ONE-SIDED RESONANCE FOR QUASILINEAR PROBLEMS WITH ASYMMETRIC NONLINEARITIES

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1. Introduction

We consider the quasilinear elliptic boundary value problem,

$$-\Delta_p u = \alpha_+(x) \left(u^+\right)^{p-1} - \alpha_-(x) \left(u^-\right)^{p-1} + f(x, u), \quad u \in W_0^{1, p}(\Omega), \tag{1.1}$$

where Ω is a bounded domain in \mathbb{R}^n , $n \ge 1$, $\Delta_p u = \operatorname{div}(|\nabla u|^{p-2}\nabla u)$ is the *p*-Laplacian, $1 , <math>u^{\pm} = \max\{\pm u, 0\}$, $\alpha_{\pm} \in L^{\infty}(\Omega)$, and *f* is a Carathéodory function on $\Omega \times \mathbb{R}$ satisfying a growth condition,

$$|f(x,t)| \le qV(x)^{p-q} |t|^{q-1} + W(x)^{p-1}, \tag{1.2}$$

with $1 \le q < p$ and $V, W \in L^p(\Omega)$. We assume that (1.1) is resonant from one side in the sense that either

$$\lambda_l \le \alpha_{\pm}(x) \le \lambda_{l+1} - \varepsilon \tag{1.3}$$

or

$$\lambda_l + \varepsilon \le \alpha_{\pm}(x) \le \lambda_{l+1}, \tag{1.4}$$

for two consecutive variational eigenvalues, $\lambda_l < \lambda_{l+1}$ of $-\Delta_p$ on $W_0^{1,p}(\Omega)$, and some $\varepsilon > 0$ (see Section 2 for the definition of the variational spectrum).

The special case where $\alpha_+(x) = \alpha_-(x) \equiv \lambda_l$ and q = 1 was recently studied by Arcoya and Orsina [1], Bouchala and Drábek [3], and Drábek and Robinson [8] (see also Cuesta et al. [6] and Dancer and Perera [7]). In the present paper, we prove a single existence theorem for the general case that includes all their results and much more.

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54 One-sided resonance

Denote by N the set of nontrivial solutions of the asymptotic problem

$$-\Delta_p u = \alpha_+(x) \left(u^+ \right)^{p-1} - \alpha_-(x) \left(u^- \right)^{p-1}, \quad u \in W_0^{1,p}(\Omega), \tag{1.5}$$

and set

$$F(x,t) := \int_0^t f(x,s) \, ds, \qquad H(x,t) := pF(x,t) - tf(x,t). \tag{1.6}$$

Our main result is the following theorem.

THEOREM 1.1. Problem (1.1) has a solution in the following cases:

- (i) equation (1.3) holds and $\int_{\Omega} H(x, u_j) \rightarrow +\infty$,
- (ii) equation (1.4) holds and ∫_Ω H(x, u_j) → -∞ for every sequence (u_j) in W₀^{1,p}(Ω) such that ||u_j|| → ∞ and u_j/||u_j|| converges to some element of N. In particular, (1.1) is solvable when (1.3) or (1.4) holds and N is empty.

As is usually the case in resonance problems, the main difficulty here is the lack of compactness of the associated variational functional, which we will overcome by constructing a sequence of approximating nonresonance problems, finding approximate solutions for them using linking and min-max type arguments, and passing to the limit (see Rabinowitz [10] for standard details of the variational theory). But first we give some corollaries and deduce the results of [1, 3, 8]. In what follows, (u_j) is as in the theorem, that is, $\rho_j := ||u_j|| \to \infty$ and $v_j := u_j/\rho_j \to v \in N$.

First, we give simple pointwise assumptions on H that imply the limits in the theorem.

COROLLARY 1.2. Problem (1.1) has a solution in the following cases:

- (i) equation (1.3) holds, $H(x, t) \rightarrow +\infty$ a.e. as $|t| \rightarrow \infty$, and $H(x, t) \ge -C(x)$,
- (ii) equation (1.4) holds, $H(x,t) \rightarrow -\infty$ a.e. as $|t| \rightarrow \infty$, and $H(x,t) \leq C(x)$ for some $C \in L^1(\Omega)$.

Note that this corollary makes no reference to N.

Proof. If (i) holds, then $H(x, u_j(x)) = H(x, \rho_j v_j(x)) \rightarrow +\infty$ for a.e. *x* such that $v(x) \neq 0$ and $H(x, u_j(x)) \ge -C(x)$, so

$$\int_{\Omega} H(x, u_j) \ge \int_{\nu \neq 0} H(x, u_j) - \int_{\nu = 0} C(x) \longrightarrow +\infty$$
(1.7)

by Fatou's lemma. Similarly, $\int_{\Omega} H(x, u_j) \rightarrow -\infty$ if (ii) holds.

Note that the above argument goes through as long as the limits in (i) and (ii) hold on subsets of $\{x \in \Omega : v(x) \neq 0\}$ with positive measure. Now, taking $w = v^{\pm}$ in

$$\int_{\Omega} |\nabla v|^{p-2} \nabla v \cdot \nabla w = \int_{\Omega} \left[\alpha_+(x) \left(v^+ \right)^{p-1} - \alpha_-(x) \left(v^- \right)^{p-1} \right] w \tag{1.8}$$

gives

$$\|v^{\pm}\|^{p} = \int_{\Omega_{\pm}} \alpha_{\pm}(x) (v^{\pm})^{p} \leq \|\alpha_{\pm}\|_{\infty} \|v^{\pm}\|_{p^{*}}^{p} \mu(\Omega_{\pm})^{p/n}$$

$$\leq \|\alpha_{\pm}\|_{\infty} S^{-1} \|v^{\pm}\|^{p} \mu(\Omega_{\pm})^{p/n},$$
(1.9)

where $\Omega_{\pm} = \{x \in \Omega : v(x) \ge 0\}$, $p^* = np/(n-p)$ is the critical Sobolev exponent, *S* is the best constant for the embedding $W_0^{1,p}(\Omega) \hookrightarrow L^{p^*}(\Omega)$, and μ is the Lebesgue measure in \mathbb{R}^n . So

$$\mu(\Omega_{\pm}) \ge \left(S \| \alpha_{\pm} \|_{\infty}^{-1} \right)^{n/p}, \tag{1.10}$$

and hence

$$\mu(\{x \in \Omega : \nu(x) = 0\}) \le \mu(\Omega) - S^{n/p} \Big(\|\alpha_+\|_{\infty}^{-n/p} + \|\alpha_-\|_{\infty}^{-n/p} \Big).$$
(1.11)

Thus, we have the following corollary.

COROLLARY 1.3. Problem (1.1) has a solution in the following cases:

- (i) equation (1.3) holds, $H(x,t) \to +\infty$ in Ω' as $|t| \to \infty$, and $H(x,t) \ge -C(x)$,
- (ii) equation (1.4) holds, $H(x,t) \to -\infty$ in Ω' as $|t| \to \infty$, and $H(x,t) \le C(x)$ for some $\Omega' \subset \Omega$ with $\mu(\Omega') > \mu(\Omega) - S^{n/p}(\|\alpha_+\|_{\infty}^{-n/p} + \|\alpha_-\|_{\infty}^{-n/p})$ and $C \in L^1(\Omega)$.

Similar conditions on *H* were recently used by Furtado and Silva [9] in the semilinear case p = 2.

Next, note that

$$\underline{H}_{+}(x)(\nu^{+}(x))^{q} + \underline{H}_{-}(x)(\nu^{-}(x))^{q} \\
\leq \liminf \frac{H(x, u_{j}(x))}{\rho_{j}^{q}} \leq \limsup \frac{H(x, u_{j}(x))}{\rho_{j}^{q}} \\
\leq \overline{H}_{+}(x)(\nu^{+}(x))^{q} + \overline{H}_{-}(x)(\nu^{-}(x))^{q},$$
(1.12)

where

$$\underline{H}_{\pm}(x) = \liminf_{t \to \pm \infty} \frac{H(x,t)}{|t|^{q}}, \qquad \overline{H}_{\pm}(x) = \limsup_{t \to \pm \infty} \frac{H(x,t)}{|t|^{q}}.$$
 (1.13)

Moreover,

$$\frac{\left|H(x,u_j(x))\right|}{\rho_j^q} \le (p+q)V(x)^{p-q} \left|v_j(x)\right|^q + \frac{(p+1)W(x)^{p-1} \left|v_j(x)\right|}{\rho_j^{q-1}} \tag{1.14}$$

56 One-sided resonance

by (1.2), so it follows that

$$\int_{\Omega} \underline{H}_{+}(v^{+})^{q} + \underline{H}_{-}(v^{-})^{q} \leq \liminf \frac{\int_{\Omega} H(x, u_{j})}{\rho_{j}^{q}}$$

$$\leq \limsup \frac{\int_{\Omega} H(x, u_{j})}{\rho_{j}^{q}} \leq \int_{\Omega} \overline{H}_{+}(v^{+})^{q} + \overline{H}_{-}(v^{-})^{q}.$$
(1.15)

Thus we have the following corollary.

COROLLARY 1.4. Problem (1.1) has a solution in the following cases:

- (i) equation (1.3) holds and $\int_{\Omega} \underline{H}_+(v^+)^q + \underline{H}_-(v^-)^q > 0$ for all $v \in N$,
- (ii) equation (1.4) holds and $\int_{\Omega} \overline{H}_+(v^+)^q + \overline{H}_-(v^-)^q < 0$ for all $v \in N$.

When $\alpha_+(x) = \alpha_-(x) \equiv \lambda_1$ and q = 1 this reduces to the result of Bouchala and Drábek [3].

Finally, we note that if

$$\frac{tf(x,t)}{|t|^q} \longrightarrow f_{\pm}(x) \quad \text{a.e. as } t \longrightarrow \pm \infty, \tag{1.16}$$

then

$$\frac{F(x,t)}{|t|^{q}} = \frac{1}{|t|^{q}} \int_{0}^{t} \left[\frac{sf(x,s)}{|s|^{q}} - f_{\pm}(x) \right] |s|^{q-2} s \, ds + \frac{f_{\pm}(x)}{q} \longrightarrow \frac{f_{\pm}(x)}{q} \tag{1.17}$$

and hence

$$\frac{H(x,t)}{|t|^q} \longrightarrow \left(\frac{p}{q} - 1\right) f_{\pm}(x), \tag{1.18}$$

so Corollary 1.4 implies the following corollary.

COROLLARY 1.5. Problem (1.1) has a solution in the following cases:

- (i) equation (1.3) holds and $\int_{\Omega} f_+(v^+)^q + f_-(v^-)^q > 0$ for all $v \in N$,
- (ii) equation (1.4) holds and $\int_{\Omega} f_{+}(v^{+})^{q} + f_{-}(v^{-})^{q} < 0$ for all $v \in N$.

This was proved in Arcoya and Orsina [1] and Drábek and Robinson [8] for the special case $\alpha_+(x) = \alpha_-(x) \equiv \lambda_l$ and q = 1.

2. Proof of Theorem 1.1

First we recall some facts about the variational spectrum of the *p*-Laplacian. It is easily seen from the Lagrange multiplier rule that the eigenvalues of $-\Delta_p$ on $W_0^{1,p}(\Omega)$ correspond to the critical values of

$$J(u) = \int_{\Omega} |\nabla u|^p, \quad u \in S := \left\{ u \in W_0^{1,p}(\Omega) : ||u||_p = 1 \right\}.$$
 (2.1)

Moreover, *J* satisfies the Palais-Smale condition (cf. Drábek and Robinson [8]). Thus, we can define an unbounded sequence of min-max eigenvalues by

$$\lambda_{l} := \inf_{A \in \mathscr{F}_{l}} \max_{u \in A} J(u), \quad l \in \mathbb{N},$$
(2.2)

where

$$\mathcal{F}_l = \left\{ A \subset S : \exists \text{ a continuous odd surjection } h : S^{l-1} \longrightarrow A \right\}$$
(2.3)

and S^{l-1} is the unit sphere in \mathbb{R}^l .

LEMMA 2.1. λ_l is an eigenvalue of $-\Delta_p$ and $\lambda_l \rightarrow \infty$.

Proof. If λ_l is a regular value of J, then there is an $\varepsilon > 0$ and an odd homeomorphism $\eta : S \to S$ such that $\eta(J^{\lambda_l+\varepsilon}) \subset J^{\lambda_l-\varepsilon}$ by [2, Theorem 2.5] (the standard first deformation lemma is not sufficient because the manifold S is not of class $C^{1,1}$ when p < 2). But then taking $A \in \mathcal{F}_l$ with max $J(A) \leq \lambda_l + \varepsilon$ and setting $\tilde{A} = \eta(A)$, we get a set in \mathcal{F}_l for which max $J(\tilde{A}) \leq \lambda_l - \varepsilon$, contradicting the definition of λ_l . Finally, denoting by $\mu_l \to \infty$ the usual Ljusternik-Schnirelmann eigenvalues, we have $\lambda_l \geq \mu_l$ since the genus of each A in \mathcal{F}_l is l, so $\lambda_l \to \infty$.

It is not known whether this is a complete list of eigenvalues when $p \neq 2$ and $n \geq 2$. However, the variational structure provided by this portion of the spectrum is sufficient to show that the associated functional admits a linking geometry in the nonresonant case. We only consider (i) as the proof for (ii) is similar. Let

$$\alpha_{\pm}^{j}(x) = \begin{cases} \alpha_{\pm}(x), & \text{if } \alpha_{\pm}(x) \ge \lambda_{l} + \frac{1}{j}, \\ \lambda_{l} + \frac{1}{j}, & \text{if } \alpha_{\pm}(x) < \lambda_{l} + \frac{1}{j}, \end{cases}$$
(2.4)

so that

$$\lambda_l + \frac{1}{j} \le \alpha_{\pm}^j(x) \le \lambda_{l+1} - \varepsilon, \qquad \left| \alpha_{\pm}^j(x) - \alpha_{\pm}(x) \right| \le \frac{1}{j}, \tag{2.5}$$

and let

$$\Phi_{j}(u) = \int_{\Omega} |\nabla u|^{p} - \alpha_{+}^{j}(x) (u^{+})^{p} - \alpha_{-}^{j}(x) (u^{-})^{p} - pF(x, u), \quad u \in W_{0}^{1, p}(\Omega).$$
(2.6)

First, we show that there is a $u_j \in W_0^{1,p}(\Omega)$ such that

$$\|u_j\|\|\Phi'_j(u_j)\| \longrightarrow 0, \quad \inf \Phi_j(u_j) > -\infty.$$
(2.7)

By (2.2), there is an $A \in \mathcal{F}_l$ such that

$$J(u) \le \lambda_l + \frac{1}{2j}, \quad u \in A.$$
(2.8)

58 One-sided resonance

For $u \in A$ and R > 0,

$$\Phi_{j}(Ru) = R^{p} \left[J(u) - \int_{\Omega} \alpha_{+}^{j}(x) (u^{+})^{p} + \alpha_{-}^{j}(x) (u^{-})^{p} \right] - \int_{\Omega} pF(x, Ru)$$

$$\leq -\frac{R^{p}}{2j} + p \left(\|V\|_{p}^{p-q} R^{q} + \|W\|_{p}^{p-1} R \right)$$
(2.9)

by (1.2), (2.5), and (2.8), so

$$\max_{u \in A} \Phi_j(Ru) \longrightarrow -\infty \quad \text{as } R \longrightarrow \infty.$$
 (2.10)

Next, let

$$\mathcal{S} = \left\{ u \in W_0^{1,p}(\Omega) : \int_{\Omega} |\nabla u|^p \ge \lambda_{l+1} \int_{\Omega} |u|^p \right\}.$$
 (2.11)

For $u \in \mathcal{G}$,

$$\Phi_{j}(u) \ge \varepsilon \|u\|_{p}^{p} - p\Big(\|V\|_{p}^{p-q}\|u\|_{p}^{q} + \|W\|_{p}^{p-1}\|u\|_{p}\Big),$$
(2.12)

so

$$\inf_{u \in \mathcal{G}} \Phi_j(u) \ge C := \min_{r \ge 0} \left[\varepsilon r^p - p \left(\|V\|_p^{p-q} r^q + \|W\|_p^{p-1} r \right) \right] > -\infty.$$
(2.13)

Now use (2.10) to fix R > 0 so large that

$$\max \Phi_j(RA) < C, \tag{2.14}$$

where $RA = \{Ru : u \in A\}$.

Since $A \in \mathcal{F}_l$, there is a continuous odd surjection $h : S^{l-1} \to A$. Let

$$\Gamma = \left\{ \varphi \in C\left(D^l, W_0^{1, p}(\Omega)\right) : \varphi|_{S^{l-1}} = Rh \right\},$$
(2.15)

where D^l is the unit disk in \mathbb{R}^l with boundary S^{l-1} . We claim that *RA* links \mathcal{G} with respect to Γ , that is,

$$\varphi(D^l) \cap \mathcal{G} \neq \emptyset \quad \forall \varphi \in \Gamma.$$
(2.16)

To see this, first note that the proof is done if $0 \in \varphi(D^l)$. Otherwise, denoting by π the radial projection onto $S, \tilde{A} := \pi(\varphi(D^l)) \cup -\pi(\varphi(D^l)) \in \mathcal{F}_{l+1}$, and hence

$$\max_{u\in\pi(\varphi(D^l))} J(u) = \max_{u\in\bar{A}} J(u) \ge \lambda_{l+1},$$
(2.17)

so $\pi(\varphi(D^l)) \cap \mathcal{G} \neq \emptyset$, which implies that $\varphi(D^l) \cap \mathcal{G} \neq \emptyset$.

Now it follows from a deformation argument of Cerami [5] that there is a u_j such that

$$\|u_j\| \|\Phi'_j(u_j)\| \longrightarrow 0, \qquad |\Phi_j(u_j) - c_j| \longrightarrow 0, \qquad (2.18)$$

where

$$c_j := \inf_{\varphi \in \Gamma} \max_{u \in \varphi(D^l)} \Phi_j(u) \ge C, \tag{2.19}$$

from which (2.7) follows.

We complete the proof by showing that a subsequence of (u_j) converges to a solution of (1.1). It is easy to see that this is the case if (u_j) is bounded, so suppose that $\rho_j := ||u_j|| \to \infty$. Setting $v_j := u_j/\rho_j$ and passing to a subsequence, we may assume that $v_j \to v$ weakly in $W_0^{1,p}(\Omega)$, strongly in $L^p(\Omega)$, and a.e. in Ω . Then

$$\int_{\Omega} |\nabla v_{j}|^{p-2} \nabla v_{j} \cdot \nabla (v_{j} - v)$$

$$= \frac{(\Phi'_{j}(u_{j}), v_{j} - v)}{p\rho_{j}^{p-1}} + \int_{\Omega} \left[\alpha_{+}^{j}(x) (v_{j}^{+})^{p-1} - \alpha_{-}^{j}(x) (v_{j}^{-})^{p-1} + \frac{f(x, u_{j})}{\rho_{j}^{p-1}} \right] (v_{j} - v) \longrightarrow 0,$$
(2.20)

and we deduce that $v_j \to v$ strongly in $W_0^{1,p}(\Omega)$ (cf. Browder [4]). In particular, ||v|| = 1, so $v \neq 0$. Moreover, for each $w \in W_0^{1,p}(\Omega)$, passing to the limit in

$$\frac{(\Phi'_{j}(u_{j}), w)}{p\rho_{j}^{p-1}} = \int_{\Omega} |\nabla v_{j}|^{p-2} \nabla v_{j} \cdot \nabla w$$
$$- \left[\alpha_{+}^{j}(x) (v_{j}^{+})^{p-1} - \alpha_{-}^{j}(x) (v_{j}^{-})^{p-1} + \frac{f(x, u_{j})}{\rho_{j}^{p-1}} \right] w$$
(2.21)

gives that

$$\int_{\Omega} |\nabla v|^{p-2} \nabla v \cdot \nabla w - \left[\alpha_{+}(x) \left(v^{+} \right)^{p-1} - \alpha_{-}(x) \left(v^{-} \right)^{p-1} \right] w = 0,$$
(2.22)

so $v \in N$. Thus,

$$\frac{(\Phi'_j(u_j), u_j)}{p} - \Phi_j(u_j) = \int_{\Omega} H(x, u_j) \longrightarrow +\infty, \qquad (2.23)$$

contradicting (2.7).

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- 60 One-sided resonance
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