# THE JAMES CONSTANT OF NORMALIZED NORMS ON $\mathbb{R}^{2}$ 

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We introduce a new class of normalized norms on $\mathbb{R}^{2}$ which properly contains all absolute normalized norms. We also give a criterion for deciding whether a given norm in this class is uniformly nonsquare. Moreover, an estimate for the James constant is presented and the exact value of some certain norms is computed. This gives a partial answer to the question raised by Kato et al.

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## 1. Introduction and preliminaries

A norm $\|\cdot\|$ on $\mathbb{C}^{2}$ (resp., $\left.\mathbb{R}^{2}\right)$ is said to be absolute if $\|(z, w)\|=\|(|z|,|w|)\|$ for all $z, w \in \mathbb{C}$ (resp., $\mathbb{R}$ ), and normalized if $\|(1,0)\|=\|(0,1)\|=1$. The $\ell_{p}$-norms $\|\cdot\|_{p}$ are such examples:

$$
\|(z, w)\|_{p}= \begin{cases}\left(|z|^{p}+|w|^{p}\right)^{1 / p} & \text { if } 1 \leq p<\infty  \tag{1.1}\\ \max \{|z|,|w|\} & \text { if } p=\infty\end{cases}
$$

Let $A N_{2}$ be the family of all absolute normalized norms on $\mathbb{C}^{2}$ (resp., $\mathbb{R}^{2}$ ), and $\Psi_{2}$ the family of all continuous convex functions $\psi$ on $[0,1]$ such that $\psi(0)=\psi(1)=1$ and $\max \{1-t, t\} \leq \psi(t) \leq 1(0 \leq t \leq 1)$. According to Bonsall and Duncan [1], $A N_{2}$ and $\Psi_{2}$ are in a one-to-one correspondence under the equation

$$
\begin{equation*}
\psi(t)=\|(1-t, t)\| \quad(0 \leq t \leq 1) . \tag{1.2}
\end{equation*}
$$

Indeed, for all $\psi \in \Psi_{2}$, let

$$
\|(z, w)\|_{\psi}= \begin{cases}(|z|+|w|) \psi\left(\frac{|w|}{|z|+|w|}\right) & \text { if }(z, w) \neq(0,0)  \tag{1.3}\\ 0 & \text { if }(z, w)=(0,0)\end{cases}
$$

Then $\|\cdot\|_{\psi} \in A N_{2}$, and $\|\cdot\|_{\psi}$ satisfies (1.2). From this result, we can consider many non- $\ell_{p}$-type norms easily. Now let

$$
\psi_{p}(t)= \begin{cases}\left((1-t)^{p}+t^{p}\right)^{1 / p} & \text { if } 1 \leq p<\infty,  \tag{1.4}\\ \max \{1-t, t\} & \text { if } p=\infty .\end{cases}
$$

Then $\psi_{p}(t) \in \Psi_{2}$ and, as is easily seen, the $\ell_{p}$-norm $\|\cdot\|_{p}$ is associated with $\psi_{p}$.
If $X$ is a Banach space, then $X$ is uniformly nonsquare if there exists $\delta \in(0,1)$ such that for any $x, y \in S_{X}$,

$$
\begin{equation*}
\text { either }\|x+y\| \leq 2(1-\delta) \quad \text { or } \quad\|x-y\| \leq 2(1-\delta) \tag{1.5}
\end{equation*}
$$

where $S_{X}=\{x \in X:\|x\|=1\}$. The James constant $J(X)$ is defined by

$$
\begin{equation*}
J(X)=\sup \left\{\min \{\|x+y\|,\|x-y\|\}: x, y \in S_{X}\right\} . \tag{1.6}
\end{equation*}
$$

The modulus of convexity of $X, \delta_{X}:[0,2] \rightarrow[0,1]$ is defined by

$$
\begin{equation*}
\delta_{X}(\varepsilon)=\inf \left\{1-\frac{1}{2}\|x+y\|: x, y \in S_{X},\|x-y\| \geq \varepsilon\right\} . \tag{1.7}
\end{equation*}
$$

The preceding parameters have been recently studied by several authors (cf. [4-6, 8, 9]). We collect together some known results.

Proposition 1.1. Let $X$ be a nontrivial Banach space, then
(i) $\sqrt{2} \leq J(X) \leq 2$ (Gao and Lau [5]),
(ii) if $X$ is a Hilbert space, then $J(X)=\sqrt{2}$; the converse is not true (Gao and Lau [5]),
(iii) $X$ is uniformly nonsquare if and only if $J(X)<2$ (Gao and Lau [5]),
(iv) $2 J(X)-2 \leq J\left(X^{*}\right) \leq J(X) / 2+1, J\left(X^{* *}\right)=J(X)$, and there exists a Banach space $X$ such that $J\left(X^{*}\right) \neq J(X)$ (Kato et al. [8]),
(v) if $2 \leq p \leq \infty$, then $\delta_{e_{p}}(\varepsilon)=1-\left(1-(\varepsilon / 2)^{p}\right)^{1 / p}$ (Hanner [6]),
(vi) $J(X)=\sup \left\{\varepsilon \in(0,2): \delta_{X}(\varepsilon) \leq 1-\varepsilon / 2\right\}$ (Gao and Lau [5]).

The paper is organized as follows. In Section 2 we introduce a new class of normalized norms on $\mathbb{R}^{2}$. This class properly contains all absolute normalized norms of Bonsall and Duncan [1]. The so-called generalized Day-James space, $\ell_{\psi}-\ell_{\varphi}$, where $\psi, \varphi \in \Psi_{2}$, is introduced and studied. More precisely, we prove that $\left(\ell_{\psi}-\ell_{\varphi}\right)^{*}=\ell_{\psi^{*}-} \ell_{\varphi^{*}}$ where $\psi^{*}$ and $\varphi^{*}$ are the dual functions of $\psi$ and $\varphi$, respectively. In Section 3, the upper bound of the James constant of the generalized Day-James space is given. Furthermore, we compute $J\left(\ell_{\psi}-\ell_{\infty}\right)$ and deduce that every generalized Day-James space except $\ell_{1}-\ell_{1}$ and $\ell_{\infty}-\ell_{\infty}$ is uniformly nonsquare. This result strengthens Corollary 3 of Saito et al. [10].

## 2. Generalized Day-James spaces

In this section, we introduce a new class of normalized norms on $\mathbb{R}^{2}$ which properly contains all absolute normalized norms of Bonsall and Duncan [1]. Moreover, we introduce a two-dimensional normed space which is a generalization of Day-James $\ell_{p}-\ell_{q}$ spaces.

Lemma 2.1. Let $\psi \in \Psi_{2}$ and let $\|\cdot\|_{\psi, \psi_{\infty}}$ be a function on $\mathbb{R}^{2}$ defined by, for all $(z, w) \in \mathbb{R}^{2}$,

$$
\begin{align*}
\|(z, w)\|_{\psi, \psi_{\infty}} & :=\max \left\{\left\|\left(z^{+}, w^{+}\right)\right\|_{\psi},\left\|\left(z^{-}, w^{-}\right)\right\|_{\psi}\right\}, \\
& = \begin{cases}\|(z, w)\|_{\psi} & \text { if } z w \geq 0, \\
\|(z, w)\|_{\infty} & \text { if } z w \leq 0,\end{cases} \tag{2.1}
\end{align*}
$$

where $x^{+}$and $x^{-}$are positive and negative parts of $x \in \mathbb{R}$, that is, $x^{+}=\max \{x, 0\}$ and $x^{-}=$ $\max \{-x, 0\}$. Then $\|\cdot\|_{\psi, \psi_{\infty}}$ is a norm on $\mathbb{R}^{2}$.

For convenience, we put $\mathscr{B}_{\psi_{1}, \psi_{2}}:=\left\{(z, w) \in \mathbb{R}^{2}:\|(z, w)\|_{\psi_{1}, \psi_{2}} \leq 1\right\}$.
Theorem 2.2. Let $\psi, \varphi \in \Psi_{2}$ and

$$
\|(z, w)\|_{\psi, \varphi}:= \begin{cases}\|(z, w)\|_{\psi} & \text { if } z w \geq 0  \tag{2.2}\\ \|(z, w)\|_{\varphi} & \text { if } z w \leq 0\end{cases}
$$

for all $(z, w) \in \mathbb{R}^{2}$. Then $\|\cdot\|_{\psi, \varphi}$ is a norm on $\mathbb{R}^{2}$. Denote by $N_{2}$ the family of all such preceding norms.
Proof. Let $\psi, \varphi \in \Psi_{2}$, we only show $\|\cdot\|_{\psi, \varphi}$ satisfies the triangle inequality. To this end, it suffices to prove that $\mathscr{B}_{\psi, \varphi}$ is convex. By Lemma 2.1, we have that $\mathscr{B}_{\psi, \psi_{\infty}}$ and $\mathscr{B}_{\varphi, \psi_{\infty}}$ are closed unit balls of $\|\cdot\|_{\psi, \psi_{\infty}}$ and $\|\cdot\|_{\varphi, \psi_{\infty}}$, respectively, and so $\mathscr{B}_{\psi, \psi_{\infty}}$ and $\mathscr{B}_{\varphi, \psi_{\infty}}$ are convex sets. We define $T: \mathbb{R}^{2} \rightarrow \mathbb{R}^{2}$ by

$$
\begin{equation*}
T((z, w))=(-z, w) \quad \forall(z, w) \in \mathbb{R}^{2} . \tag{2.3}
\end{equation*}
$$

Then $T$ is a linear operator and $T\left(\mathscr{B}_{\varphi, \psi_{\infty}}\right)=\mathscr{B}_{\psi_{\infty}, \varphi}$, which implies that $\mathscr{B}_{\psi_{\infty}, \varphi}$ is convex and so $\mathscr{B}_{\psi, \varphi}=\mathscr{B}_{\psi_{\infty}, \varphi} \cap \mathscr{B}_{\psi, \psi_{\infty}}$ is convex.

Taking $\psi=\psi_{p}$ and $\varphi=\psi_{q}(1 \leq p, q \leq \infty)$ in Theorem 2.2, we obtain the following.
Corollary 2.3 (Day-James $\ell_{p}-\ell_{q}$ spaces). For $1 \leq p, q \leq \infty$, denote by $\ell_{p}-\ell_{q}$ the DayJames space, that is, $\mathbb{R}^{2}$ with the norm defined by, for all $(z, w) \in \mathbb{R}^{2}$,

$$
\|(z, w)\|_{p, q}= \begin{cases}\|(z, w)\|_{p} & \text { if } z w \geq 0  \tag{2.4}\\ \|(z, w)\|_{q} & \text { if } z w \leq 0\end{cases}
$$

James [7] considered the $\ell_{p}-\ell_{p^{\prime}}$ space as an example of a Banach space which is isometric to its dual but which is not given by a Hilbert norm when $p \neq 2$. Day [2] considered even more general spaces, namely, if $(X,\|\cdot\|)$ is a two-dimensional Banach space and $\left(X^{*},\|\cdot\|^{*}\right)$ its dual, then the $X-X^{*}$ space is the space $X$ with the norm defined by, for all $(z, w) \in \mathbb{R}^{2}$,

$$
\|(z, w)\|_{X_{, X^{*}}}= \begin{cases}\|(z, w)\|^{\prime} & \text { if } z w \geq 0  \tag{2.5}\\ \|(z, w)\|^{*} & \text { if } z w \leq 0\end{cases}
$$

For $\psi, \varphi \in \Psi_{2}$, denote by $\ell_{\psi}-\ell_{\varphi}$ the generalized Day-James space, that is, $\mathbb{R}^{2}$ with the norm $\|\cdot\|_{\psi, \varphi}$ defined by (2.2). For $\psi_{p}$ defined by (1.4), we write $\ell_{\psi}-\ell_{p}$ for $\ell_{\psi}-\ell_{\psi_{p}}$. For example, if $1 \leq p, q \leq \infty, \ell_{p}-\ell_{q}$ means $\ell_{\psi_{p}}-\ell_{\psi_{q}}$.

It is worthwhile to mention that there is a normalized norm which is not absolute.
Proposition 2.4. There is $\psi \in \Psi_{2}$ such that $\ell_{\psi}-\ell_{\infty}$ is not isometrically isomorphic to $\ell_{\varphi}-\ell_{\varphi}$ for all $\varphi \in \Psi_{2}$.

Proof. Let

$$
\psi(t):= \begin{cases}1-t & \text { if } 0 \leq t \leq \frac{1}{8}  \tag{2.6}\\ \frac{11-4 t}{12} & \text { if } \frac{1}{8} \leq t \leq \frac{1}{2} \\ \frac{1+t}{2} & \text { if } \frac{1}{2} \leq t \leq 1\end{cases}
$$

We observe that the sphere of $\ell_{\psi}-\ell_{\infty}$ is the octagon whose right half consists of 4 segments of different lengths. Suppose that there are $\varphi \in \Psi_{2}$ and an isometric isomorphism from $\ell_{\psi}-\ell_{\infty}$ onto $\ell_{\varphi}-\ell_{\varphi}$. Since the image of each segment in $\ell_{\psi}-\ell_{\infty}$ is again a segment of the same length in $\ell_{\varphi}-\ell_{\varphi}$, the sphere of $\ell_{\varphi}-\ell_{\varphi}$ must be the octagon whose each corresponding side has the same length (measured by $\|\cdot\|_{\varphi}$ ). We show that this cannot happen. Consider $(1,0) \in S_{\ell_{\varphi}-\ell_{\varphi}}$. If $(1,0)$ is an extreme point of $B_{\ell_{\varphi}-\ell_{\varphi}}$, then $S_{\ell_{\varphi}-\ell_{\varphi}}$ contains 4 segments of same lengths since $\|\cdot\|_{\varphi}$ is absolute. On the other hand, if $(1,0)$ is an not extreme point of $B_{\ell_{\varphi}-\ell_{\varphi}}$, again $S_{\ell_{\varphi}-\ell_{\varphi}}$ contains 4 segments of same lengths.

Next, we prove that the dual of a generalized Day-James space is again a generalized Day-James space. Recall that, for $\psi \in \Psi_{2}$, the dual function $\psi^{*}$ of $\psi$ is defined by

$$
\begin{equation*}
\psi^{*}(s)=\max _{0 \leq t \leq 1} \frac{(1-s)(1-t)+s t}{\psi(t)} \tag{2.7}
\end{equation*}
$$

for all $s \in[0,1]$. It was proved that $\psi^{*} \in \Psi_{2}$ and $\left(\ell_{\psi^{-}} \ell_{\psi}\right)^{*}=\ell_{\psi^{*}} \ell_{\psi^{*}}$ (see [3, Proposition 1 and Theorem 2]). We generalize this result to our spaces as follows.

Theorem 2.5. For $\psi, \varphi \in \Psi_{2}$, there is an isometric isomorphism that identifies $\left(\ell_{\psi}-\ell_{\varphi}\right)^{*}$ with $\ell_{\psi^{*}}-\ell_{\varphi^{*}}$ such that if $f \in\left(\ell_{\psi}-\ell_{\varphi}\right)^{*}$ is identified with the element $(z, w) \in \ell_{\psi^{*}}-\ell_{\varphi^{*}}$, then

$$
\begin{equation*}
f(u, v)=z u+w v \tag{2.8}
\end{equation*}
$$

for all $(u, v) \in \mathbb{R}^{2}$.
Proof. We can prove analogous to [3, Theorem 2].

## 3. The James constant and uniform nonsquareness

The next lemmas are crucial for proving the main theorems.
Lemma 3.1. Let $\psi, \varphi \in \Psi_{2}$. Then
(i) $\|\cdot\|_{\infty} \leq\|\cdot\|_{\psi, \varphi} \leq\|\cdot\|_{1}$,
(ii) $\left(1 / M_{\psi, \varphi}\right)\|\cdot\|_{\psi} \leq\|\cdot\|_{\psi, \varphi} \leq M_{\varphi, \psi}\|\cdot\|_{\psi}$,
(iii) $\left(1 / M_{\varphi, \psi}\right)\|\cdot\|_{\varphi} \leq\|\cdot\|_{\psi, \varphi} \leq M_{\psi, \varphi}\|\cdot\|_{\varphi}$,
where $M_{\varphi, \psi}=\max _{0 \leq t \leq 1} \varphi(t) / \psi(t)$ and $M_{\psi, \varphi}=\max _{0 \leq t \leq 1} \psi(t) / \varphi(t)$.
Lemma 3.2. Let $\psi, \varphi \in \Psi_{2}$ and let $Q_{i}(i=1, \ldots, 4)$ denote the ith quadrant in $\mathbb{R}^{2}$. Suppose that $x, y \in S_{e_{\psi}-\ell_{\varphi}}$, then the following statements are true.
(i) If $x, y \in Q_{1}$, then $x+y \in Q_{1}$ and $x-y \in Q_{2} \cup Q_{4}$.
(ii) If $x, y \in Q_{2}$, then $x+y \in Q_{2}$ and $x-y \in Q_{1} \cup Q_{3}$.
(iii) If $\psi(t) \leq \varphi(t)$ for all $t \in[0,1]$ and $x-y \in Q_{2}^{\circ} \cup Q_{4}^{\circ}$, where $Q_{2}^{\circ}$ and $Q_{4}^{\circ}$ are the interiors of $Q_{2}$ and $Q_{4}$, respectively, then $x+y \in Q_{1} \cup Q_{3}$.

We will estimate the James constant of $\ell_{\psi}-\ell_{\varphi}$.
Theorem 3.3. Let $\psi, \varphi \in \Psi_{2}$ with $\psi(t) \leq \varphi(t)$ for all $t \in[0,1]$, let $M_{\varphi, \psi}=\max _{0 \leq t \leq 1} \varphi(t) / \psi(t)$, and let $\delta_{\psi}(\cdot)$ be the modulus of convexity of $\ell_{\psi}-\ell_{\psi}$. Then for $\varepsilon \in[0,2]$,

$$
\begin{equation*}
\delta_{\psi, \varphi}(\varepsilon) \geq \min \left\{1-M_{\varphi, \psi}\left(1-\delta_{\psi}(\varepsilon)\right), \delta_{\psi}\left(\frac{\varepsilon}{M_{\varphi, \psi}}\right)\right\}, \tag{3.1}
\end{equation*}
$$

where $\delta_{\psi, \varphi}(\cdot)$ is the modulus of convexity of $\ell_{\psi}-\ell_{\varphi}$. Consequently,

$$
\begin{equation*}
J\left(\ell_{\psi}-\ell_{\varphi}\right) \leq \sup \left\{\varepsilon \in(0,2): \varepsilon \leq 2 M_{\varphi, \psi}\left(1-\delta_{\psi}(\varepsilon)\right) \text { or } \varepsilon \leq 2\left(1-\delta_{\psi}\left(\frac{\varepsilon}{M_{\varphi, \psi}}\right)\right)\right\} . \tag{3.2}
\end{equation*}
$$

Proof. By Lemma 3.1(ii), we have

$$
\begin{equation*}
\|\cdot\|_{\psi} \leq\|\cdot\|_{\psi, \varphi} \leq M_{\varphi, \psi}\|\cdot\|_{\psi} . \tag{3.3}
\end{equation*}
$$

We now evaluate the modulus of convexity $\delta_{\psi, \varphi}$ for $\ell_{\psi}-\ell_{\varphi}$. We consider two cases.
Case 1. Take $\|x\|_{\psi, \varphi}=\|y\|_{\psi, \varphi}=1$ with $\|x-y\|_{\psi, \varphi} \geq \varepsilon$, where $x-y \in Q_{1} \cup Q_{3}$. Thus $\|x\|_{\psi} \leq 1,\|y\|_{\psi} \leq 1$, and $\|x-y\|_{\psi} \geq \varepsilon$, which implies that

$$
\begin{equation*}
\frac{1}{2}\|x+y\|_{\psi} \leq 1-\delta_{\psi}(\varepsilon) \tag{3.4}
\end{equation*}
$$

This in turn implies

$$
\begin{equation*}
\frac{1}{2}\|x+y\|_{\psi, \varphi} \leq \frac{1}{2} M_{\varphi, \psi}\|x+y\|_{\psi} \leq M_{\varphi, \psi}\left(1-\delta_{\psi}(\varepsilon)\right), \tag{3.5}
\end{equation*}
$$

thus

$$
\begin{equation*}
1-\frac{1}{2}\|x+y\|_{\psi, \varphi} \geq 1-M_{\varphi, \psi}\left(1-\delta_{\psi}(\varepsilon)\right) . \tag{3.6}
\end{equation*}
$$

Case 2. Now take $x, y$ as above, but with $x-y \in Q_{2}^{\circ} \cup Q_{4}^{\circ}$. By Lemma 3.2(iii), $x+y \in$ $Q_{1} \cup Q_{3}$. Since $\|x-y\|_{\psi, \varphi} \geq \varepsilon$,

$$
\begin{equation*}
\|x-y\|_{\psi} \geq \frac{\|x-y\|_{\psi, \varphi}}{M_{\varphi, \psi}} \geq \frac{\varepsilon}{M_{\varphi, \psi}} . \tag{3.7}
\end{equation*}
$$

Then

$$
\begin{equation*}
\frac{1}{2}\|x+y\|_{\psi, \varphi}=\frac{1}{2}\|x+y\|_{\psi} \leq 1-\delta_{\psi}\left(\frac{\varepsilon}{M_{\varphi, \psi}}\right) \tag{3.8}
\end{equation*}
$$

and so

$$
\begin{equation*}
1-\frac{1}{2}\|x+y\|_{\psi, \varphi} \geq \delta_{\psi}\left(\frac{\varepsilon}{M_{\varphi, \psi}}\right) . \tag{3.9}
\end{equation*}
$$

Hence we obtain (3.1). By Proposition 1.1(vi), (3.2) follows.
The following corollary shows that we can have equality in (3.2).
Corollary 3.4 [4, 8]. If $1 \leq q \leq p<\infty$ and $p \geq 2$, then

$$
\begin{equation*}
J\left(\ell_{p}-\ell_{q}\right) \leq 2\left(\frac{2^{p / q}}{2^{p / q}+2}\right)^{1 / p} \tag{3.10}
\end{equation*}
$$

In particular, if $p=2$ and $q=1$, then $J\left(\ell_{2}-\ell_{1}\right)=\sqrt{8 / 3}$.
Proof. It follows that since

$$
\begin{equation*}
M_{\psi_{q}, \psi_{p}}=2^{1 / q-1 / p}, \quad \delta_{\ell_{p}-\ell_{p}}(\varepsilon)=1-\left(1-\left(\frac{\varepsilon}{2}\right)^{p}\right)^{1 / p} \tag{3.11}
\end{equation*}
$$

Moreover, if $p=2$ and $q=1$, then $J\left(\ell_{2}-\ell_{1}\right) \leq \sqrt{8 / 3}$. Now we put

$$
\begin{equation*}
x_{0}=\left(\frac{2+\sqrt{2}}{2 \sqrt{3}}, \frac{2-\sqrt{2}}{2 \sqrt{3}}\right), \quad y_{0}=\left(\frac{2-\sqrt{2}}{2 \sqrt{3}}, \frac{2+\sqrt{2}}{2 \sqrt{3}}\right) . \tag{3.12}
\end{equation*}
$$

Then

$$
\begin{equation*}
\left\|x_{0}\right\|_{2,1}=\left\|y_{0}\right\|_{2,1}=1, \quad\left\|x_{0} \pm y_{0}\right\|_{2,1}=\sqrt{\frac{8}{3}} . \tag{3.13}
\end{equation*}
$$

Theorem 3.5. Let $\psi, \varphi \in \Psi_{2}$ with $\psi(t) \leq \varphi(t)$ for all $t \in[0,1]$, let $M_{\varphi, \psi}=\max _{0 \leq t \leq 1} \varphi(t) / \psi(t)$, and let $\delta_{\varphi}(\cdot)$ be the modulus of convexity of $\ell_{\varphi}-\ell_{\varphi}$. Then for $\varepsilon \in[0,2]$,

$$
\begin{equation*}
\delta_{\psi, \varphi}(\varepsilon) \geq 1-M_{\varphi, \psi}\left(1-\delta_{\varphi}\left(\frac{\varepsilon}{M_{\varphi, \psi}}\right)\right), \tag{3.14}
\end{equation*}
$$

where $\delta_{\psi, \varphi}(\cdot)$ is the modulus of convexity of $\ell_{\psi}-\ell_{\varphi}$. Consequently,

$$
\begin{equation*}
J\left(\ell_{\psi}-\ell_{\varphi}\right) \leq \sup \left\{\varepsilon \in(0,2): \varepsilon \leq 2 M_{\varphi, \psi}\left(1-\delta_{\varphi}\left(\frac{\varepsilon}{M_{\varphi, \psi}}\right)\right)\right\} . \tag{3.15}
\end{equation*}
$$

Proof. By Lemma 3.1(iii), we have

$$
\begin{equation*}
\frac{1}{M_{\varphi, \psi}}\|\cdot\|_{\varphi} \leq\|\cdot\|_{\psi, \varphi} \leq\|\cdot\|_{\varphi} . \tag{3.16}
\end{equation*}
$$

We now evaluate the modulus of convexity $\delta_{\psi, \varphi}$ for $\ell_{\psi}-\ell_{\varphi}$. Let

$$
\begin{equation*}
\|x\|_{\psi, \varphi}=\|y\|_{\psi, \varphi}=1 \quad \text { with }\|x-y\|_{\psi, \varphi} \geq \varepsilon . \tag{3.17}
\end{equation*}
$$

Then

$$
\begin{align*}
\frac{1}{M_{\varphi, \psi}}\|x\|_{\varphi} \leq 1, \quad \frac{1}{M_{\varphi, \psi}}\|y\|_{\varphi} \leq 1, \\
\frac{1}{M_{\varphi, \psi}}\|x-y\|_{\varphi} \geq \frac{1}{M_{\varphi, \psi}}\|x-y\|_{\psi, \varphi} \geq \frac{\varepsilon}{M_{\varphi, \psi}}, \tag{3.18}
\end{align*}
$$

which implies that

$$
\begin{equation*}
\frac{1}{2 M_{\varphi, \psi}}\|x+y\|_{\varphi} \leq 1-\delta_{\varphi}\left(\frac{\varepsilon}{M_{\varphi, \psi}}\right) . \tag{3.19}
\end{equation*}
$$

This in turn implies that

$$
\begin{equation*}
\frac{1}{2 M_{\varphi, \psi}}\|x+y\|_{\psi, \varphi} \leq \frac{1}{2 M_{\varphi, \psi}}\|x+y\|_{\varphi} \leq 1-\delta_{\varphi}\left(\frac{\varepsilon}{M_{\varphi, \psi}}\right) \tag{3.20}
\end{equation*}
$$

thus

$$
\begin{equation*}
1-\frac{1}{2}\|x+y\|_{\psi, \varphi} \geq 1-M_{\varphi, \psi}\left(1-\delta_{\varphi}\left(\frac{\varepsilon}{M_{\varphi, \psi}}\right)\right) \tag{3.21}
\end{equation*}
$$

Hence we obtain (3.14). By Proposition 1.1(vi), (3.15) follows.
Corollary 3.6. If $2 \leq q \leq p<\infty$, then

$$
\begin{equation*}
J\left(\ell_{p}-\ell_{q}\right) \leq 2^{1-1 / p} . \tag{3.22}
\end{equation*}
$$

It is easy to see that the estimate (3.22) is better than one obtained in [4, Example 2.4(3)].

For some generalized Day-James spaces, [8, Corollary 4] of Kato et al. gives only rough result for the estimate of the James constant, that is, for $\psi \in \Psi_{2}$,

$$
\begin{equation*}
\frac{2}{M} \leq J\left(\ell_{\psi}-\ell_{\infty}\right) \leq 2 M \tag{3.23}
\end{equation*}
$$

where $M=\max _{0 \leq t \leq 1} \psi_{\infty}(t) / \psi(t)$.
However, the following theorem gives the exact value of the James constant of these spaces.

Theorem 3.7. Let $\psi \in \Psi_{2}$. Then

$$
\begin{equation*}
J\left(\ell_{\psi}-\ell_{\infty}\right)=1+\frac{1 / 2}{\psi(1 / 2)} . \tag{3.24}
\end{equation*}
$$

Proof. For our convenience, we write $\|\cdot\|$ instead of $\|\cdot\|_{\psi, \psi_{\infty}}$. Let $x, y \in S_{\ell_{\psi}-e_{\infty}}$. We prove that

$$
\begin{equation*}
\text { either }\|x+y\| \leq 1+\frac{1 / 2}{\psi(1 / 2)} \quad \text { or } \quad\|x-y\| \leq 1+\frac{1 / 2}{\psi(1 / 2)} . \tag{3.25}
\end{equation*}
$$

Let us consider the following cases.
Case 1. $x, y \in Q_{1}$. Let $x=(a, b)$ and $y=(c, d)$ where $a, b, c, d \in[0,1]$. By Lemma 3.2(i), we have $x-y \in Q_{2} \cup Q_{4}$. Then

$$
\begin{equation*}
\|x-y\|=\max \{|a-c|,|b-d|\} \leq 1 \leq 1+\frac{1 / 2}{\psi(1 / 2)} \tag{3.26}
\end{equation*}
$$

Case 2. $x, y \in Q_{2}$. If $x, y$ lies in the same segment, then $\|x-y\| \leq 1$. We now suppose that $x=(-1, a)$ and $y=(-c, 1)$ where $a, c \in[0,1]$.

Subcase 2.1. $a \leq(1 / 2) / \psi(1 / 2)$ and $c \leq(1 / 2) / \psi(1 / 2)$. Then

$$
\begin{equation*}
\|x+y\|=\|(-1-c, 1+a)\|_{\infty}=\max \{1+c, 1+a\} \leq 1+\frac{1 / 2}{\psi(1 / 2)} \tag{3.27}
\end{equation*}
$$

Subcase 2.2. $a \geq(1 / 2) / \psi(1 / 2)$ or $c \geq(1 / 2) / \psi(1 / 2)$. Put $z=(-1,1)$, then

$$
\begin{equation*}
\|x-y\| \leq\|x-z\|+\|z-y\|=1-a+1-c \leq 1+1-\frac{1 / 2}{\psi(1 / 2)} \leq 1+\frac{1 / 2}{\psi(1 / 2)} \tag{3.28}
\end{equation*}
$$

From now on, we may assume without loss of generality that there is $\beta \in[1 / 2,1]$ such that $\psi(\beta) \leq \psi(t)$ for all $t \in[0,1]$. Indeed, $J\left(\ell_{\psi}-\ell_{\infty}\right)=J\left(\ell_{\tilde{\psi}}-\ell_{\infty}\right)$ where $\tilde{\psi}(t)=\psi(1-t)$ for all $t \in[0,1]$.

Case 3. $x \in Q_{1}$ and $y \in Q_{2}$. Let $x=(a, b), y=(-c, 1)$ where $a, b, c \in[0,1]$. We consider three subcases.

Subcase 3.1. $a \leq(1 / 2) / \psi(1 / 2)$ or $c \leq(1 / 2) / \psi(1 / 2)$. Then

$$
\begin{equation*}
\|x-y\|=\|(a+c, b-1)\|_{\infty}=\max \{a+c, 1-b\} \leq 1+\frac{1 / 2}{\psi(1 / 2)} . \tag{3.29}
\end{equation*}
$$

Subcase 3.2. $(1 / 2) / \psi(1 / 2) \leq a \leq c$. Then $b \leq(1 / 2) / \psi(1 / 2)$ and

$$
\begin{equation*}
\|x+y\|=\|(a-c, b+1)\|_{\infty}=\max \{c-a, 1+b\} \leq 1+\frac{1 / 2}{\psi(1 / 2)} \tag{3.30}
\end{equation*}
$$

Subcase 3.3. $(1 / 2) / \psi(1 / 2)<c \leq a$. We write $a=\left(1-t_{0}\right) / \psi\left(t_{0}\right), b=t_{0} / \psi\left(t_{0}\right)$ where $t_{0}=$ $b /(a+b)$ and $0 \leq t_{0} \leq 1 / 2$. By the convexity of $\psi$ and $\psi(t) \geq \psi(\beta)$ for all $0 \leq t \leq 1$, we
have $\psi\left(t_{0}\right) \geq \psi(1 / 2)$ and so $1 / \psi\left(t_{0}\right) \leq 1 / \psi(1 / 2)$. By Lemma 3.1(i),

$$
\begin{align*}
\|x+y\| & =\|(a, b)+(-c, 1)\| \leq\|(a-c, b+1)\|_{1} \\
& =a-c+b+1=\frac{1}{\psi\left(t_{0}\right)}+1-c  \tag{3.31}\\
& \leq \frac{1}{\psi(1 / 2)}+1-\frac{1 / 2}{\psi(1 / 2)}=1+\frac{1 / 2}{\psi(1 / 2)} .
\end{align*}
$$

Case 4. $x \in Q_{1}$ and $y \in Q_{2}$. Let $x=(a, b), y=(-1, c)$ where $a, b, c \in[0,1]$. We consider three subcases.

Subcase 4.1. $b \leq(1 / 2) / \psi(1 / 2)$ or $c \leq(1 / 2) / \psi(1 / 2)$. Then

$$
\begin{equation*}
\|x+y\|=\|(a-1, b+c)\|_{\infty}=\max \{1-a, b+c\} \leq 1+\frac{1 / 2}{\psi(1 / 2)} . \tag{3.32}
\end{equation*}
$$

Subcase 4.2. $(1 / 2) / \psi(1 / 2)<b \leq c$. Then $a \leq(1 / 2) / \psi(1 / 2)$ and

$$
\begin{equation*}
\|x-y\|=\|(1+a, b-c)\|_{\infty}=\max \{1+a, c-b\} \leq 1+\frac{1 / 2}{\psi(1 / 2)} \tag{3.33}
\end{equation*}
$$

Subcase 4.3. $(1 / 2) / \psi(1 / 2)<c \leq b$. We write $a=\left(1-t_{0}\right) / \psi\left(t_{0}\right), b=t_{0} / \psi\left(t_{0}\right)$, where $t_{0}=$ $b /(a+b)$ and $1 / 2 \leq t_{0} \leq 1$. We choose $\alpha=b /(a+2 b-1)$, then

$$
\begin{equation*}
\frac{1}{2} \leq \alpha \leq 1, \quad a=\frac{1-2 \alpha}{\alpha} b+1 \tag{3.34}
\end{equation*}
$$

Since $b-c \leq 1+a$ and $b \leq 1$,

$$
\begin{equation*}
\frac{b-c}{1+a+b-c} \leq \frac{1}{2} \leq t_{0} \leq \alpha . \tag{3.35}
\end{equation*}
$$

Let

$$
\psi_{\alpha}(t)= \begin{cases}\frac{\alpha-1}{\alpha} t+1 & \text { if } 0 \leq t \leq \alpha  \tag{3.36}\\ t & \text { if } \alpha \leq t \leq 1\end{cases}
$$

We see that $\psi_{\alpha}\left(t_{0}\right)=\psi\left(t_{0}\right)$. By the convexity of $\psi$, we have

$$
\begin{equation*}
\psi(t) \leq \psi_{\alpha}(t) \quad \forall t \leq t_{0} . \tag{3.37}
\end{equation*}
$$

Therefore,

$$
\begin{align*}
\|x-y\| & =\|(a+1, b-c)\|_{\psi}=(1+a+b-c) \psi\left(\frac{b-c}{1+a+b-c}\right) \\
& \leq(1+a+b-c) \psi_{\alpha}\left(\frac{b-c}{1+a+b-c}\right)=\frac{\alpha-1}{\alpha}(b-c)+1+a+b-c \\
& =1+a+\frac{2 \alpha-1}{\alpha} b-\frac{2 \alpha-1}{\alpha} c=1+1-\frac{2 \alpha-1}{\alpha} c  \tag{3.38}\\
& <1+1-\frac{2 \alpha-1}{\alpha} \frac{1 / 2}{\psi(1 / 2)}=1+\frac{1 / 2}{\psi(1 / 2)}+1-\frac{3 \alpha-1}{2 \alpha} \frac{1}{\psi(1 / 2)} \\
& =1+\frac{1 / 2}{\psi(1 / 2)}+1-\frac{\psi_{\alpha}(1 / 2)}{\psi(1 / 2)} \leq 1+\frac{1 / 2}{\psi(1 / 2)} .
\end{align*}
$$

Finally, we conclude that

$$
\begin{equation*}
J\left(\ell_{\psi}-\ell_{\infty}\right) \leq 1+\frac{1 / 2}{\psi(1 / 2)} . \tag{3.3}
\end{equation*}
$$

Now, we put $x_{0}=((1 / 2) / \psi(1 / 2),(1 / 2) / \psi(1 / 2))$ and $y_{0}=(-1,1)$, then

$$
\begin{equation*}
\left\|x_{0}\right\|=\left\|y_{0}\right\|=1, \quad\left\|x_{0} \pm y_{0}\right\|=1+\frac{1 / 2}{\psi(1 / 2)} . \tag{3.40}
\end{equation*}
$$

Thus,

$$
\begin{equation*}
J\left(\ell_{\psi}-\ell_{\infty}\right) \geq \min \left\{\left\|x_{0}-y_{0}\right\|,\left\|x_{0}+y_{0}\right\|\right\}=1+\frac{1 / 2}{\psi(1 / 2)} \tag{3.41}
\end{equation*}
$$

This together with (3.39) completes the proof.
Corollary 3.8 [4, Example 2.4(2)]. Let $1 \leq p \leq \infty$, then

$$
\begin{equation*}
J\left(\ell_{p}-\ell_{\infty}\right)=1+\left(\frac{1}{2}\right)^{1 / p} . \tag{3.42}
\end{equation*}
$$

Indeed, $\psi_{p}(1 / 2)=2^{1 / p-1}$.
We now obtain the bounds for $J\left(\ell_{\psi}-\ell_{1}\right)$.
Corollary 3.9. Let $\psi \in \Psi_{2}$. Then

$$
\begin{equation*}
2 \min _{0 \leq t \leq 1} \psi(t) \leq J\left(\ell_{\psi}-\ell_{1}\right) \leq \frac{3}{2}+\frac{1}{2} \min _{0 \leq t \leq 1} \psi(t) . \tag{3.43}
\end{equation*}
$$

Proof. Note that $\psi^{*}(1 / 2)=\max _{0 \leq t \leq 1}(1 / 2) / \psi(t)=1 / 2 \min _{0 \leq t \leq 1} \psi(t)$. By Theorem 3.7, we have $J\left(\ell_{\psi^{*}-\ell_{\infty}}\right)=1+\min _{0 \leq t \leq 1} \psi(t)$. Applying Proposition 1.1(iv), the assertion is obtained.

We now improve the upper bound for $J\left(\ell_{p}-\ell_{1}\right)$ (see also Corollary 3.4).

Corollary 3.10. Let $1 \leq p<\infty$. Then

$$
\begin{equation*}
J\left(\ell_{p}-\ell_{1}\right) \leq \frac{3}{2}+\left(\frac{1}{2}\right)^{2-1 / p} . \tag{3.44}
\end{equation*}
$$

In particular, if $p \geq 2$, then

$$
\begin{equation*}
J\left(\ell_{p}-\ell_{1}\right) \leq \min \left\{\frac{4}{\left(2^{p}+2\right)^{1 / p}}, \frac{3}{2}+\left(\frac{1}{2}\right)^{2-1 / p}\right\} . \tag{3.45}
\end{equation*}
$$

The following corollary follows by Theorem 3.7 and Corollary 3.9.
Corollary 3.11. Let $\psi \in \Psi_{2}$. Then
(i) $\ell_{\psi}-\ell_{\infty}$ is uniformly nonsquare if and only if $\psi \neq \psi_{\infty}$,
(ii) $\ell_{\psi}-\ell_{1}$ is uniformly nonsquare if and only if $\psi \neq \psi_{1}$.

We can say more about the uniform nonsquareness of $\ell_{\psi}-\ell_{\varphi}$.
Theorem 3.12. Let $\psi, \varphi \in \Psi_{2}$. Then all $\ell_{\psi}-\ell_{\varphi}$ except $\ell_{1}-\ell_{1}$ and $\ell_{\infty}-\ell_{\infty}$ are uniformly nonsquare.

Proof. If $\psi=\varphi$, we are done by [10, Corollary 3]. Assume that $\psi \neq \varphi$. We prove that $\ell_{\psi}-\ell_{\varphi}$ is uniformly nonsquare. Suppose not, that is, there are $x, y \in S_{e_{\psi}-\ell_{\varphi}}$ such that $\|x \pm y\|_{\psi, \varphi}=$ 2. We consider three cases.

Case 1. $x, y \in Q_{1}$. Then

$$
\begin{align*}
& \|x\|_{\psi, 1}=\|x\|_{\psi}=\|x\|_{\psi, \varphi}=1, \\
& \|y\|_{\psi, 1}=\|y\|_{\psi}=\|y\|_{\psi, \varphi}=1 . \tag{3.46}
\end{align*}
$$

It follows by Lemma 3.2(i) that $x+y \in Q_{1}$ and $x-y \in Q_{2} \cup Q_{4}$. Therefore

$$
\begin{gather*}
\|x+y\|_{\psi, 1}=\|x+y\|_{\psi, \varphi}=2, \\
2=\|x-y\|_{\psi, \varphi} \leq\|x-y\|_{1}=\|x-y\|_{\psi, 1} \leq 2 . \tag{3.47}
\end{gather*}
$$

Hence $\|x \pm y\|_{\psi, 1}=2$ and this implies that $\ell_{\psi}-\ell_{1}$ is not uniformly nonsquare. By Corollary 3.11(ii), we have $\psi=\psi_{1}$. Again, since $\ell_{\psi}-\ell_{\varphi}=\ell_{1}-\ell_{\varphi}$ is not uniformly nonsquare, $\varphi=\psi_{1}=$ $\psi$; a contradiction.

Case 2. $x, y \in Q_{2}$. It is similar to Case 1, so we omit the proof.
Case 3. $x:=(a, b) \in Q_{1}$ and $y:=(-c, d) \in Q_{2}$ where $a, b, c, d \in[0,1]$. Since $\|x+y\|_{\psi, \varphi}=2$, the line segment joining $x$ and $y$ must lie in the sphere. In particular, there is $\alpha \in[0,1]$ such that

$$
\begin{equation*}
(0,1)=\alpha x+(1-\alpha) y . \tag{3.48}
\end{equation*}
$$

It follows that $b=1$ since $b, d \leq 1$. Similarly consider $x$ and $-y$ instead of $x$ and $y$, we can also conclude that $a=1$. Hence $\|(1,1)\|_{\psi}=\|(1,1)\|_{\psi, \varphi}=1$, that is, $\psi(1 / 2)=1 / 2$. Then $\psi=\psi_{\infty}$ and so $\ell_{\psi}-\ell_{\varphi}=\ell_{\infty}-\ell_{\varphi}$ is not uniformly nonsquare. By Corollary 3.11(i), we have $\varphi=\psi_{\infty}=\psi$; a contradiction.

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