# ON THE EXISTENCE OF SOLUTIONS TO A FOURTH-ORDER QUASILINEAR RESONANT PROBLEM 

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By means of Morse theory we prove the existence of a nontrivial solution to a superlinear $p$-harmonic elliptic problem with Navier boundary conditions having a linking structure around the origin. Moreover, in case of both resonance near zero and nonresonance at $+\infty$ the existence of two nontrivial solutions is shown.

## 1. Introduction and main results

Let $p>1$ and $\Omega \subset \mathbb{R}^{n}$ be a smooth bounded domain with $n \geqslant 2 p+1$. We are concerned with the existence of nontrivial solutions to the $p$-harmonic equation

$$
\begin{equation*}
\Delta\left(|\Delta u|^{p-2} \Delta u\right)=g(x, u) \quad \text { in } \Omega \tag{1.1}
\end{equation*}
$$

with Navier boundary conditions

$$
\begin{equation*}
u=\Delta u=0 \quad \text { on } \partial \Omega, \tag{1.2}
\end{equation*}
$$

where $g: \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ is a Carathéodory function such that for some $C>0$,

$$
\begin{equation*}
|g(x, s)| \leqslant C\left(1+|s|^{q-1}\right) \tag{1.3}
\end{equation*}
$$

for a.e. $x \in \Omega$ and all $s \in \mathbb{R}$, being $1 \leqslant q<p_{*}$ and $p_{*}=n p /(n-2 p)$.
It is well known that the functional $\Phi: W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega) \rightarrow \mathbb{R}$

$$
\begin{equation*}
\Phi(u)=\frac{1}{p} \int_{\Omega}|\Delta u|^{p} d x-\int_{\Omega} G(x, u) d x, \tag{1.4}
\end{equation*}
$$

with $G(x, s)=\int_{0}^{s} g(x, t) d t$, is of class $C^{1}$ and

$$
\begin{equation*}
\left\langle\Phi^{\prime}(u), \varphi\right\rangle=\int_{\Omega}|\Delta u|^{p-2} \Delta u \Delta \varphi d x-\int_{\Omega} g(x, u) \varphi d x \tag{1.5}
\end{equation*}
$$

for each $\varphi \in W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega)$. Moreover, the critical points of $\Phi$ are weak solutions for (1.1). Notice that for the eigenvalue problem

$$
\begin{equation*}
\Delta\left(|\Delta u|^{p-2} \Delta u\right)=\lambda|u|^{p-2} u \quad \text { in } \Omega \tag{1.6}
\end{equation*}
$$

with boundary data (1.2), as for the $p$-Laplacian eigenvalue problem with Dirichlet boundary data,

$$
\begin{equation*}
\lambda_{n}=\inf _{A \in I_{n}} \sup _{u \in A} \int_{\Omega}|\Delta u|^{p} d x, \quad n=1,2, \ldots \tag{1.7}
\end{equation*}
$$

is the sequence of eigenvalues, where

$$
\begin{equation*}
\Gamma_{n}=\left\{A \subseteq W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega) \backslash\{0\}: A=-A, \gamma(A) \geqslant n\right\} \tag{1.8}
\end{equation*}
$$

being $\gamma(A)$ the Krasnoselski's genus of the set $A$. This follows by the LjusternikSchnirelman theory for $C^{1}$-manifolds proved in [13] applied to the functional

$$
\begin{gather*}
\left.J\right|_{\mathcal{M}}(u)=\int_{\Omega}|\Delta u|^{p} d x  \tag{1.9}\\
\mathcal{M}=\left\{u \in W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega): \int_{\Omega}|u|^{p} d x=1\right\},
\end{gather*}
$$

since $\mathcal{M}$ is a $C^{1}$-manifold with tangent space

$$
\begin{equation*}
T_{u} \mathcal{M}=\left\{w \in W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega): \int_{\Omega}|u|^{p-2} u w d x=0\right\} . \tag{1.10}
\end{equation*}
$$

The next remark is the starting point of our paper.
Remark 1.1. It has been recently proved by Drábek and Ôtani [4] that (1.6) with boundary data (1.2) has the least eigenvalue

$$
\begin{equation*}
\lambda_{1}(p)=\inf \left\{\int_{\Omega}|\Delta u|^{p} d x: u \in W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega),\|u\|_{p}^{p}=1\right\} \tag{1.11}
\end{equation*}
$$

which is simple, positive, and isolated in the sense that the solutions of (1.6) with $\lambda=\lambda_{1}(p)$ form a one-dimensional linear space spanned by a positive eigenfunction $\phi_{1}(p)$ associated with $\lambda_{1}(p)$ and there exists $\delta>0$ so that $\left(\lambda_{1}(p), \lambda_{1}(p)+\delta\right)$ does not contain other eigenvalues. The situation is actually more involved with Dirichlet boundary conditions

$$
\begin{equation*}
u=\nabla u=0 \quad \text { on } \partial \Omega \tag{1.12}
\end{equation*}
$$

and, to our knowledge, it is not clear whether the first eigenspace has the previous good properties; the fact is that while Navier boundary conditions allow to reduce the fourth-order problem into a system of two second-order problems, Dirichlet boundary conditions do not. Some pathologies are indeed known, for instance, the first eigenfunction of $\Delta^{2} u=\lambda u$ with boundary data (1.12) may change sign [12].

Remark 1.2. Let $V=\operatorname{span}\left\{\phi_{1}\right\}$ be the eigenspace associated with $\lambda_{1}$, where $\left\|\phi_{1}\right\|_{2, p}=1$. Taking a subspace $W \subset W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega)$ complementing $V$, that is, $W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega)=V \oplus W$, there exists $\hat{\lambda}>\lambda_{1}$ with

$$
\begin{equation*}
\int_{\Omega}|\Delta u|^{p} d x \geqslant \hat{\lambda} \int_{\Omega}|u|^{p} d x \tag{1.13}
\end{equation*}
$$

for each $u \in W$ (in case $p=2$, one may take $\hat{\lambda}=\lambda_{2}$ ).
We may now assume the following conditions:
$\left(\mathscr{H}_{1}\right)$ there exist $R>0$ and $\left.\bar{\lambda} \in\right] \lambda_{1}, \hat{\lambda}[$ such that

$$
\begin{equation*}
|s| \leqslant R \Longrightarrow \lambda_{1}|s|^{p} \leqslant p G(x, s) \leqslant \bar{\lambda}|s|^{p} \tag{1.14}
\end{equation*}
$$

for a.e. $x \in \Omega$ and each $s \in \mathbb{R}$;
$\left(\mathscr{H}_{2}\right)$ there exist $\vartheta>p$ and $M>0$ such that

$$
\begin{equation*}
|s| \geqslant M \Longrightarrow 0<\vartheta G(x, s) \leqslant \operatorname{sg}(x, s) \tag{1.15}
\end{equation*}
$$

for a.e. $x \in \Omega$ and each $s \in \mathbb{R}$.
Assumption $\left(\mathscr{H}_{1}\right)$ corresponds to a resonance condition around the origin while $\left(\mathscr{H}_{2}\right)$ is the standard condition of Ambrosetti-Rabinowitz type.

Theorem 1.3. Assume that conditions ( $\mathscr{H}_{1}$ ) and ( $\mathscr{H}_{2}$ ) hold. Then problem (1.1) with boundary conditions (1.2) admits a nontrivial solution in $W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega)$.

Now replace $\left(\mathscr{H}_{2}\right)$ with a nonresonance condition at $+\infty$.
Theorem 1.4. Assume that condition $\left(\mathscr{H}_{1}\right)$ holds and that for a.e. $x \in \Omega$

$$
\begin{equation*}
\lim _{|s| \rightarrow+\infty} \frac{p G(x, s)}{|s|^{p}}<\lambda_{1} . \tag{1.16}
\end{equation*}
$$

Then problem (1.1) with boundary conditions (1.2) admits two nontrivial solutions in $W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega)$.

We use variational methods to prove Theorems 1.3 and 1.4. Usually, one uses a minimax type argument of mountain pass type to prove the existence of solutions of equations with a variational structure. However, it seems difficult to use minimax theorems in our situation. Thus we will adopt an approach based on Morse theory. Notice that there were a few works using Morse theory to treat $p$-Laplacian problems with Dirichlet boundary conditions (see [9] and the references therein). Moreover, to the authors' knowledge, (1.1) has a very poor literature; the only papers in which a $p$-harmonic equation is mentioned are [1, Section 8] and [4].

The existence of multiple solutions depends mainly on the behaviour of $G(x, s)$ near 0 and at $+\infty$. Without the above resonant or nonresonant conditions to obtain multiple solutions seems hard even in the semilinear case $p=2$.

Remark 1.5. Arguing as in [9], it is possible to prove Theorem 1.4 by replacing assumption (1.16) with the following conditions:

$$
\begin{equation*}
\lim _{|s| \rightarrow+\infty} \frac{p G(x, s)}{|s|^{p}}=\lambda_{1}, \quad \lim _{|s| \rightarrow+\infty}\{g(x, s) s-p G(x, s)\}=+\infty \tag{1.17}
\end{equation*}
$$

for a.e. $x \in \Omega$ (resonance condition at $+\infty$ ).
Remark 1.6. The existence of solutions $u \in W_{0}^{2, p}(\Omega)$ of the quasilinear problem

$$
\begin{align*}
\Delta\left(|\Delta u|^{p-2} \Delta u\right) & =g(x, u) & & \text { in } \Omega,  \tag{1.18}\\
u & =\nabla u=0 & & \text { on } \partial \Omega
\end{align*}
$$

under the previous assumptions $\left(\mathscr{H}_{j}\right)$ is, to our knowledge, an open problem.

## 2. Proofs of Theorems 1.3 and 1.4

In this section, we give the proof of our main results. It is readily seen that

$$
\begin{equation*}
\|u\|_{2, p}=\left(\int_{\Omega}|\Delta u|^{p} d x\right)^{1 / p} \tag{2.1}
\end{equation*}
$$

is an equivalent norm of the standard space norm of $W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega)$. For $\Phi$ a continuously Fréchet differentiable map, let $\Phi^{\prime}$ denote its Fréchet derivative.
Lemma 2.1. The functional $\Phi$ satisfies the Palais-Smale condition.
Proof. Let $\left(u_{h}\right) \subset W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega)$ be such that $\left|\Phi\left(u_{h}\right)\right| \leqslant B$, for some $B>0$ and $\Phi^{\prime}\left(u_{h}\right) \rightarrow 0$ as $h \rightarrow+\infty$. Let $d=\sup _{h \geqslant 0} \Phi\left(u_{h}\right)$. Then we have

$$
\begin{align*}
\vartheta d+\left\|u_{h}\right\|_{2, p} \geqslant & \vartheta \Phi\left(u_{h}\right)-\left\langle\Phi^{\prime}\left(u_{h}\right), u_{h}\right\rangle \\
= & \left(\frac{\vartheta}{p}-1\right)\left\|u_{h}\right\|_{2, p}^{p}-\int_{\left\{\left|u_{h}\right| \geqslant M\right\}}\left[\vartheta G\left(x, u_{h}\right)-g\left(x, u_{h}\right) u_{h}\right] d x \\
& -\int_{\left\{\left|u_{h}\right| \leqslant M\right\}}\left[\vartheta G\left(x, u_{h}\right)-g\left(x, u_{h}\right) u_{h}\right] d x  \tag{2.2}\\
\geqslant & \left(\frac{\vartheta}{p}-1\right)\left\|u_{h}\right\|_{2, p}^{p}-\int_{\left\{\left|u_{h}\right| \leqslant M\right\}}\left[\vartheta G\left(x, u_{h}\right)-g\left(x, u_{h}\right) u_{h}\right] d x \\
\geqslant & \left(\frac{\vartheta}{p}-1\right)\left\|u_{h}\right\|_{2, p}^{p}-D,
\end{align*}
$$

for some $D \in \mathbb{R}$. Thus ( $u_{h}$ ) is bounded and, up to a subsequence, we may assume that $u_{h} \rightharpoonup u$ is, for some $u$, in $W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega)$. Since the embedding $W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega) \hookrightarrow L^{q}(\Omega)$ is compact, then a standard argument shows that $u_{h} \rightarrow u$ strongly and the proof is complete.

Now recall the notion of "Local Linking," which was initially introduced by Liu and $\mathrm{Li}[8]$ and has been used in a vast amount of literature (cf. [2, 5, 6, 11]).

Definition 2.2. Let $E$ be a real Banach space such that $E=V \oplus W$, where $V$ and $W$ are closed subspaces of $E$. Let $\Phi: E \rightarrow \mathbb{R}$ be a $C^{1}$-functional. We say that $\Phi$ has a local linking near the origin 0 (with respect to the decomposition $E=V \oplus W$ ), if there exists $\varphi>0$ such that

$$
\begin{align*}
& u \in V:\|u\| \leqslant \varrho \Longrightarrow \Phi(u) \leqslant 0 \\
& u \in W: 0<\|u\| \leqslant \varrho \Longrightarrow \Phi(u)>0 \tag{2.3}
\end{align*}
$$

We now show that our functional $\Phi$ has a local linking near the origin with respect to the space decomposition $W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega)=V \oplus W$, according to Remark 1.2.

Lemma 2.3. There exists $\varphi>0$ such that conditions (2.3) hold with respect to the decomposition $W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega)=V \oplus W$.

Proof. For $u \in V$, the condition $\|u\|_{2, p} \leqslant \rho$ implies $u(x) \leqslant R$ for a.e. $x \in \Omega$ if $\varrho>0$ is small enough, being $R>0$ as in assumption $\left(\mathscr{H}_{1}\right)$. Thus for $u \in V$,

$$
\begin{align*}
\Phi(u) & =\frac{1}{p} \int_{\Omega}|\Delta u|^{p} d x-\int_{\Omega} G(x, u) d x \\
& =\frac{\lambda_{1}}{p} \int_{\Omega}|u|^{p} d x-\int_{\Omega} G(x, u) d x=\int_{\{|u| \leqslant R\}}\left[\frac{\lambda_{1}}{p}|u|^{p}-G(x, u)\right] d x \leqslant 0 \tag{2.4}
\end{align*}
$$

provided that $\|u\|_{2, p} \leqslant \rho$ and $\rho$ is small.
To prove the second assertion, take $u \in W$. In view of (1.3) and (1.13) we have

$$
\begin{align*}
\Phi(u)= & \frac{1}{p} \int_{\Omega}|\Delta u|^{p} d x-\int_{\Omega} G(x, u) d x \\
= & \frac{1}{p} \int_{\Omega}\left(|\Delta u|^{p}-\bar{\lambda}|u|^{p}\right) d x \\
& -\left(\int_{\{|u| \leqslant R\}}+\int_{\{|u| \geqslant R\}}\right)\left(G(x, u)-\frac{\bar{\lambda}}{p}|u|^{p}\right) d x  \tag{2.5}\\
\geqslant & \frac{1}{p}\left(1-\frac{\bar{\lambda}}{\hat{\lambda}}\right)\|u\|_{2, p}^{p}-c \int_{\Omega}|u|^{s} d x \geqslant \frac{1}{p}\left(1-\frac{\bar{\lambda}}{\hat{\lambda}}\right)\|u\|_{2, p}^{p}-C\|u\|_{2, p}^{s},
\end{align*}
$$

where $p<s \leqslant p_{*}$ and $c, C$ are positive constants. Since $s>p$, it follows that $\Phi(u)>0$ for $\varphi>0$ sufficiently small.

Assume that $u$ is an isolated critical point of $\Phi$ such that $\Phi(u)=c$. We define the critical group of $\Phi$ at $u$ by setting for each $q \in \mathbb{Z}$

$$
\begin{equation*}
C_{q}(\Phi, u)=H_{q}\left(\Phi_{c}, \Phi_{c} \backslash\{u\}\right), \tag{2.6}
\end{equation*}
$$

being $H_{q}(X, Y)$ the $q$ th homology group of the topological pair $(X, Y)$ over the ring $\mathbb{Z}$ and $\Phi_{c}$ the $c$-sublevel of $\Phi$. For the detail of Morse theory and critical groups, we refer the reader to [3].

Since $\operatorname{dim} V=1<+\infty$, by combining Lemma 2.3 and [7, Theorem 2.1], we obtain the following result.

Lemma 2.4. The point 0 is a critical point of $\Phi$ and $C_{1}(\Phi, 0) \neq\{0\}$.
We now investigate the behavior of $\Phi$ near infinity.
Lemma 2.5. There exists a constant $A>0$ such that

$$
\begin{equation*}
a<-A \Longrightarrow \Phi_{a} \simeq S^{\infty} \tag{2.7}
\end{equation*}
$$

where $S^{\infty}=\left\{u \in W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega):\|u\|_{2, p}=1\right\}$.
Proof. By integrating inequality (1.15), we obtain a constant $C_{1}>0$ with

$$
\begin{equation*}
|s| \geqslant M \Longrightarrow G(x, s) \geqslant C_{1}|s|^{9} \tag{2.8}
\end{equation*}
$$

a.e. in $\Omega$ and for each $s \in \mathbb{R}$. Thus, for $u \in S^{\infty}$, we have $\Phi(t u) \rightarrow-\infty$, as $t$ goes to $+\infty$. Set

$$
\begin{equation*}
A=\left(1+\frac{1}{p}\right) M \mathscr{L}^{n}(\Omega) \max _{\Omega \times[-M, M]}|g(x, s)|+1 \tag{2.9}
\end{equation*}
$$

being $\mathscr{L}^{n}$ the Lebesgue measure. As in the proof of [10, Lemma 2.4] we obtain

$$
\begin{align*}
\int_{\Omega} G(x, u) & d x-\frac{1}{p} \int_{\Omega} g(x, u) u d x \\
& \leqslant\left(\frac{1}{9}-\frac{1}{p}\right) \int_{\{|u| \geqslant M\}} g(x, u) u d x+A-1 \tag{2.10}
\end{align*}
$$

For $a<-A$ and

$$
\begin{equation*}
\Phi(t u)=\frac{|t|^{p}}{p}-\int_{\Omega} G(x, t u) d x \leqslant a \quad\left(u \in S^{\infty}\right) \tag{2.11}
\end{equation*}
$$

in view of (2.8) and (2.10), arguing as in the proof of [10, Lemma 2.4],

$$
\begin{equation*}
\frac{d}{d t} \Phi(t u)<0 \tag{2.12}
\end{equation*}
$$

By the implicit function theorem, there is a unique $T \in C\left(S^{\infty}, \mathbb{R}\right)$ such that

$$
\begin{equation*}
\forall u \in S^{\infty}, \quad \Phi(T(u) u)=a . \tag{2.13}
\end{equation*}
$$

For $u \neq 0$, set $\tilde{T}(u)=\left(1 /\|u\|_{2, p}\right) T\left(u /\|u\|_{2, p}\right)$. Then $\tilde{T} \in C\left(W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega) \backslash\{0\}\right.$, $\mathbb{R}$ ) and

$$
\begin{equation*}
\forall u \in W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega) \backslash\{0\}, \quad \Phi(\tilde{T}(u) u)=a . \tag{2.14}
\end{equation*}
$$

We define now a functional $\hat{T}: W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega) \backslash\{0\} \rightarrow \mathbb{R}$ by setting

$$
\hat{T}(u)= \begin{cases}\tilde{T}(u) & \text { if } \Phi(u) \geqslant a  \tag{2.15}\\ 1 & \text { if } \Phi(u) \leqslant a\end{cases}
$$

Since $\Phi(u)=a$ implies $\tilde{T}(u)=1$, we conclude that

$$
\begin{equation*}
\hat{T} \in C\left(W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega) \backslash\{0\}, \mathbb{R}\right) . \tag{2.16}
\end{equation*}
$$

Finally, let $\eta:[0,1] \times W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega) \backslash\{0\} \rightarrow W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega) \backslash\{0\}$,

$$
\begin{equation*}
\eta(s, u)=(1-s) u+s \hat{T}(u) u \tag{2.17}
\end{equation*}
$$

It results that $\eta$ is a strong deformation retract from $W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega) \backslash\{0\}$ to $\Phi_{a}$. Thus $\Phi_{a} \simeq W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega) \backslash\{0\} \simeq S^{\infty}$.

Remark 2.6. A result similar to Lemma 2.5 has been proved for the Laplacian $-\Delta$ in [3, 14], under the additional conditions

$$
\begin{equation*}
g \in C^{1}(\Omega \times \mathbb{R}, \mathbb{R}), \quad g_{t}(x, 0)=\left.\frac{\partial g(x, t)}{\partial t}\right|_{t=0}=0 . \tag{2.18}
\end{equation*}
$$

We recall the following topological result due to Perera [11].
Lemma 2.7. Let $Y \subset B \subset A \subset X$ be topological spaces and $q \in \mathbb{Z}$. If

$$
\begin{equation*}
H_{q}(A, B) \neq\{0\}, \quad H_{q}(X, Y)=\{0\} \tag{2.19}
\end{equation*}
$$

then it results that

$$
\begin{equation*}
H_{q+1}(X, A) \neq\{0\} \quad \text { or } \quad H_{q-1}(B, Y) \neq\{0\} . \tag{2.20}
\end{equation*}
$$

Proof of Theorem 1.3. By Lemma 2.1, $\Phi$ satisfies the Palais-Smale condition. Note that $\Phi(0)=0$, by [3, Chapter I, Theorem 4.2], there exists $\varepsilon>0$ with

$$
\begin{equation*}
H_{1}\left(\Phi_{\varepsilon}, \Phi_{-\varepsilon}\right)=C_{1}(\Phi, 0) \neq\{0\} . \tag{2.21}
\end{equation*}
$$

If $A$ is as in Lemma 2.5, for $a<-A$ we have $\Phi_{a} \simeq S^{\infty}$, which yields

$$
\begin{equation*}
H_{1}\left(W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega), \Phi_{a}\right)=H_{1}\left(W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega), S^{\infty}\right)=\{0\} \tag{2.22}
\end{equation*}
$$

Therefore, being $\Phi_{a} \subset \Phi_{-\varepsilon} \subset \Phi_{\varepsilon}$, Lemma 2.7 yields

$$
\begin{equation*}
H_{2}\left(W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega), \Phi_{\varepsilon}\right) \neq\{0\} \quad \text { or } \quad H_{0}\left(\Phi_{-\varepsilon}, \Phi_{a}\right) \neq\{0\} \tag{2.23}
\end{equation*}
$$

It follows that $\Phi$ has a critical point $u$ for which

$$
\begin{equation*}
\Phi(u)>\varepsilon \quad \text { or } \quad-\varepsilon>\Phi(u)>a \text {. } \tag{2.24}
\end{equation*}
$$

Therefore, $u \neq 0$ and (1.1), (1.2) possess a nontrivial solution.

Recall from [9] the following three-critical point theorem.
Lemma 2.8. Let $X$ be a real Banach space and let $\Phi \in C^{1}(X, \mathbb{R})$ be bounded from below and satisfying the Palais-Smale condition. Assume that $\Phi$ has a critical point $u$ which is homologically nontrivial, that is, $C_{j}(\Phi, u) \neq\{0\}$ for some $j$, and it is not a minimizer for $\Phi$. Then $\Phi$ admits at least three critical points.

Proof of Theorem 1.4. By Lemma 2.8, taking into account Lemma 2.4, it suffices to show that $\Phi$ is bounded from below. Indeed, by (1.16) there exist $\varepsilon>0$ small and $C>0$ such that

$$
\begin{equation*}
G(x, s) \leqslant \frac{\lambda_{1}-\varepsilon}{p}|s|^{p}+C \tag{2.25}
\end{equation*}
$$

for a.e. $x \in \Omega$ and each $s \in \mathbb{R}$. This, by (1.11), immediately yields

$$
\begin{align*}
\Phi(u) & \geqslant \frac{1}{p}\|u\|_{2, p}^{p}-\frac{1}{p}\left(\lambda_{1}-\varepsilon\right)\|u\|_{p}^{p}-C \mathscr{L}^{n}(\Omega) \\
& \geqslant \frac{1}{p}\left(1-\frac{\lambda_{1}-\varepsilon}{\lambda_{1}}\right)\|u\|_{2, p}^{p}-C \mathscr{L}^{n}(\Omega) \longrightarrow+\infty \tag{2.26}
\end{align*}
$$

as $\|u\|_{2, p} \rightarrow+\infty$. Then $\Phi$ is coercive and satisfies the Palais-Smale condition. In particular Lemma 2.8 provides the existence of at least two nontrivial critical points of $\Phi$.

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