THE JAMES CONSTANT OF NORMALIZED NORMS ON \mathbb{R}^2

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We introduce a new class of normalized norms on \mathbb{R}^2 which properly contains all absolute normalized norms. We also give a criterion for deciding whether a given norm in this class is uniformly nonsquare. Moreover, an estimate for the James constant is presented and the exact value of some certain norms is computed. This gives a partial answer to the question raised by Kato et al.

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1. Introduction and preliminaries

A norm $\|\cdot\|$ on \mathbb{C}^2 (resp., \mathbb{R}^2) is said to be *absolute* if $\|(z,w)\| = \|(|z|,|w|)\|$ for all $z, w \in \mathbb{C}$ (resp., \mathbb{R}), and *normalized* if $\|(1,0)\| = \|(0,1)\| = 1$. The ℓ_p -norms $\|\cdot\|_p$ are such examples:

$$\left| \left| (z,w) \right| \right|_{p} = \begin{cases} \left(|z|^{p} + |w|^{p} \right)^{1/p} & \text{if } 1 \le p < \infty, \\ \max\left\{ |z|, |w| \right\} & \text{if } p = \infty. \end{cases}$$
(1.1)

Let AN_2 be the family of all absolute normalized norms on \mathbb{C}^2 (resp., \mathbb{R}^2), and Ψ_2 the family of all continuous convex functions ψ on [0,1] such that $\psi(0) = \psi(1) = 1$ and $\max\{1-t, t\} \le \psi(t) \le 1$ ($0 \le t \le 1$). According to Bonsall and Duncan [1], AN_2 and Ψ_2 are in a one-to-one correspondence under the equation

$$\psi(t) = ||(1-t,t)|| \quad (0 \le t \le 1).$$
(1.2)

Indeed, for all $\psi \in \Psi_2$, let

$$\left| \left| (z,w) \right| \right|_{\psi} = \begin{cases} \left(|z| + |w| \right) \psi \left(\frac{|w|}{|z| + |w|} \right) & \text{if } (z,w) \neq (0,0), \\ 0 & \text{if } (z,w) = (0,0). \end{cases}$$
(1.3)

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Then $\|\cdot\|_{\psi} \in AN_2$, and $\|\cdot\|_{\psi}$ satisfies (1.2). From this result, we can consider many non- ℓ_{p} -type norms easily. Now let

$$\psi_p(t) = \begin{cases} \left((1-t)^p + t^p \right)^{1/p} & \text{if } 1 \le p < \infty, \\ \max\{1-t, t\} & \text{if } p = \infty. \end{cases}$$
(1.4)

Then $\psi_p(t) \in \Psi_2$ and, as is easily seen, the ℓ_p -norm $\|\cdot\|_p$ is associated with ψ_p .

If *X* is a Banach space, then *X* is *uniformly nonsquare* if there exists $\delta \in (0, 1)$ such that for any $x, y \in S_X$,

either
$$||x + y|| \le 2(1 - \delta)$$
 or $||x - y|| \le 2(1 - \delta)$, (1.5)

where $S_X = \{x \in X : ||x|| = 1\}$. The *James constant* J(X) is defined by

$$J(X) = \sup \{ \min \{ \|x + y\|, \|x - y\| \} : x, y \in S_X \}.$$
(1.6)

The modulus of convexity of *X*, $\delta_X : [0,2] \rightarrow [0,1]$ is defined by

$$\delta_X(\varepsilon) = \inf \left\{ 1 - \frac{1}{2} \|x + y\| : x, \ y \in S_X, \ \|x - y\| \ge \varepsilon \right\}.$$
(1.7)

The preceding parameters have been recently studied by several authors (cf. [4–6, 8, 9]). We collect together some known results.

PROPOSITION 1.1. Let X be a nontrivial Banach space, then

- (i) $\sqrt{2} \leq J(X) \leq 2$ (Gao and Lau [5]),
- (ii) if X is a Hilbert space, then $J(X) = \sqrt{2}$; the converse is not true (Gao and Lau [5]),
- (iii) X is uniformly nonsquare if and only if J(X) < 2 (Gao and Lau [5]),
- (iv) $2J(X) 2 \le J(X^*) \le J(X)/2 + 1$, $J(X^{**}) = J(X)$, and there exists a Banach space X such that $J(X^*) \ne J(X)$ (Kato et al. [8]),
- (v) if $2 \le p \le \infty$, then $\delta_{\ell_p}(\varepsilon) = 1 (1 (\varepsilon/2)^p)^{1/p}$ (Hanner [6]),
- (vi) $J(X) = \sup \{ \varepsilon \in (0,2) : \delta_X(\varepsilon) \le 1 \varepsilon/2 \}$ (Gao and Lau [5]).

The paper is organized as follows. In Section 2 we introduce a new class of normalized norms on \mathbb{R}^2 . This class properly contains all absolute normalized norms of Bonsall and Duncan [1]. The so-called generalized Day-James space, $\ell_{\psi}-\ell_{\varphi}$, where $\psi, \varphi \in \Psi_2$, is introduced and studied. More precisely, we prove that $(\ell_{\psi}-\ell_{\varphi})^* = \ell_{\psi^*}-\ell_{\varphi^*}$ where ψ^* and φ^* are the dual functions of ψ and φ , respectively. In Section 3, the upper bound of the James constant of the generalized Day-James space is given. Furthermore, we compute $J(\ell_{\psi}-\ell_{\infty})$ and deduce that every generalized Day-James space except $\ell_1-\ell_1$ and $\ell_{\infty}-\ell_{\infty}$ is uniformly nonsquare. This result strengthens Corollary 3 of Saito et al. [10].

2. Generalized Day-James spaces

In this section, we introduce a new class of normalized norms on \mathbb{R}^2 which properly contains all absolute normalized norms of Bonsall and Duncan [1]. Moreover, we introduce a two-dimensional normed space which is a generalization of Day-James ℓ_p - ℓ_q spaces. LEMMA 2.1. Let $\psi \in \Psi_2$ and let $\|\cdot\|_{\psi,\psi_{\infty}}$ be a function on \mathbb{R}^2 defined by, for all $(z,w) \in \mathbb{R}^2$,

$$\begin{aligned} ||(z,w)||_{\psi,\psi_{\infty}} &:= \max \{ ||(z^{+},w^{+})||_{\psi}, ||(z^{-},w^{-})||_{\psi} \}, \\ &= \begin{cases} ||(z,w)||_{\psi} & \text{if } zw \ge 0, \\ ||(z,w)||_{\infty} & \text{if } zw \le 0, \end{cases} \end{aligned}$$
(2.1)

where x^+ and x^- are positive and negative parts of $x \in \mathbb{R}$, that is, $x^+ = \max\{x, 0\}$ and $x^- = \max\{-x, 0\}$. Then $\|\cdot\|_{\psi, \psi_{\infty}}$ is a norm on \mathbb{R}^2 .

For convenience, we put $\mathscr{B}_{\psi_1,\psi_2} := \{(z,w) \in \mathbb{R}^2 : ||(z,w)||_{\psi_1,\psi_2} \le 1\}.$

THEOREM 2.2. Let $\psi, \varphi \in \Psi_2$ and

$$||(z,w)||_{\psi,\varphi} := \begin{cases} ||(z,w)||_{\psi} & \text{if } zw \ge 0, \\ ||(z,w)||_{\varphi} & \text{if } zw \le 0 \end{cases}$$
(2.2)

for all $(z, w) \in \mathbb{R}^2$. Then $\|\cdot\|_{\psi, \varphi}$ is a norm on \mathbb{R}^2 . Denote by N_2 the family of all such preceding norms.

Proof. Let $\psi, \varphi \in \Psi_2$, we only show $\|\cdot\|_{\psi,\varphi}$ satisfies the triangle inequality. To this end, it suffices to prove that $\mathcal{B}_{\psi,\varphi}$ is convex. By Lemma 2.1, we have that $\mathcal{B}_{\psi,\psi_{\infty}}$ and $\mathcal{B}_{\varphi,\psi_{\infty}}$ are closed unit balls of $\|\cdot\|_{\psi,\psi_{\infty}}$ and $\|\cdot\|_{\varphi,\psi_{\infty}}$, respectively, and so $\mathcal{B}_{\psi,\psi_{\infty}}$ and $\mathcal{B}_{\varphi,\psi_{\infty}}$ are convex sets. We define $T : \mathbb{R}^2 \to \mathbb{R}^2$ by

$$T((z,w)) = (-z,w) \quad \forall (z,w) \in \mathbb{R}^2.$$
(2.3)

Then *T* is a linear operator and $T(\mathcal{B}_{\varphi,\psi_{\infty}}) = \mathcal{B}_{\psi_{\infty},\varphi}$, which implies that $\mathcal{B}_{\psi_{\infty},\varphi}$ is convex and so $\mathcal{B}_{\psi,\varphi} = \mathcal{B}_{\psi_{\infty},\varphi} \cap \mathcal{B}_{\psi,\psi_{\infty}}$ is convex.

Taking $\psi = \psi_p$ and $\varphi = \psi_q$ $(1 \le p, q \le \infty)$ in Theorem 2.2, we obtain the following.

COROLLARY 2.3 (Day-James $\ell_p - \ell_q$ spaces). For $1 \le p$, $q \le \infty$, denote by $\ell_p - \ell_q$ the Day-James space, that is, \mathbb{R}^2 with the norm defined by, for all $(z, w) \in \mathbb{R}^2$,

$$||(z,w)||_{p,q} = \begin{cases} ||(z,w)||_p & \text{if } zw \ge 0, \\ ||(z,w)||_q & \text{if } zw \le 0. \end{cases}$$
(2.4)

James [7] considered the $\ell_p - \ell_{p'}$ space as an example of a Banach space which is isometric to its dual but which is not given by a Hilbert norm when $p \neq 2$. Day [2] considered even more general spaces, namely, if $(X, \|\cdot\|)$ is a two-dimensional Banach space and $(X^*, \|\cdot\|^*)$ its dual, then the *X*-*X*^{*} space is the space *X* with the norm defined by, for all $(z, w) \in \mathbb{R}^2$,

$$||(z,w)||_{X,X^*} = \begin{cases} ||(z,w)|| & \text{if } zw \ge 0, \\ ||(z,w)||^* & \text{if } zw \le 0. \end{cases}$$
(2.5)

For $\psi, \varphi \in \Psi_2$, denote by $\ell_{\psi} - \ell_{\varphi}$ the generalized Day-James space, that is, \mathbb{R}^2 with the norm $\|\cdot\|_{\psi,\varphi}$ defined by (2.2). For ψ_p defined by (1.4), we write $\ell_{\psi} - \ell_p$ for $\ell_{\psi} - \ell_{\psi_p}$. For example, if $1 \le p, q \le \infty$, $\ell_p - \ell_q$ means $\ell_{\psi_p} - \ell_{\psi_q}$.

It is worthwhile to mention that there is a normalized norm which is not absolute.

PROPOSITION 2.4. There is $\psi \in \Psi_2$ such that $\ell_{\psi} - \ell_{\infty}$ is not isometrically isomorphic to $\ell_{\varphi} - \ell_{\varphi}$ for all $\varphi \in \Psi_2$.

Proof. Let

$$\psi(t) := \begin{cases} 1-t & \text{if } 0 \le t \le \frac{1}{8}, \\ \frac{11-4t}{12} & \text{if } \frac{1}{8} \le t \le \frac{1}{2}, \\ \frac{1+t}{2} & \text{if } \frac{1}{2} \le t \le 1. \end{cases}$$
(2.6)

We observe that the sphere of $\ell_{\psi} - \ell_{\infty}$ is the octagon whose right half consists of 4 segments of different lengths. Suppose that there are $\varphi \in \Psi_2$ and an isometric isomorphism from $\ell_{\psi} - \ell_{\infty}$ onto $\ell_{\varphi} - \ell_{\varphi}$. Since the image of each segment in $\ell_{\psi} - \ell_{\infty}$ is again a segment of the same length in $\ell_{\varphi} - \ell_{\varphi}$, the sphere of $\ell_{\varphi} - \ell_{\varphi}$ must be the octagon whose each corresponding side has the same length (measured by $\|\cdot\|_{\varphi}$). We show that this cannot happen. Consider $(1,0) \in S_{\ell_{\varphi} - \ell_{\varphi}}$. If (1,0) is an extreme point of $B_{\ell_{\varphi} - \ell_{\varphi}}$, then $S_{\ell_{\varphi} - \ell_{\varphi}}$ contains 4 segments of same lengths since $\|\cdot\|_{\varphi}$ is absolute. On the other hand, if (1,0) is an not extreme point of $B_{\ell_{\varphi} - \ell_{\varphi}}$, again $S_{\ell_{\varphi} - \ell_{\varphi}}$ contains 4 segments of same lengths.

Next, we prove that the dual of a generalized Day-James space is again a generalized Day-James space. Recall that, for $\psi \in \Psi_2$, the *dual function* ψ^* of ψ is defined by

$$\psi^*(s) = \max_{0 \le t \le 1} \frac{(1-s)(1-t) + st}{\psi(t)}$$
(2.7)

for all $s \in [0,1]$. It was proved that $\psi^* \in \Psi_2$ and $(\ell_{\psi} - \ell_{\psi})^* = \ell_{\psi^*} - \ell_{\psi^*}$ (see [3, Proposition 1 and Theorem 2]). We generalize this result to our spaces as follows.

THEOREM 2.5. For $\psi, \varphi \in \Psi_2$, there is an isometric isomorphism that identifies $(\ell_{\psi}-\ell_{\varphi})^*$ with $\ell_{\psi^*}-\ell_{\varphi^*}$ such that if $f \in (\ell_{\psi}-\ell_{\varphi})^*$ is identified with the element $(z,w) \in \ell_{\psi^*}-\ell_{\varphi^*}$, then

$$f(u,v) = zu + wv \tag{2.8}$$

for all $(u, v) \in \mathbb{R}^2$.

Proof. We can prove analogous to [3, Theorem 2].

3. The James constant and uniform nonsquareness

The next lemmas are crucial for proving the main theorems.

LEMMA 3.1. Let $\psi, \varphi \in \Psi_2$. Then (i) $\|\cdot\|_{\infty} \leq \|\cdot\|_{\psi,\varphi} \leq \|\cdot\|_1$,

- (ii) $(1/M_{\psi,\varphi}) \| \cdot \|_{\psi} \le \| \cdot \|_{\psi,\varphi} \le M_{\varphi,\psi} \| \cdot \|_{\psi},$
- (iii) $(1/M_{\varphi,\psi}) \| \cdot \|_{\varphi} \le \| \cdot \|_{\psi,\varphi} \le M_{\psi,\varphi} \| \cdot \|_{\varphi},$
- where $M_{\varphi,\psi} = \max_{0 \le t \le 1} \varphi(t)/\psi(t)$ and $M_{\psi,\varphi} = \max_{0 \le t \le 1} \psi(t)/\varphi(t)$.

LEMMA 3.2. Let $\psi, \varphi \in \Psi_2$ and let Q_i (i = 1, ..., 4) denote the *i*th quadrant in \mathbb{R}^2 . Suppose that $x, y \in S_{\ell_w - \ell_w}$, then the following statements are true.

- (i) If $x, y \in Q_1$, then $x + y \in Q_1$ and $x y \in Q_2 \cup Q_4$.
- (ii) If $x, y \in Q_2$, then $x + y \in Q_2$ and $x y \in Q_1 \cup Q_3$.
- (iii) If $\psi(t) \le \varphi(t)$ for all $t \in [0,1]$ and $x y \in Q_2^\circ \cup Q_4^\circ$, where Q_2° and Q_4° are the interiors of Q_2 and Q_4 , respectively, then $x + y \in Q_1 \cup Q_3$.

We will estimate the James constant of ℓ_{ψ} - ℓ_{φ} .

THEOREM 3.3. Let $\psi, \varphi \in \Psi_2$ with $\psi(t) \le \varphi(t)$ for all $t \in [0, 1]$, let $M_{\varphi, \psi} = \max_{0 \le t \le 1} \varphi(t)/\psi(t)$, and let $\delta_{\psi}(\cdot)$ be the modulus of convexity of ℓ_{ψ} - ℓ_{ψ} . Then for $\varepsilon \in [0, 2]$,

$$\delta_{\psi,\varphi}(\varepsilon) \ge \min\left\{1 - M_{\varphi,\psi}(1 - \delta_{\psi}(\varepsilon)), \ \delta_{\psi}\left(\frac{\varepsilon}{M_{\varphi,\psi}}\right)\right\},\tag{3.1}$$

where $\delta_{\psi,\varphi}(\cdot)$ is the modulus of convexity of ℓ_{ψ} - ℓ_{φ} . Consequently,

$$J(\ell_{\psi}-\ell_{\varphi}) \leq \sup\left\{\varepsilon \in (0,2) : \varepsilon \leq 2M_{\varphi,\psi}(1-\delta_{\psi}(\varepsilon)) \text{ or } \varepsilon \leq 2\left(1-\delta_{\psi}\left(\frac{\varepsilon}{M_{\varphi,\psi}}\right)\right)\right\}.$$
 (3.2)

Proof. By Lemma 3.1(ii), we have

$$\|\cdot\|_{\psi} \le \|\cdot\|_{\psi,\varphi} \le M_{\varphi,\psi}\|\cdot\|_{\psi}.$$
(3.3)

We now evaluate the modulus of convexity $\delta_{\psi,\varphi}$ for ℓ_{ψ} - ℓ_{φ} . We consider two cases.

Case 1. Take $||x||_{\psi,\varphi} = ||y||_{\psi,\varphi} = 1$ with $||x - y||_{\psi,\varphi} \ge \varepsilon$, where $x - y \in Q_1 \cup Q_3$. Thus $||x||_{\psi} \le 1$, $||y||_{\psi} \le 1$, and $||x - y||_{\psi} \ge \varepsilon$, which implies that

$$\frac{1}{2}\|x+y\|_{\psi} \le 1 - \delta_{\psi}(\varepsilon). \tag{3.4}$$

This in turn implies

$$\frac{1}{2} \|x+y\|_{\psi,\varphi} \le \frac{1}{2} M_{\varphi,\psi} \|x+y\|_{\psi} \le M_{\varphi,\psi} (1-\delta_{\psi}(\varepsilon)),$$
(3.5)

thus

$$1 - \frac{1}{2} \| x + y \|_{\psi, \varphi} \ge 1 - M_{\varphi, \psi} (1 - \delta_{\psi}(\varepsilon)).$$
(3.6)

Case 2. Now take x, y as above, but with $x - y \in Q_2^\circ \cup Q_4^\circ$. By Lemma 3.2(iii), $x + y \in Q_1 \cup Q_3$. Since $||x - y||_{\psi, \varphi} \ge \varepsilon$,

$$\|x - y\|_{\psi} \ge \frac{\|x - y\|_{\psi,\varphi}}{M_{\varphi,\psi}} \ge \frac{\varepsilon}{M_{\varphi,\psi}}.$$
(3.7)

Then

$$\frac{1}{2} \|x + y\|_{\psi,\varphi} = \frac{1}{2} \|x + y\|_{\psi} \le 1 - \delta_{\psi} \left(\frac{\varepsilon}{M_{\varphi,\psi}}\right), \tag{3.8}$$

and so

$$1 - \frac{1}{2} \|x + y\|_{\psi, \varphi} \ge \delta_{\psi} \left(\frac{\varepsilon}{M_{\varphi, \psi}}\right).$$
(3.9)

 \Box

Hence we obtain (3.1). By Proposition 1.1(vi), (3.2) follows.

The following corollary shows that we can have equality in (3.2).

COROLLARY 3.4 [4, 8]. If $1 \le q \le p < \infty$ and $p \ge 2$, then

$$J(\ell_p - \ell_q) \le 2 \left(\frac{2^{p/q}}{2^{p/q} + 2}\right)^{1/p}.$$
(3.10)

In particular, if p = 2 and q = 1, then $J(\ell_2 - \ell_1) = \sqrt{8/3}$. *Proof.* It follows that since

$$M_{\psi_q,\psi_p} = 2^{1/q-1/p}, \qquad \delta_{\ell_p-\ell_p}(\varepsilon) = 1 - \left(1 - \left(\frac{\varepsilon}{2}\right)^p\right)^{1/p}.$$
(3.11)

Moreover, if p = 2 and q = 1, then $J(\ell_2 - \ell_1) \le \sqrt{8/3}$. Now we put

$$x_0 = \left(\frac{2+\sqrt{2}}{2\sqrt{3}}, \frac{2-\sqrt{2}}{2\sqrt{3}}\right), \qquad y_0 = \left(\frac{2-\sqrt{2}}{2\sqrt{3}}, \frac{2+\sqrt{2}}{2\sqrt{3}}\right). \tag{3.12}$$

Then

$$||x_0||_{2,1} = ||y_0||_{2,1} = 1, \qquad ||x_0 \pm y_0||_{2,1} = \sqrt{\frac{8}{3}}.$$
 (3.13)

THEOREM 3.5. Let $\psi, \varphi \in \Psi_2$ with $\psi(t) \leq \varphi(t)$ for all $t \in [0,1]$, let $M_{\varphi,\psi} = \max_{0 \leq t \leq 1} \varphi(t)/\psi(t)$, and let $\delta_{\varphi}(\cdot)$ be the modulus of convexity of ℓ_{φ} - ℓ_{φ} . Then for $\varepsilon \in [0,2]$,

$$\delta_{\psi,\varphi}(\varepsilon) \ge 1 - M_{\varphi,\psi} \left(1 - \delta_{\varphi} \left(\frac{\varepsilon}{M_{\varphi,\psi}} \right) \right), \tag{3.14}$$

where $\delta_{\psi,\varphi}(\cdot)$ is the modulus of convexity of ℓ_{ψ} - ℓ_{φ} . Consequently,

$$J(\ell_{\psi} - \ell_{\varphi}) \le \sup \left\{ \varepsilon \in (0, 2) : \varepsilon \le 2M_{\varphi, \psi} \left(1 - \delta_{\varphi} \left(\frac{\varepsilon}{M_{\varphi, \psi}} \right) \right) \right\}.$$
(3.15)

Proof. By Lemma 3.1(iii), we have

$$\frac{1}{M_{\varphi,\psi}} \|\cdot\|_{\varphi} \le \|\cdot\|_{\psi,\varphi} \le \|\cdot\|_{\varphi}.$$
(3.16)

We now evaluate the modulus of convexity $\delta_{\psi,\varphi}$ for $\ell_{\psi}-\ell_{\varphi}$. Let

$$\|x\|_{\psi,\varphi} = \|y\|_{\psi,\varphi} = 1 \quad \text{with } \|x - y\|_{\psi,\varphi} \ge \varepsilon.$$
(3.17)

Then

$$\frac{1}{M_{\varphi,\psi}} \|x\|_{\varphi} \le 1, \qquad \frac{1}{M_{\varphi,\psi}} \|y\|_{\varphi} \le 1,$$

$$\frac{1}{M_{\varphi,\psi}} \|x-y\|_{\varphi} \ge \frac{1}{M_{\varphi,\psi}} \|x-y\|_{\psi,\varphi} \ge \frac{\varepsilon}{M_{\varphi,\psi}},$$
(3.18)

which implies that

$$\frac{1}{2M_{\varphi,\psi}} \|x+y\|_{\varphi} \le 1 - \delta_{\varphi} \left(\frac{\varepsilon}{M_{\varphi,\psi}}\right).$$
(3.19)

This in turn implies that

$$\frac{1}{2M_{\varphi,\psi}} \|x+y\|_{\psi,\varphi} \le \frac{1}{2M_{\varphi,\psi}} \|x+y\|_{\varphi} \le 1 - \delta_{\varphi} \left(\frac{\varepsilon}{M_{\varphi,\psi}}\right),\tag{3.20}$$

thus

$$1 - \frac{1}{2} \|x + y\|_{\psi,\varphi} \ge 1 - M_{\varphi,\psi} \left(1 - \delta_{\varphi} \left(\frac{\varepsilon}{M_{\varphi,\psi}} \right) \right).$$
(3.21)

Hence we obtain (3.14). By Proposition 1.1(vi), (3.15) follows. \Box

Corollary 3.6. If $2 \le q \le p < \infty$, then

$$J(\ell_p - \ell_q) \le 2^{1 - 1/p}.$$
(3.22)

It is easy to see that the estimate (3.22) is better than one obtained in [4, Example 2.4(3)].

For some generalized Day-James spaces, [8, Corollary 4] of Kato et al. gives only rough result for the estimate of the James constant, that is, for $\psi \in \Psi_2$,

$$\frac{2}{M} \le J(\ell_{\psi} - \ell_{\infty}) \le 2M, \tag{3.23}$$

where $M = \max_{0 \le t \le 1} \psi_{\infty}(t) / \psi(t)$.

However, the following theorem gives the exact value of the James constant of these spaces.

Theorem 3.7. Let $\psi \in \Psi_2$. Then

$$J(\ell_{\psi} - \ell_{\infty}) = 1 + \frac{1/2}{\psi(1/2)}.$$
(3.24)

Proof. For our convenience, we write $\|\cdot\|$ instead of $\|\cdot\|_{\psi,\psi_{\infty}}$. Let $x, y \in S_{\ell_{\psi}-\ell_{\infty}}$. We prove that

either
$$||x + y|| \le 1 + \frac{1/2}{\psi(1/2)}$$
 or $||x - y|| \le 1 + \frac{1/2}{\psi(1/2)}$. (3.25)

Let us consider the following cases.

Case 1. $x, y \in Q_1$. Let x = (a, b) and y = (c, d) where $a, b, c, d \in [0, 1]$. By Lemma 3.2(i), we have $x - y \in Q_2 \cup Q_4$. Then

$$\|x - y\| = \max\left\{|a - c|, |b - d|\right\} \le 1 \le 1 + \frac{1/2}{\psi(1/2)}.$$
(3.26)

Case 2. $x, y \in Q_2$. If x, y lies in the same segment, then $||x - y|| \le 1$. We now suppose that x = (-1, a) and y = (-c, 1) where $a, c \in [0, 1]$.

Subcase 2.1. $a \le (1/2)/\psi(1/2)$ and $c \le (1/2)/\psi(1/2)$. Then

$$\|x+y\| = \|(-1-c,1+a)\|_{\infty} = \max\{1+c,1+a\} \le 1 + \frac{1/2}{\psi(1/2)}.$$
(3.27)

Subcase 2.2. $a \ge (1/2)/\psi(1/2)$ or $c \ge (1/2)/\psi(1/2)$. Put z = (-1, 1), then

$$\|x - y\| \le \|x - z\| + \|z - y\| = 1 - a + 1 - c \le 1 + 1 - \frac{1/2}{\psi(1/2)} \le 1 + \frac{1/2}{\psi(1/2)}.$$
 (3.28)

From now on, we may assume without loss of generality that there is $\beta \in [1/2, 1]$ such that $\psi(\beta) \leq \psi(t)$ for all $t \in [0, 1]$. Indeed, $J(\ell_{\psi} - \ell_{\infty}) = J(\ell_{\tilde{\psi}} - \ell_{\infty})$ where $\tilde{\psi}(t) = \psi(1 - t)$ for all $t \in [0, 1]$.

Case 3. $x \in Q_1$ and $y \in Q_2$. Let x = (a,b), y = (-c,1) where $a,b,c \in [0,1]$. We consider three subcases.

Subcase 3.1. $a \le (1/2)/\psi(1/2)$ or $c \le (1/2)/\psi(1/2)$. Then

$$\|x - y\| = \left\| (a + c, b - 1) \right\|_{\infty} = \max\{a + c, 1 - b\} \le 1 + \frac{1/2}{\psi(1/2)}.$$
(3.29)

Subcase 3.2. $(1/2)/\psi(1/2) \le a \le c$. Then $b \le (1/2)/\psi(1/2)$ and

$$\|x+y\| = \left| \left| (a-c,b+1) \right| \right|_{\infty} = \max\{c-a, 1+b\} \le 1 + \frac{1/2}{\psi(1/2)}.$$
(3.30)

Subcase 3.3. $(1/2)/\psi(1/2) < c \le a$. We write $a = (1 - t_0)/\psi(t_0)$, $b = t_0/\psi(t_0)$ where $t_0 = b/(a+b)$ and $0 \le t_0 \le 1/2$. By the convexity of ψ and $\psi(t) \ge \psi(\beta)$ for all $0 \le t \le 1$, we

have $\psi(t_0) \ge \psi(1/2)$ and so $1/\psi(t_0) \le 1/\psi(1/2)$. By Lemma 3.1(i),

$$\begin{aligned} \|x+y\| &= \left| \left| (a,b) + (-c,1) \right| \right| \le \left| \left| (a-c,b+1) \right| \right|_1 \\ &= a-c+b+1 = \frac{1}{\psi(t_0)} + 1 - c \\ &\le \frac{1}{\psi(1/2)} + 1 - \frac{1/2}{\psi(1/2)} = 1 + \frac{1/2}{\psi(1/2)}. \end{aligned}$$
(3.31)

Case 4. $x \in Q_1$ and $y \in Q_2$. Let x = (a,b), y = (-1,c) where $a, b, c \in [0,1]$. We consider three subcases.

Subcase 4.1. $b \le (1/2)/\psi(1/2)$ or $c \le (1/2)/\psi(1/2)$. Then

$$\|x+y\| = \left\| (a-1,b+c) \right\|_{\infty} = \max\{1-a, b+c\} \le 1 + \frac{1/2}{\psi(1/2)}.$$
(3.32)

Subcase 4.2. $(1/2)/\psi(1/2) < b \le c$. Then $a \le (1/2)/\psi(1/2)$ and

$$\|x - y\| = \|(1 + a, b - c)\|_{\infty} = \max\{1 + a, c - b\} \le 1 + \frac{1/2}{\psi(1/2)}.$$
(3.33)

Subcase 4.3. $(1/2)/\psi(1/2) < c \le b$. We write $a = (1 - t_0)/\psi(t_0)$, $b = t_0/\psi(t_0)$, where $t_0 = b/(a+b)$ and $1/2 \le t_0 \le 1$. We choose $\alpha = b/(a+2b-1)$, then

$$\frac{1}{2} \le \alpha \le 1, \qquad a = \frac{1 - 2\alpha}{\alpha}b + 1.$$
 (3.34)

Since $b - c \le 1 + a$ and $b \le 1$,

$$\frac{b-c}{1+a+b-c} \le \frac{1}{2} \le t_0 \le \alpha.$$
(3.35)

Let

$$\psi_{\alpha}(t) = \begin{cases} \frac{\alpha - 1}{\alpha} t + 1 & \text{if } 0 \le t \le \alpha, \\ t & \text{if } \alpha \le t \le 1. \end{cases}$$
(3.36)

We see that $\psi_{\alpha}(t_0) = \psi(t_0)$. By the convexity of ψ , we have

$$\psi(t) \le \psi_{\alpha}(t) \quad \forall t \le t_0. \tag{3.37}$$

Therefore,

$$\begin{aligned} \|x - y\| &= \left\| (a + 1, b - c) \right\|_{\psi} = (1 + a + b - c)\psi\left(\frac{b - c}{1 + a + b - c}\right) \\ &\leq (1 + a + b - c)\psi_{\alpha}\left(\frac{b - c}{1 + a + b - c}\right) = \frac{\alpha - 1}{\alpha}(b - c) + 1 + a + b - c \\ &= 1 + a + \frac{2\alpha - 1}{\alpha}b - \frac{2\alpha - 1}{\alpha}c = 1 + 1 - \frac{2\alpha - 1}{\alpha}c \\ &< 1 + 1 - \frac{2\alpha - 1}{\alpha}\frac{1/2}{\psi(1/2)} = 1 + \frac{1/2}{\psi(1/2)} + 1 - \frac{3\alpha - 1}{2\alpha}\frac{1}{\psi(1/2)} \\ &= 1 + \frac{1/2}{\psi(1/2)} + 1 - \frac{\psi_{\alpha}(1/2)}{\psi(1/2)} \leq 1 + \frac{1/2}{\psi(1/2)}. \end{aligned}$$
(3.38)

Finally, we conclude that

$$J(\ell_{\psi} - \ell_{\infty}) \le 1 + \frac{1/2}{\psi(1/2)}.$$
(3.39)

Now, we put $x_0 = ((1/2)/\psi(1/2), (1/2)/\psi(1/2))$ and $y_0 = (-1, 1)$, then

$$||x_0|| = ||y_0|| = 1, \qquad ||x_0 \pm y_0|| = 1 + \frac{1/2}{\psi(1/2)}.$$
 (3.40)

Thus,

$$J(\ell_{\psi} - \ell_{\infty}) \ge \min\{||x_0 - y_0||, ||x_0 + y_0||\} = 1 + \frac{1/2}{\psi(1/2)}.$$
(3.41)

This together with (3.39) completes the proof.

COROLLARY 3.8 [4, Example 2.4(2)]. Let $1 \le p \le \infty$, then

$$J(\ell_p - \ell_{\infty}) = 1 + \left(\frac{1}{2}\right)^{1/p}.$$
(3.42)

Indeed, $\psi_p(1/2) = 2^{1/p-1}$.

We now obtain the bounds for $J(\ell_{\psi} - \ell_1)$.

Corollary 3.9. Let $\psi \in \Psi_2$. Then

$$2\min_{0 \le t \le 1} \psi(t) \le J(\ell_{\psi} - \ell_1) \le \frac{3}{2} + \frac{1}{2}\min_{0 \le t \le 1} \psi(t).$$
(3.43)

Proof. Note that $\psi^*(1/2) = \max_{0 \le t \le 1} (1/2)/\psi(t) = 1/2 \min_{0 \le t \le 1} \psi(t)$. By Theorem 3.7, we have $J(\ell_{\psi^*} - \ell_{\infty}) = 1 + \min_{0 \le t \le 1} \psi(t)$. Applying Proposition 1.1(iv), the assertion is obtained.

We now improve the upper bound for $J(\ell_p - \ell_1)$ (see also Corollary 3.4).

COROLLARY 3.10. Let $1 \le p < \infty$. Then

$$J(\ell_p - \ell_1) \le \frac{3}{2} + \left(\frac{1}{2}\right)^{2 - 1/p}.$$
(3.44)

In particular, if $p \ge 2$, then

$$J(\ell_p - \ell_1) \le \min\left\{\frac{4}{(2^p + 2)^{1/p}}, \frac{3}{2} + \left(\frac{1}{2}\right)^{2-1/p}\right\}.$$
(3.45)

The following corollary follows by Theorem 3.7 and Corollary 3.9.

Corollary 3.11. Let $\psi \in \Psi_2$. Then

(i) $\ell_{\psi} - \ell_{\infty}$ is uniformly nonsquare if and only if $\psi \neq \psi_{\infty}$,

(ii) ℓ_{ψ} - ℓ_1 is uniformly nonsquare if and only if $\psi \neq \psi_1$.

We can say more about the uniform nonsquareness of ℓ_{ψ} - ℓ_{φ} .

THEOREM 3.12. Let $\psi, \varphi \in \Psi_2$. Then all $\ell_{\psi} - \ell_{\varphi}$ except $\ell_1 - \ell_1$ and $\ell_{\infty} - \ell_{\infty}$ are uniformly non-square.

Proof. If $\psi = \varphi$, we are done by [10, Corollary 3]. Assume that $\psi \neq \varphi$. We prove that $\ell_{\psi} - \ell_{\varphi}$ is uniformly nonsquare. Suppose not, that is, there are $x, y \in S_{\ell_{\psi} - \ell_{\varphi}}$ such that $||x \pm y||_{\psi,\varphi} = 2$. We consider three cases.

Case 1. $x, y \in Q_1$. Then

$$\begin{aligned} \|x\|_{\psi,1} &= \|x\|_{\psi} = \|x\|_{\psi,\varphi} = 1, \\ \|y\|_{\psi,1} &= \|y\|_{\psi} = \|y\|_{\psi,\varphi} = 1. \end{aligned} (3.46)$$

It follows by Lemma 3.2(i) that $x + y \in Q_1$ and $x - y \in Q_2 \cup Q_4$. Therefore

$$\|x+y\|_{\psi,1} = \|x+y\|_{\psi,\varphi} = 2,$$

$$2 = \|x-y\|_{\psi,\varphi} \le \|x-y\|_1 = \|x-y\|_{\psi,1} \le 2.$$
(3.47)

Hence $||x \pm y||_{\psi,1} = 2$ and this implies that $\ell_{\psi} - \ell_1$ is not uniformly nonsquare. By Corollary 3.11(ii), we have $\psi = \psi_1$. Again, since $\ell_{\psi} - \ell_{\varphi} = \ell_1 - \ell_{\varphi}$ is not uniformly nonsquare, $\varphi = \psi_1 = \psi$; a contradiction.

Case 2. $x, y \in Q_2$. It is similar to Case 1, so we omit the proof.

Case 3. $x := (a,b) \in Q_1$ and $y := (-c,d) \in Q_2$ where $a,b,c,d \in [0,1]$. Since $||x + y||_{\psi,\varphi} = 2$, the line segment joining *x* and *y* must lie in the sphere. In particular, there is $\alpha \in [0,1]$ such that

$$(0,1) = \alpha x + (1-\alpha)y. \tag{3.48}$$

It follows that b = 1 since $b, d \le 1$. Similarly consider x and -y instead of x and y, we can also conclude that a = 1. Hence $||(1,1)||_{\psi} = ||(1,1)||_{\psi,\varphi} = 1$, that is, $\psi(1/2) = 1/2$. Then $\psi = \psi_{\infty}$ and so $\ell_{\psi} - \ell_{\varphi} = \ell_{\infty} - \ell_{\varphi}$ is not uniformly nonsquare. By Corollary 3.11(i), we have $\varphi = \psi_{\infty} = \psi$; a contradiction.

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