24A^3\beta_2a^2\sigma_1 - 18A^3b\delta \\ &- 6A^2b\beta_2 + 18A^3b\delta\sigma_2 - 24A^3\beta_2a^2 + 6Ab\sigma_2 \Big) \Big\} t^4 \\ &+ \Big\{ -\frac{A(-48A^2a\gamma - 48A^2a\gamma\sigma_1)\cos^3(x)}{24(1+\sigma_1)^2} - \frac{A(-6A\beta_1a - 8A\beta_1a\sigma_1)\cos^2(x)}{24(1+\sigma_1)^2} \\ &- \frac{A(-48A^2\beta_2\sigma_1a - 48A^2\beta_2\alpha)}{24(1+\sigma_1)^2} - \frac{A(-6A\beta_1a - 8A\beta_1a\sigma_1)\cos^2(x)}{24(1+\sigma_1)^2} \\ &+ \Big\{ -\frac{A\cos(x)}{24(1+\sigma_1)^2} \Big(8aA + 96aA^3\delta\sigma_1 - 8aA\sigma_2\sigma_1 - 8aA\sigma_2 + 8A^2\beta_2a + 96aA^3\delta \\ &+ 8A^2\beta_2\sigma_1a + 8aA\sigma_1 \Big) \\ &- \frac{A\cos(x)}{24(1+\sigma_1)^2} \Big(12-48A^2b\beta_2 + 120A^2a^2\sigma_1^2 - 48A^2\sigma_1b\beta_2 + 24\sigma_1 + 12\sigma_1^2 \\ &- 12\sigma_1a\beta_1 + 240A^2a^2\sigma_1 - 12a\beta_1 + 120A^2a^2 \Big) \\ &- \frac{A\cos^3(x)}{24(1+\sigma_1)^2} \Big(48A^2b\beta_2 - 144A^2a^2\sigma_1^2 - 288A^2a^2\sigma_1 - 144A^2a^2 + 48A^2\sigma_1b\beta_2 \Big) \Big\}$$

$$-\frac{A}{24(1+\sigma_{1})^{2}} \Big(24Aa\gamma + 36A^{3}b\delta + 12A^{2}b\beta_{2} - 12Ab\sigma_{2} + 12Ab + 36A^{3}\sigma_{1}b\delta + 24A\sigma_{1}a\gamma + 12A\sigma_{1}b - 12A\sigma_{1}b\sigma_{2} + 12A^{2}\sigma_{1}b\beta_{2} \Big) + \Big[-\frac{A\cos^{2}(x)}{24(1+\sigma_{1})^{2}} \Big(-48A^{2}b\gamma - 108a\delta A^{2} - 108\sigma_{1}a\delta A^{2} - 48A^{2}\sigma_{1}b\gamma \Big) - \frac{A}{24(1+\sigma_{1})^{2}} \Big(12\sigma_{1}a\beta_{2}A + 36\sigma_{1}a\delta A^{2} - 12a\sigma_{2} + 12a\beta_{2}A - 12\sigma_{1}a\sigma_{2} + 12\sigma_{1}a + 24A^{2}b\gamma + 36a\delta A^{2} + 12a + 24A^{2}\sigma_{1}b\gamma \Big) - \frac{A\cos(x)}{24(1+\sigma_{1})^{2}} \Big(240aA\sigma_{1} + 120aA - 48a\beta_{2}A - 12bA\beta_{1} - 48\sigma_{1}a\beta_{2}A + 120aA\sigma_{1}^{2} - 12\sigma_{1}bA\beta_{1} \Big) \Big] \sin(x) - \frac{A\cos^{2}(x)}{24(1+\sigma_{1})^{2}} \Big(-12Ab + 12Ab\sigma_{2} - 144A^{3}\sigma_{1}b\delta - 144A^{3}b\delta - 12A^{2}b\beta_{2} - 48Aa\gamma - 12A^{2}\sigma_{1}b\beta_{2} - 48A\sigma_{1}a\gamma - 12A\sigma_{1}b + 12A\sigma_{1}b\sigma_{2} \Big) - \frac{A(108A^{3}b\delta + 108A^{3}\sigma_{1}b\delta)\cos^{4}(x)}{24(1+\sigma_{1})^{2}} \Big\} t^{2} \times \Big\{ \frac{A(-48aA - 96aA\sigma_{1} - 48aA\sigma_{1}^{2})\cos(x)\sin(x)}{24(1+\sigma_{1})^{2}} + \frac{A(-24 - 48\sigma_{1} - 24\sigma_{1}^{2})\cos(x)}{24(1+\sigma_{1})^{2}} \Big\} t - \frac{A(24 + 48\sigma_{1} + 24\sigma_{1}^{2})\cos(x)}{24(1+\sigma_{1})^{2}} - \frac{A(-24 - 48\sigma_{1} - 24\sigma_{1}^{2})}{24(1+\sigma_{1})^{2}}.$$
(5.5)

Now we make calculations for the results obtained by the variational iteration method using the Maple software package with the following arbitrary constants:

 $a = 0.5, A = 0.001, b = 0.5, \alpha = 1, \beta_1 = \beta_2 = 0.05, \gamma = 1, \delta = 0.8.$ (5.6)

5.1. Special Case

- (1) If we take into our consideration the first iteration (i.e., $u = u_1$ and $\theta = \theta_1$). See Figures 15, 16, 17, and 18.
- (2) If the magnetic field and rotation are neglected, the components of the displacement u_2 and temperature θ_2 take the following forms. See Figures 19 and 20.



Figure 11: Variations of the displacement u_2 and temperature θ_2 for various values of the axis x and time t when $\Omega = 0.1$, $\sigma_1 = 0.2$, $\sigma_2 = 0.1$.



Figure 12: Variations of the displacement u_2 and temperature θ_2 for various values of the axis x and rotation Ω when t = 0.1, $\sigma_1 = 0.2$, $\sigma_2 = 0.1$.

6. Discussion

With the view of illustrating the theoretical results obtained in the preceding sections, a numerical result is calculated for the homotopy perturbation method and variational iteration method.

Figures (1–10) illustrate the influences of time *t*, rotation Ω , and sensitive pats of the magnetic field σ_1 and σ_2 for the iterations ($u = V_0 + V_1 + V_2$ and $\theta = \Theta_0 + \Theta_1 + \Theta_2$) and ($u = V_0 + V_1$ and $\theta = \Theta_0 + \Theta_1$), and if the rotation and magnetic field neglected, respectively, respect to the coordinate *x* for the homotopy perturbation method. Figures (11–20) illustrate the influences of time *t*, rotation Ω , and sensitive pats of the magnetic field σ_1 and σ_2 for the iterations ($u = V_2$, $\theta = \theta_2$ and $u = V_1$ and $\theta = \theta_1$), and if the rotation and magnetic field have been neglected, respectively, respect to the coordinate *x* for the variational iteration method.

From Figures 1 and 11, it is concluded that the displacement u and temperature θ start from their maximum values, decrease and increase periodically with an increasing of



Figure 13: Variations of the displacement u_2 and temperature θ_2 for various values of the axis x and magnetic field σ_1 when t = 0.1, $\Omega = 0.1$, $\sigma_2 = 0.1$.



Figure 14: Variations of the displacement u_2 and temperature θ_2 for various values of the axis x and magnetic field σ_2 when t = 0.1, $\Omega = 0.1$, $\sigma_1 = 0.1$.



Figure 15: Variations of the displacement u_1 and temperature θ_1 for various values of the axis x and time t when $\Omega = 0.1$, $\sigma_1 = 0.2$, $\sigma_2 = 0.1$.



Figure 16: Variations of the displacement u_1 and temperature θ_1 for various values of the axis x and rotation Ω when t = 0.1, $\sigma_1 = 0.2$, $\sigma_2 = 0.1$.



Figure 17: Variations of the displacement u_1 and temperature θ_1 for various values of the axis x and magnetic field σ_1 when t = 0.1, $\Omega = 0.1$, $\sigma_2 = 0.1$.

the coordinate x, also, it is obvious that their values take the minimum values and increases with the increasing values of the time t. From Figures 2, 3, 4, 12, 13, 14, it is seen that the components of the displacement u and temperature θ begin from the minimum values near zero increase and then decrease periodically with the coordinate x, it is clear also that there are a sligh increasing with an increasing of the sensitive parts of the magnetic field, also, one can see that u and θ decrease with an increasing of the rotation Ω .

Figures 5–8 and 15–18 display the first iteration with respect to the homotopy perturbation method and variational iteration method on the influences of the parameters time *t*, rotation Ω , and sensitive pats of the magnetic field σ_1 and σ_2 to obtain the displacement and the temperature components on the medium due to the harmonic wave propagation. It is shown that the increasing of the coordinate *x* sensitive an increasing and dereasing on them periodically due to appearance of the pairs (cos, sin) in the initial condition and the approximate solutions; it is also clear that the components begin from their minimum values and increase absolutely with the variation of the time *t*. With the variations of the rotation and magnetic field tends to slightly affect on the displacement and the temperature.



Figure 18: Variations of the displacement u_1 and temperature θ_1 for various values of the axis x and magnetic field σ_2 when t = 0.1, $\Omega = 0.1$, $\sigma_1 = 0.1$.



Figure 19: Variations of the displacement u_2 and temperature θ_2 for various values of the axis x and time t ($u = u_2$ and $\theta = \theta_2$) when $\Omega = \sigma_1 = \sigma_2 = 0$.



Figure 20: Variations of the displacement u_1 and temperature θ_1 for various values of the axis x and time t ($u = u_1$ and $\theta = \theta_1$) when $\Omega = \sigma_1 = \sigma_2 = 0$.

It seems too that there are a clear differs between the results obtained by the HPM and the corresponding results obtained by VIM resultant to the appearance of the high order of time in VIM tends to the high values of the approximate solution comparing with the results obtained by HPM. Because of the results obtained, we concluded that the homotopy perturbation method is more effective and powerful than the variational iteration method.

On the other hand, from Figures 9, 10, 19, and 20, it is obvious that if the rotaion and magnetic field are neglected, the approximate solutions by HPM and VIM in first iteration are the same in both methods and agree with the results obtained by Sweilam and Khader [1].

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