## Research Article

# Oscillation Properties for Second-Order Partial Differential Equations with Damping and Functional Arguments 

Run Xu, Yuhua Lu, and Fanwei Meng<br>Department of Mathematics, Qufu Normal University, Shandong, Qufu 273165, China

Correspondence should be addressed to Run Xu, xurun_2005@163.com
Received 19 September 2011; Accepted 21 October 2011
Academic Editor: Norio Yoshida
Copyright © 2011 Run Xu et al. This is an open access article distributed under the Creative Commons Attribution License, which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited.

Using an integral averaging method and generalized Riccati technique, by introducing a parameter $\beta \geq 1$, we derive new oscillation criteria for second-order partial differential equations with damping. The results are of high degree of generality and sharper than most known ones.

## 1. Introduction

Consider the second-order partial delay differential equation

$$
\begin{align*}
\frac{\partial}{\partial t}\left(r(t) \frac{\partial}{\partial t} u(x, t)\right)+p(t) \frac{\partial u(x, t)}{\partial t}= & a(t) \Delta u(x, t)+\sum_{k=1}^{s} a_{k}(t) \Delta u\left(x, t-\rho_{k}(t)\right)-q(x, t) f(u(x, t)) \\
& -\sum_{j=1}^{m} q_{j}(x, t) f_{j}\left(u\left(x, t-\sigma_{j}\right)\right), \quad(x, t) \in \Omega \times R_{+} \equiv G \tag{1.1}
\end{align*}
$$

where $\Delta$ is the Laplacian in $R^{N}, R_{+}=[0, \infty)$ and $\Omega$ is a bounded domain in $R^{N}$ with a piecewise smooth boundary $\partial \Omega$.

Throughout this paper, we assume that
$\left(\mathrm{H}_{1}\right) r(t) \in C^{1}\left(R_{+},(0, \infty)\right), p(t) \in C\left(R_{+}, R\right) ;$
$\left(\mathrm{H}_{2}\right) q(x, t), q_{j}(x, t) \in C\left(\overline{\mathrm{G}}, R_{+}\right), q(t)=\min _{x \in \overline{\mathrm{G}}} q(x, t), q_{j}(t)=\min _{x \in \overline{\mathrm{G}}} q_{j}(x, t), j \in I_{m}=\{1$, $2, \ldots, m\} ;$
$\left(\mathrm{H}_{3}\right) a(t), a_{k}(t), \rho_{k}(t) \in C\left(R_{+}, R_{+}\right), \lim _{t \rightarrow \infty}\left(t-\rho_{k}(t)\right)=\infty$, and $\sigma_{j}$ are nonnegative constants, $j \in I_{m}, k \in I_{s}=\{1,2, \ldots, s\} ;$
$\left(\mathrm{H}_{4}\right) f(u) \in C^{1}(R, R), f_{j}(u) \in C(R, R)$ are convex in $R_{+}$with $u f_{j}(u)>0, u f(u)>0$, and $f^{\prime}(u) \geq \mu>0,(u \neq 0)$.
We say that a continuous function $H(t, s)$ belongs to the function class $\omega$, denoted by $H \in \omega$, if $H \in C\left(D, R_{+}\right)$, where $D=\{(t, s):-\infty<s \leq t<+\infty\}$, satisfy

$$
\begin{equation*}
H(t, t)=0, \quad H(t, s)>0, \quad-\infty<s<t<+\infty \tag{1.2}
\end{equation*}
$$

Furthermore, the continuous partial derivative $\partial H / \partial S$ exists on $D$, and there is $h \in$ $L_{\text {loc }}(D, R)$, such that

$$
\begin{equation*}
\frac{\partial H}{\partial s}=-h(t, s) \sqrt{H(t, s)} \tag{1.3}
\end{equation*}
$$

Various results on the oscillation for the partial functional differential equation have been obtained recently. We refer the reader to [1-3] for parabolic equations and to [4-11] for hyperbolic equations.

Recently, Li and Cui [12] studied the equation of the form

$$
\begin{align*}
\frac{\partial}{\partial t}\left[p(t) \frac{\partial}{\partial t}\left(u(x, t)+\sum_{i=1}^{l} \lambda_{i}(t) u\left(x, t-\tau_{i}\right)\right)\right]= & a(t) \Delta u(x, t)+\sum_{k=1}^{s} a_{k}(t) \Delta u\left(x, t-\rho_{k}(t)\right)-q(t) u(x, t) \\
& -\sum_{j=1}^{m} q_{j}(x, t) u\left(x, t-\sigma_{j}\right), \quad(x, t) \in \Omega \times R_{+} \equiv G \tag{1.4}
\end{align*}
$$

with Robin boundary condition

$$
\begin{equation*}
\frac{\partial u(x, t)}{\partial \gamma}+g(x, t) u(x, t)=0, \quad(x, t) \in \partial \Omega \times R_{+} \tag{1.5}
\end{equation*}
$$

where $\gamma$ is the unit exterior vector to $\partial \Omega$ and $g(x, t)$ is a nonnegative continuous function on $\partial \Omega \times R_{+}$and obtained the following result.

Theorem A (see [12, Theorem 2.2]). Suppose that $H \in \omega$, let
$\left(\mathrm{C}_{1}\right) 0<\inf _{s \geq t_{0}}\left\{\liminf _{t \rightarrow \infty}\left(H(t, s) / H\left(t, t_{0}\right)\right)\right\} \leq \infty$, suppose that there exists some $j_{0} \in I_{m}$ and there exist two functions $\phi \in C^{1}\left[t_{0}, \infty\right), A \in C\left[t_{0}, \infty\right)$ satisfying,
$\left(\mathrm{C}_{2}\right) \lim \sup _{t \rightarrow \infty}\left(1 / H\left(t, t_{0}\right)\right) \int_{t_{0}}^{t} p\left(s-\sigma_{j_{0}}\right) \phi(s) h^{2}(t, s) d s<\infty$,
$\left(C_{3}\right) \int_{t_{0}}^{\infty}\left(A_{+}^{2}(s) / p\left(s-\sigma_{j_{0}}\right) \phi(s)\right) d s=\infty$, and for every $t_{1} \geq t_{0}$,
$\left(\mathrm{C}_{4}\right) \lim \sup _{t \rightarrow \infty}\left(1 / H\left(t, t_{1}\right)\right) \int_{t_{1}}^{t}\left[H(t, s) \psi(s)-(1 / 4) \phi(s) p\left(s-\sigma_{j_{0}}\right) h^{2}(t, s)\right] d s \geq A\left(t_{1}\right)$,
where $\phi(s)=\exp \left\{-2 \int^{s} \phi(\xi) d \xi\right\}, A_{+}(s)=\max \{A(s), 0\}$, and $\psi(s)=\phi(s)\left\{\alpha_{j_{0}} q_{j_{0}}(s)\left[1-\sum_{i=1}^{l} \lambda_{i}(s-\right.\right.$ $\left.\left.\left.\sigma_{j_{0}}\right)\right]+p\left(s-\sigma_{j_{0}}\right) \phi^{2}(s)-\left[p\left(s-\sigma_{j_{0}}\right) \phi(s)\right]^{\prime}\right\}$. Then every solution $u(x, t)$ of the problem (1.4), (1.5) is oscillatory in $G$.

In 2008, Rogovchenko and Tuncay [13] established new oscillation criteria for secondorder nonlinear differential equations with damping term

$$
\begin{equation*}
\left(r(t) x^{\prime}(t)\right)^{\prime}+p(t) x^{\prime}(t)+q(t) f(x(t))=0 \tag{1.6}
\end{equation*}
$$

without an assumption that has been required in related results reported in the literature over the last two decades. Motivated by the ideas in $[12,13]$, by introducing a Parameter $\beta \geq 1$, we will further improve Theorems A and derive new interval criteria for oscillation of (1.1). We suggest two different approaches which allow one to remove condition $\left(\mathrm{C}_{2}\right)$ in Theorem A. A modified integral averaging technique enables one to simplify essentially the proofs of oscillation criteria.

## 2. Main Results

Theorem 2.1. Suppose that there exists a function $y \in C^{1}\left[t_{0}, \infty\right)$ such that for some $\beta \geq 1$ and for some $H \in \omega$,

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \sup \frac{1}{H\left(t, t_{0}\right)} \int_{t_{0}}^{t}\left(H(t, s) \psi(s)-\frac{\beta}{4 \mu} v(s) r(s) h^{2}(t, s)\right) d s=\infty, \quad t \geq 0 \tag{2.1}
\end{equation*}
$$

where

$$
\begin{gather*}
v(t)=\exp \left(-2 \int^{t}\left(\mu y(s)-\frac{p(s)}{2 r(s)}\right) d s\right),  \tag{2.2}\\
\psi(t)=v(t)\left[q(t)+\mu r(t) y^{2}(t)-p(t) y(t)-(r(t) y(t))^{\prime}\right], \tag{2.3}
\end{gather*}
$$

then every solution $u(x, t)$ of the problem (1.1), (1.5) is oscillatory in $G$.
Proof. Suppose to the contrary that there is a nonoscillatory solution $u(x, t)$ of the problem (1.1), (1.5) which has no zero on $\Omega \times\left[t_{0}, \infty\right)$ for some $t_{0}>0$. Without loss of generality, we assume that $u(x, t)>0, u\left(x, t-\rho_{k}(t)\right)>0$ and $u\left(x, t-\sigma_{j}\right)>0$ in $\Omega \times\left[t_{1}, \infty\right), t_{1} \geq t_{0}, k \in I_{s}, j \in I_{m}$. Integrating (1.1) with respect to $x$ over the domain $\Omega$, we have

$$
\begin{align*}
& \frac{d}{d t}\left(r(t) \frac{d}{d t} \int_{\Omega} u(x, t) d x\right)+p(t) \frac{d}{d t} \int_{\Omega} u(x, t) d x \\
& \quad=a(t) \int_{\Omega} \Delta u(x, t) d x+\sum_{k=1}^{s} a_{k}(t) \int_{\Omega} \Delta u\left(x, t-\rho_{k}(t)\right) d x  \tag{2.4}\\
& \quad-\int_{\Omega} q(x, t) f(u(x, t)) d x-\sum_{j=1}^{m} \int_{\Omega} q_{j}(x, t) f_{j}\left(u\left(x, t-\sigma_{j}\right)\right) d x, \quad t \geq t_{1} .
\end{align*}
$$

From Green's formula and the boundary condition (1.5), we have

$$
\begin{align*}
\int_{\Omega} \Delta u(x, t) d x & =\int_{\partial \Omega} \frac{\partial u(x, t)}{\partial \gamma} d s=-\int_{\partial \Omega} g(x, t) u(x, t) d s \leq 0 \\
\int_{\Omega} \Delta u\left(x, t-\rho_{k}(t)\right) d x & =\int_{\partial \Omega} \frac{\partial u\left(x, t-\rho_{k}(t)\right)}{\partial \gamma} d s  \tag{2.5}\\
& =-\int_{\partial \Omega} g\left(x, t-\rho_{k}(t)\right) u\left(x, t-\rho_{k}(t)\right) d s \leq 0, \quad t \geq t_{1}, k \in I_{s},
\end{align*}
$$

where $d s$ denotes the surface element on $\partial \Omega$. Moreover, from $\left(\mathrm{H}_{2}\right),\left(\mathrm{H}_{4}\right)$ and Jensen's inequality, we have

$$
\begin{align*}
\int_{\Omega} q(x, t) f(u(x, t)) d x & \geq q(t) \int_{\Omega} f(u(x, t)) d x \geq|\Omega| q(t) f\left(\frac{1}{|\Omega|} \int_{\Omega} u(x, t) d x\right) \\
\int_{\Omega} q_{j}(x, t) f_{j}\left(u\left(x, t-\sigma_{j}\right)\right) d x & \geq q_{j}(t) \int_{\Omega} f_{j}\left(u\left(x, t-\sigma_{j}\right)\right) d x  \tag{2.6}\\
& \geq|\Omega| q_{j}(t) f_{j}\left(\frac{1}{|\Omega|} \int_{\Omega} u\left(x, t-\sigma_{j}\right)\right) d x
\end{align*}
$$

where $|\Omega|=\int_{\Omega} d x$.
Set

$$
\begin{equation*}
U(t)=\frac{1}{|\Omega|} \int_{\Omega} u(x, t) d x, \quad t \geq t_{1} \tag{2.7}
\end{equation*}
$$

In view of (2.5)-(2.7), (2.4) yields that

$$
\begin{equation*}
\left(r(t) U^{\prime}(t)\right)^{\prime}+p(t) U^{\prime}(t)+q(t) f(U(t))+\sum_{j=1}^{m} q_{j}(t) f_{j}\left(U\left(t-\sigma_{j}\right)\right) \leq 0, \quad t \geq t_{1} \tag{2.8}
\end{equation*}
$$

Note that $\left(\mathrm{H}_{4}\right),(2.8)$ yields that

$$
\begin{equation*}
\left(r(t) U^{\prime}(t)\right)^{\prime}+p(t) U^{\prime}(t)+q(t) f(U(t)) \leq 0, \quad t \geq t_{1} . \tag{2.9}
\end{equation*}
$$

Put

$$
\begin{equation*}
w(t)=v(t) r(t)\left[\frac{U^{\prime}(t)}{f(U(t))}+y(t)\right], \quad t \geq t_{1} \tag{2.10}
\end{equation*}
$$

where $v(t)$ is given by (2.2), then

$$
\begin{align*}
& w^{\prime}(t)= {\left[-2 \mu y(t)+\frac{p(t)}{r(t)}\right] w(t)+v(t)\left[\frac{\left(r(t) U^{\prime}(t)\right)^{\prime}}{f(U(t))}-\frac{r(t) f^{\prime}(U(t))\left(U^{\prime}(t)\right)^{2}}{f^{2}(U(t))}+(r(t) y(t))^{\prime}\right] } \\
& \begin{aligned}
\leq & {\left[-2 \mu y(t)+\frac{p(t)}{r(t)}\right] w(t)+v(t)\left[\frac{\left(r(t) U^{\prime}(t)\right)^{\prime}}{f(U(t))}-\frac{\mu r(t)\left(U^{\prime}(t)\right)^{2}}{f^{2}(U(t))}+(r(t) y(t))^{\prime}\right] } \\
\leq & {\left[-2 \mu y(t)+\frac{p(t)}{r(t)}\right] w(t)-v(t)\left[\frac{p(t) U^{\prime}(t)}{f(U(t))}+q(t)+\mu r(t)\left(\frac{U^{\prime}(t)}{f(U(t))}\right)^{2}-(r(t) y(t))^{\prime}\right] } \\
= & {\left[-2 \mu y(t)+\frac{p(t)}{r(t)}\right] w(t)-v(t)\left[p(t)\left(\frac{w(t)}{v(t) r(t)}-y(t)\right)+q(t)\right.} \\
= & {\left.\left[-2 \mu y(t)+\frac{p(t)}{r(t)}\right] w(t) \quad+\mu r(t)\left[\frac{w(t)}{v(t) r(t)}-y(t)\right]^{2}-(r(t) y(t))^{\prime}\right] } \\
& -v(t)\left[\frac{p(t)}{v(t) r(t)} w(t)-p(t) y(t)+q(t)\right. \\
& \left.\quad+\mu r(t) \frac{w^{2}(t)}{v^{2}(t) r^{2}(t)}-2 \mu \frac{w(t) y(t)}{v(t)}+\mu r(t) y^{2}(t)-(r(t) y(t))^{\prime}\right] \\
= & -2 \mu y(t) w(t)+\frac{p(t)}{r(t)} w(t)-\frac{p(t)}{r(t)} w(t)-v(t)\left[-p(t) y(t)+q(t)+\mu r(t) y^{2}(t)-(r(t) y(t))^{\prime}\right] \\
& +2 \mu y(t) w(t)-\mu \frac{w^{2}(t)}{v(t) r(t)} \\
= & -v(t)\left[-p(t) y(t)+q(t)+\mu r(t) y^{2}(t)-(r(t) y(t))^{\prime}\right]-\mu \frac{w^{2}(t)}{v(t) r(t)} \\
= & -\psi(t)-\mu \frac{w^{2}(t)}{v(t) r(t)},
\end{aligned}
\end{align*}
$$

that is,

$$
\begin{equation*}
\psi(t) \leq-w^{\prime}(t)-\mu \frac{w^{2}(t)}{v(t) r(t)} \tag{2.12}
\end{equation*}
$$

where $\psi(t)$ is defined by (2.3). Multiplying (2.12) by $H(t, s)$ and integrating from $T$ to $t$, we have, for some $\beta \geq 1$ and for all $t \geq T \geq t_{1}$,

$$
\begin{align*}
\int_{T}^{t} H(t, s) \psi(s) d s \leq & -\int_{T}^{t} H(t, s) w^{\prime}(s) d s-\int_{T}^{t} H(t, s) \frac{\mu}{v(s) r(s)} w^{2}(s) d s \\
= & H(t, T) w(T)-\int_{T}^{t}\left(h(t, s) \sqrt{H(t, s)} w(s)+H(t, s) \frac{\mu w^{2}(s)}{v(s) r(s)}\right) d s \\
= & H(t, T) w(T)-\int_{T}^{t}\left[\sqrt{\frac{\mu H(t, s)}{\beta v(s) r(s)}} w(s)+\sqrt{\frac{\beta v(s) r(s)}{4 \mu}} h(t, s)\right]^{2} d s  \tag{2.13}\\
& +\frac{\beta}{4 \mu} \int_{T}^{t} v(s) r(s) h^{2}(t, s) d s-\int_{T}^{t} \frac{(\beta-1) \mu}{\beta v(s) r(s)} H(t, s) w^{2}(s) d s
\end{align*}
$$

Writing the latter inequality in the form

$$
\begin{align*}
& \int_{T}^{t}\left[H(t, s) \psi(s)-\frac{\beta}{4 \mu} v(s) r(s) h^{2}(t, s)\right] d s \\
& \quad \leq H(t, T) w(T)-\int_{T}^{t}\left[\sqrt{\frac{\mu H(t, s)}{\beta v(s) r(s)}} w(s)+\sqrt{\frac{\beta v(s) r(s)}{4 \mu}} h(t, s)\right]^{2} d s  \tag{2.14}\\
& \quad-\int_{T}^{t} \frac{(\beta-1) \mu}{\beta v(s) r(s)} H(t, s) w^{2}(s) d s .
\end{align*}
$$

Using the properties of $H(t, s)$, we have

$$
\begin{equation*}
\int_{t_{1}}^{t}\left[H(t, s) \psi(s)-\frac{\beta}{4 \mu} v(s) r(s) h^{2}(t, s)\right] d s \leq H\left(t, t_{1}\right)\left|w\left(t_{1}\right)\right| \leq H\left(t, t_{0}\right)\left|w\left(t_{1}\right)\right|, \quad t \geq t_{1} \tag{2.15}
\end{equation*}
$$

and for all $t \geq t_{1} \geq t_{0}$,

$$
\begin{equation*}
\int_{t_{0}}^{t}\left[H(t, s) \psi(s)-\frac{\beta}{4 \mu} v(s) r(s) h^{2}(t, s)\right] d s \leq H\left(t, t_{0}\right)\left[\int_{t_{0}}^{t_{1}}|\psi(s)| d s+\left|w\left(t_{1}\right)\right|\right] \tag{2.16}
\end{equation*}
$$

By (2.16),

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \sup \frac{1}{H\left(t, t_{0}\right)} \int_{t_{0}}^{t}\left(H(t, s) \psi(s)-\frac{\beta}{4 \mu} v(s) r(s) h^{2}(t, s)\right) d s \leq \int_{t_{0}}^{t_{1}}|\psi(s)| d s+\left|w\left(t_{1}\right)\right|<\infty \tag{2.17}
\end{equation*}
$$

which contradicts (2.1). This proves Theorem 2.1.

Consider a Kamenev-type function $H(t, s)$ defined by $H(t, s)=(t-s)^{n-1},(t, s) \in D$, where $n>2$ is an integer. Obviously, $H$ belongs to the class $\omega$, and $h(t, s)=(n-1)(t-$ $s)^{(n-3) / 2},(t, s) \in D$. Then, we can get the following results.

Corollary 2.2. Suppose that there exists a function $y(t) \in C^{1}\left(\left[t_{0}, \infty\right) ; R\right)$ such that for some integer $n>2$ and some $\beta \geq 1$,

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \sup \frac{1}{t^{n-1}} \int_{t_{0}}^{t}(t-s)^{n-3}\left[(t-s)^{2} \psi(s)-\frac{\beta(n-1)^{2}}{4 \mu} v(s) r(s)\right] d s=\infty \tag{2.18}
\end{equation*}
$$

where $v(t)$ and $\psi(t)$ are as defined in Theorem 2.1. Then every solution $u(x, t)$ of the problem (1.1), (1.5) is oscillatory in $G$.

Theorem 2.3. Suppose that

$$
\begin{equation*}
0<\inf _{s \geq t_{0}}\left(\lim _{t \rightarrow \infty} \inf \frac{H(t, s)}{H\left(t, t_{0}\right)}\right) \leq \infty \tag{2.19}
\end{equation*}
$$

Assume that there exist functions $f \in C^{1}\left(\left[t_{0}, \infty\right) ; R\right)$ and $\phi \in C\left(\left[t_{0}, \infty\right) ; R\right)$ such that, for all $t \geq T \geq$ $t_{0}$ and for some $\beta>1$,

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \sup \frac{1}{H(t, T)} \int_{T}^{t}\left(H(t, s) \psi(s)-\frac{\beta}{4 \mu} v(s) r(s) h^{2}(t, s)\right) d s \geq \phi(T) \tag{2.20}
\end{equation*}
$$

where $v(t), \psi(t)$ are as defined in Theorem 2.1 and suppose further that

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \sup \int_{t_{0}}^{t} \frac{\phi_{+}^{2}(s)}{v(s) r(s)}=\infty \tag{2.21}
\end{equation*}
$$

where $\phi_{+}(t)=\max (\phi(t), 0)$. Then every solution $u(x, t)$ of the problem (1.1), (1.5) is oscillatory in $G$.

Proof. Suppose to the contrary that there is a nonoscillatory solution $u(x, t)$ of the problem (1.1), (1.5) which has no zero on $\Omega \times\left[t_{1}, \infty\right)$ for some $t_{1}>t_{0}$, without loss of generality, we assume that $u(x, t)>0, u\left(x, t-\rho_{k}(t)\right)>0$ and $u\left(x, t-\sigma_{j}\right)>0$ in $\Omega \times\left[t_{1}, \infty\right), t \geq t_{1} \geq t_{0}, k \in$ $I_{s}, j \in I_{m}$.

As in the proof of Theorem 2.1, (2.14) holds for all $t \geq T \geq t_{1}$, we have

$$
\begin{align*}
& \frac{1}{H(t, T)} \int_{T}^{t}\left[H(t, s) \psi(s)-\frac{\beta}{4 \mu} v(s) r(s) h^{2}(t, s)\right] d s \\
& \quad \leq w(T)-\frac{1}{H(t, T)} \int_{T}^{t}\left[\sqrt{\frac{\mu H(t, s)}{\beta v(s) r(s)}} w(s)+\sqrt{\frac{\beta v(s) r(s)}{4 \mu}} h(t, s)\right]^{2} d s  \tag{2.22}\\
& \quad-\frac{1}{H(t, T)} \int_{T}^{t} \frac{(\beta-1) \mu}{\beta v(s) r(s)} H(t, s) w^{2}(s) d s \\
& \quad \leq w(T)-\frac{1}{H(t, T)} \int_{T}^{t} \frac{(\beta-1) \mu}{\beta v(s) r(s)} H(t, s) w^{2}(s) d s
\end{align*}
$$

Therefore, for $t>T \geq T_{0}$,

$$
\begin{align*}
& \lim _{t \rightarrow \infty} \sup \frac{1}{H(t, T)} \int_{T}^{t}\left[H(t, s) \psi(s)-\frac{\beta}{4 \mu} v(s) r(s) h^{2}(t, s)\right] d s \\
& \quad \leq w(T)-\lim _{t \rightarrow \infty} \inf \frac{1}{H(t, T)} \int_{T}^{t} \frac{(\beta-1) \mu}{\beta v(s) r(s)} H(t, s) w^{2}(s) d s . \tag{2.23}
\end{align*}
$$

It follows from (2.20) that

$$
\begin{equation*}
w(T) \geq \phi(T)+\lim _{t \rightarrow \infty} \inf \frac{1}{H(t, T)} \int_{T}^{t} \frac{(\beta-1) \mu}{\beta v(s) r(s)} H(t, s) w^{2}(s) d s \tag{2.24}
\end{equation*}
$$

for all $T \geq t_{1}$ and for any $\beta>1$. Then, for all $T \geq t_{1}$,

$$
\begin{gather*}
w(T) \geq \phi(T)  \tag{2.25}\\
\lim _{t \rightarrow \infty} \inf \frac{1}{H\left(t, t_{1}\right)} \int_{t_{1}}^{t} \frac{H(t, s)}{v(s) r(s)} w^{2}(s) d s \leq \frac{\beta}{\mu(\beta-1)}\left(w\left(t_{1}\right)-\phi\left(t_{1}\right)\right) \tag{2.26}
\end{gather*}
$$

Now, we claim that

$$
\begin{equation*}
\int_{t_{1}}^{\infty} \frac{w^{2}(s)}{v(s) r(s)} d s<\infty \tag{2.27}
\end{equation*}
$$

Suppose the contrary, that is,

$$
\begin{equation*}
\int_{t_{1}}^{\infty} \frac{w^{2}(s)}{v(s) r(s)} d s=\infty \tag{2.28}
\end{equation*}
$$

By (2.19), there is a positive constant $M_{1}$, satisfying

$$
\begin{equation*}
\inf _{s \geq t_{0}}\left(\lim _{t \rightarrow \infty} \inf \frac{H(t, s)}{H\left(t, t_{0}\right)}\right)>M_{1}>0 \tag{2.29}
\end{equation*}
$$

Let $M$ be any arbitrary positive number, then from (2.28) we get that there exists a $T_{1}>t_{1}$ such that, for all $t \geq T_{1}$,

$$
\begin{equation*}
\int_{t_{1}}^{t} \frac{w^{2}(s)}{v(s) r(s)} d s \geq \frac{M}{M_{1}} \tag{2.30}
\end{equation*}
$$

Using integration by parts, for all $t \geq T_{1}$, we get

$$
\begin{align*}
\frac{1}{H\left(t, t_{1}\right)} \int_{t_{1}}^{t} H(t, s) \frac{w^{2}(s)}{v(s) r(s)} d s & =\frac{1}{H\left(t, t_{1}\right)} \int_{t_{1}}^{t}\left(-\frac{\partial H(t, s)}{\partial s}\right)\left(\int_{t_{1}}^{s} \frac{w^{2}(\tau)}{v(\tau) r(\tau)} d \tau\right) d s \\
& \geq \frac{1}{H\left(t, t_{1}\right)} \int_{T_{1}}^{t}\left(-\frac{\partial H(t, s)}{\partial s}\right)\left(\int_{t_{1}}^{s} \frac{w^{2}(\tau)}{v(\tau) r(\tau)} d \tau\right) d s  \tag{2.31}\\
& \geq \frac{M}{M_{1}} \frac{1}{H\left(t, t_{1}\right)} \int_{T_{1}}^{t}\left(-\frac{\partial H(t, s)}{\partial s}\right) d s \\
& =\frac{M}{M_{1}} \frac{H\left(t, T_{1}\right)}{H\left(t, t_{1}\right)}
\end{align*}
$$

By (2.29), there exists a $T_{2}>T_{1}$ such that, for all $t \geq T_{2}$,

$$
\begin{equation*}
\frac{H\left(t, T_{1}\right)}{H\left(t, t_{1}\right)} \geq M_{1} \tag{2.32}
\end{equation*}
$$

It follows from (2.31) that for all $t \geq T_{2}$,

$$
\begin{equation*}
\frac{1}{H\left(t, t_{1}\right)} \int_{t_{1}}^{t} H(t, s) \frac{w^{2}(s)}{v(s) r(s)} d s \geq M \tag{2.33}
\end{equation*}
$$

Since $M$ is an arbitrary positive constant,

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \inf \frac{1}{H\left(t, t_{1}\right)} \int_{t_{1}}^{t} H(t, s) \frac{w^{2}(s)}{v(s) r(s)} d s=\infty \tag{2.34}
\end{equation*}
$$

which contradicts (2.26). Consequently, (2.27) holds. And from (2.25), we obtain

$$
\begin{equation*}
\int_{t_{1}}^{\infty} \frac{\phi_{+}^{2}(s)}{v(s) r(s)} d s \leq \int_{t_{1}}^{\infty} \frac{w^{2}(s)}{v(s) r(s)} d s<\infty \tag{2.35}
\end{equation*}
$$

which contradicts (2.21). This completes the proof of Theorem 2.3.

Choosing H as in Corollary 2.2, by Theorem 2.3, we can obtain the following corollary.
Corollary 2.4. Let $v(t)$ and $\psi(t)$ be as in Theorem 2.1, assume further that there exist functions $f \in C^{1}\left(\left[t_{0}, \infty\right) ; R\right)$ and $\phi \in C\left(\left[t_{0}, \infty\right) ; R\right)$ such that, for all $T \geq t_{0}$, for some integer $n>2$, and for some $\beta>1$,

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \sup \frac{1}{t^{n-1}} \int_{T}^{t}(t-s)^{n-3}\left[(t-s)^{2} \psi(s)-\frac{\beta(n-1)^{2}}{4 \mu} v(s) r(s)\right] d s \geq \phi(T) \tag{2.36}
\end{equation*}
$$

and (2.21) hold. Then every solution $u(x, t)$ of the problem (1.1), (1.5) is oscillatory in $G$.
Theorem 2.5. Suppose that there exists a function $f \in C^{1}\left(\left[t_{0}, \infty\right) ; R\right)$ such that for some $\beta \geq 1$ and for some $H \in \omega$,

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \sup \frac{1}{H\left(t, t_{0}\right)} \int_{t_{0}}^{t}\left[H(t, s)\left(\bar{\psi}(s)-\frac{p^{2}(s) \bar{v}(s)}{2 \mu r(s)}\right)-\frac{\mu \beta}{2} \bar{v}(s) r(s) h^{2}(t, s)\right] d s=\infty, \tag{2.37}
\end{equation*}
$$

where

$$
\begin{gather*}
\bar{v}(t)=\exp \left(-2 \mu \int^{t} y(s) d s\right)  \tag{2.38}\\
\bar{\psi}(t)=\bar{v}(t)\left(q(t)+\mu r(t) y^{2}(t)-p(t) y(t)-(r(t) y(t))^{\prime}\right)
\end{gather*}
$$

Then every solution $u(x, t)$ of the problem (1.1), (1.5) is oscillatory in $G$.
Proof. As in Theorem 2.1, without loss of generality, we assume that a nonoscillatory solution $u(x, t)$ of the problem (1.1), (1.5) satisfies $u(x, t)>0, u\left(x, t-\rho_{k}(t)\right)>0$ and $u\left(x, t-\sigma_{j}\right)>0$ in $\Omega \times\left[t_{1}, \infty\right), t_{1} \geq t_{0}, k \in I_{s}, j \in I_{m}$. Define a generalized Riccati transformation

$$
\begin{equation*}
\bar{w}(t)=\bar{v}(t) r(t)\left[\frac{U^{\prime}(t)}{f(U(t))}+y(t)\right], \quad t \geq t_{1} \tag{2.39}
\end{equation*}
$$

where $\bar{v}(t)$ is given by (2.38). Then

$$
\begin{align*}
\bar{w}^{\prime}(t) & \leq-2 \mu y(t) \bar{w}(t)+\bar{v}(t)\left\{-q(t)+(r(t) y(t))^{\prime}-p(t)\left[\frac{\bar{w}(t)}{\bar{v}(t) r(t)}-y(t)\right]-\mu r(t)\left[\frac{\bar{w}(t)}{\bar{v}(t) r(t)}-y(t)\right]^{2}\right\} \\
& =-\bar{\psi}(t)-\frac{p(t)}{r(t)} \bar{w}(t)-\mu \frac{1}{r(t) \bar{v}(t)} \bar{w}^{2}(t), \quad t \geq t_{1} . \tag{2.40}
\end{align*}
$$

Using an elementary inequality

$$
\begin{equation*}
-a x^{2}+b x \leq-\frac{a}{2} x^{2}+\frac{b^{2}}{2 a} \tag{2.41}
\end{equation*}
$$

for all $a>0$ and for all $b, x \in R$, we conclude from (2.40) that

$$
\begin{equation*}
\bar{\psi}(t)-\frac{p^{2}(t) \bar{v}(t)}{2 \mu r(t)} \leq-\bar{w}^{\prime}(t)-\frac{\mu}{2 \bar{v}(t) r(t)} \bar{w}^{2}(t), \quad t \geq t_{1} . \tag{2.42}
\end{equation*}
$$

Multiplying (2.42) by $H(t, s)$ and integrating from $T<t$, we obtain, for some $\beta \geq 1$ and for all $t \geq T \geq t_{1}$,

$$
\begin{align*}
\int_{T}^{t} H(t, s)\left(\bar{\psi}(s)-\frac{p^{2}(s) \bar{v}(s)}{2 \mu r(s)}\right) d s \leq & H(t, T) \bar{w}(T)-\int_{T}^{t} h(t, s) \sqrt{H(t, s)} \bar{w}(s) d s-\int_{T}^{t} \frac{\mu \bar{w}^{2}(s)}{2 \bar{v}(s) r(s)} d s \\
\leq & H(t, T) \bar{w}(T)+\frac{\mu \beta}{2} \int_{T}^{t} \bar{v}(s) r(s) h^{2}(t, s) d s \\
& -\mu \int_{T}^{t} \frac{(\beta-1) H(t, s)}{2 \beta \bar{v}(s) r(s)} \bar{w}^{2}(s) d s \\
& -\frac{\mu}{2} \int_{T}^{t}\left(\sqrt{\frac{H(t, s)}{\beta \bar{v}(s) r(s)}} \bar{w}(s)+\sqrt{\beta \bar{v}(s) r(s)} h(t, s)\right)^{2} d s . \tag{2.43}
\end{align*}
$$

Therefore, for all $t \geq T \geq t_{1}$, we have

$$
\begin{align*}
& \int_{T}^{t}\left[H(t, s) \bar{\psi}(s)-H(t, s) \frac{p^{2}(s) \bar{v}(s)}{2 \mu r(s)}-\frac{\mu \beta}{2} \bar{v}(s) r(s) h^{2}(t, s)\right] d s \\
& \leq H(t, T) \bar{w}(T)-\mu \int_{T}^{t} \frac{(\beta-1) H(t, s)}{2 \beta \bar{v}(s) r(s)} \bar{w}^{2}(s) d s  \tag{2.44}\\
& \quad-\frac{\mu}{2} \int_{T}^{t}\left(\sqrt{\left.\frac{H(t, s)}{\beta \bar{v}(s) r(s)} \bar{w}(s)+\sqrt{\beta \bar{v}(s) r(s)} h(t, s)\right)^{2} d s} .\right.
\end{align*}
$$

Following the same lines as in the proof of Theorem 2.1, we have

$$
\begin{align*}
& \lim _{t \rightarrow \infty} \sup \frac{1}{H\left(t, t_{0}\right)} \int_{t_{0}}^{t}\left(H(t, s) \bar{\psi}(s)-H(t, s) \frac{p^{2}(s) \bar{v}(s)}{2 \mu r(s)}-\frac{\mu \beta}{2} \bar{v}(s) r(s) h^{2}(t, s)\right) d s  \tag{2.45}\\
& \quad \leq \int_{t_{0}}^{t_{1}}|\bar{\psi}(s)| d s+\left|\bar{w}\left(t_{1}\right)\right|<\infty
\end{align*}
$$

which contradicts the assumption (2.19).
This completes the proof.

Theorem 2.6. Let (2.19) holds. Assume that there exist functions $f \in C^{1}\left(\left[t_{0}, \infty\right) ; R\right)$ and $\phi \in$ $C\left(\left[t_{0}, \infty\right), R\right)$ such that, for all $t \geq t_{0}$, any $T \geq t_{0}$, and for some $\beta>1$,

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \sup \frac{1}{H(t, T)} \int_{T}^{t}\left[H(t, s)\left(\bar{\psi}(s)-\frac{p^{2}(s) \bar{v}(s)}{2 \mu r(s)}\right)-\frac{\mu \beta}{2} \bar{v}(s) r(s) h^{2}(t, s)\right] d s \geq \phi(T) \tag{2.46}
\end{equation*}
$$

and (2.21) holds, where $\psi(t), \bar{v}(t)$ are defined as in Theorem 2.6 and $\phi_{+}(t)=\max (\phi(t), 0)$, then every solution $u(x, t)$ of the problem (1.1), (1.5) is oscillatory in $G$.

Theorem 2.7. Let all assumptions of Theorem 2.6 be satisfied except that condition (2.46) be replaced by

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \inf \frac{1}{H(t, T)} \int_{T}^{t}\left[H(t, s)\left(\bar{\psi}(s)-\frac{p^{2}(s) \bar{v}(s)}{2 \mu r(s)}\right)-\frac{\mu \beta}{2} \bar{v}(s) r(s) h^{2}(t, s)\right] d s \geq \phi(T) \tag{2.47}
\end{equation*}
$$

Then every solution $u(x, t)$ of the problem (1.1), (1.5) is oscillatory in $G$.
Remark 2.8. By introducing the parameter $\beta$ in Theorem 2.3, we derive new oscillation criteria of the problem (1.1), (1.5) which are simpler than that in Theorem A; furthermore, modifications of the proofs through the refinement of the standard integral averaging method allowed us to shorten significantly the proofs of Theorem 2.3. We can also derive a number of oscillation criteria with the appropriate choice of the function $H$ and $\rho$, here, we omit the details.

## 3. Examples

Now, we consider these following examples.
Example 3.1. Consider the partial differential equation

$$
\begin{align*}
\frac{\partial}{\partial t}\left[\frac{1}{t}\right. & \left.\frac{\partial}{\partial t}\left(u(x, t)+\frac{1}{2} u(x, t-\pi)\right)\right]+2 \cos t \frac{\partial u(x, t)}{\partial t} \\
& =\Delta u(x, t)+\frac{1}{t^{2}} \Delta u\left(x, t-\frac{3}{2} \pi\right)  \tag{3.1}\\
& -\left(\frac{2}{t^{3}}+t \cos ^{2} t+\sin t\right) f(u(x, t))-\frac{1}{t^{2}} f_{1}(u(x, t-\pi)), \quad(x, t) \in(0, \pi) \times(0, \infty),
\end{align*}
$$

with the boundary condition

$$
\begin{equation*}
u_{x}(0, t)=u_{x}(\pi, t)=0, \quad t>0, \tag{3.2}
\end{equation*}
$$

where $f(u)=u^{3}+u, f_{1}(u)=u e^{u}+u$.
Here, $N=1, l=1, s=1, m=1, \mu=1, r(t)=1 / t, p(t)=2 \cos t, q(x, t)=q(t)=\left(2 / t^{3}+\right.$ $\left.t \cos ^{2} t+\sin t\right), q_{1}(x, t)=1 / t^{2}, f(u)=f_{j}(u)=u, a(t)=1, a_{1}(t)=1 / t^{2}, \rho_{1}(t)=(3 / 2) \pi, \sigma_{1}=$ $\pi, \tau_{1}=\pi$.

Let

$$
\begin{equation*}
y(t)=\frac{1}{t}+t \cos t \tag{3.3}
\end{equation*}
$$

then

$$
\begin{equation*}
v(t)=t^{2}, \quad \psi(t)=t^{-1} \tag{3.4}
\end{equation*}
$$

Let $n=3$, for any $\beta \geq 1$,

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \sup \frac{1}{t^{2}} \int_{1}^{t}\left[(t-s)^{2} s^{-1}-\beta s^{2} \frac{1}{s}\right] d s=\lim _{t \rightarrow \infty} \sup \frac{1}{t^{2}} \int_{1}^{t}\left[(t-s)^{2} s^{-1}-\beta s\right] d s=\infty \tag{3.5}
\end{equation*}
$$

Therefore, Corollary 2.2 holds, then every solution $u(x, t)$ of the problem (3.1), (3.2) oscillates in $(0, \pi) \times(0, \infty)$.

Example 3.2. Consider the partial differential equation

$$
\begin{align*}
& \frac{\partial}{\partial t}\left[\left(1+\frac{1}{2 t^{3}}\right)(2+\sin t) \frac{\partial u(x, t)}{\partial t}\right]+\frac{3}{t}\left(1+\frac{1}{2 t^{3}}\right)(2+\sin t) \frac{\partial u(x, t)}{\partial t} \\
&= 3 \Delta u(x, t)+(2-\cos t) \Delta u\left(x, t-\frac{3}{2} \pi\right)-t^{-3}\left(\left(1-t^{3}+2 t^{2}-6 t\right) \sin t+12 t\right) f(u(x, t)) \\
&-(2 t+\sin t) f_{1}(u(x, t-\pi))-2 f_{2}\left(u\left(x, t-\frac{\pi}{2}\right)\right), \quad(x, t) \in(0, \pi) \times(0, \infty) \tag{3.6}
\end{align*}
$$

with the boundary condition (3.2), where $f(u)=u^{5}+u, f_{1}(u)=u \sin ^{2} u, f_{2}(u)=u^{3} \cos ^{2} u$.
Here $N=1, s=1, m=2, \mu=1, r(t)=\left(1+1 / 2 t^{3}\right)(2+\sin t), p(t)=(3 / t)\left(1+1 / 2 t^{3}\right)(2+$ $\sin t), q(t)=t^{-3}\left[\left(1-t^{3}+2 t^{2}-6 t\right) \sin t+12 t\right], a(t)=3, a_{1}(t)=2-\cos t, q_{1}(x, t)=2+\sin t$, $q_{2}(x, t)=2, \rho_{1}(t)=(3 / 2) \pi, \sigma_{1}=\pi, \sigma_{2}=\pi / 2$.

Let $y(t)=0$, then $v(t)=t^{3}$ and $\psi(t)=v(t) q(t)=\left(1-t^{3}+2 t^{2}-6 t\right) \sin t+12 t$.
Choose $\beta=2, n=3$, a straightforward computation yields

$$
\begin{gather*}
\limsup _{t \rightarrow \infty} \frac{1}{t^{2}} \int_{T}^{t}\left[(t-s)^{2}\left(\left(1-s^{3}+2 s^{2}-6 s\right) \sin s+12 s\right)-\left(2 s^{3}+1\right)(2+\sin s)\right] d s  \tag{3.7}\\
=16-T^{3} \cos T+T^{2}(2 \cos T-6+3 \sin T)-4 T \sin T-3 \cos T=\phi(T) .
\end{gather*}
$$

Let $\phi_{+}(t)=\max (\phi(t), 0)$. It is not difficult to see that

$$
\begin{equation*}
\limsup _{t \rightarrow \infty} \int_{1}^{t} \frac{\phi_{+}^{2}(s)}{\left(s^{3}+(1 / 2)\right)(2+\sin s)} d s \geq \limsup _{t \rightarrow \infty} \int_{1}^{t} \frac{\phi_{+}^{2}(s)}{3\left(s^{3}+(1 / 2)\right)} d s=\infty \tag{3.8}
\end{equation*}
$$

By Corollary 2.4, we obtain that every solution of problem (3.6), (3.2) oscillates in $(0, \pi) \times$ $(0, \infty)$.

Note that in this example,

$$
\begin{equation*}
\limsup _{t \rightarrow \infty} \frac{1}{t^{2}} \int_{1}^{t} 4\left(s^{3}+\frac{1}{2}+(2+\sin s)\right) d s=\infty \tag{3.9}
\end{equation*}
$$

so the condition $\left(\mathrm{C}_{2}\right)$ would not have been satisfied with the same choices of $v(t)$.

## Acknowledgments

This research was supported by National Science Foundation of China (11171178), the fund of subject for doctor of ministry of education (20103705110003), and the Natural Science Foundations of Shandong Province of China (ZR2009AM011 and ZR2009AL015).

## References

[1] D. P. Mishev and D. D. Bănov, "Oscillation of the solutions of parabolic differential equations of neutral type," Applied Mathematics and Computation, vol. 28, no. 2, pp. 97-111, 1988.
[2] X. L. Fu and W. Zhuang, "Oscillation of certain neutral delay parabolic equations," Journal of Mathematical Analysis and Applications, vol. 191, no. 3, pp. 473-489, 1995.
[3] B. T. Cui, "Oscillation properties for parabolic equations of neutral type," Journal of Computational and Applied Mathematics, vol. 33, no. 4, pp. 581-588, 1992.
[4] B. T. Cui, Y. H. Yu, and S. Z. Lin, "Oscillation of solutions to hyperbolic differential equations with delays," Acta Mathematicae Applicatae Sinica, vol. 19, no. 1, pp. 80-88, 1996 (Chinese).
[5] B. S. Lalli, Y. H. Yu, and B. T. Cui, "Oscillation of hyperbolic equations with functional arguments," Applied Mathematics and Computation, vol. 53, no. 2-3, pp. 97-110, 1993.
[6] W. N. Li, "Oscillation for solutions of partial differential equations with delays," Demonstratio Mathematica, vol. 33, no. 2, pp. 319-332, 2000.
[7] R. P. Agarwal, F. W. Meng, and W. N. Li, "Oscillation of solutions of systems of neutral type partial functional differential equations," Computers and Mathematics with Applications, vol. 44, no. 5-6, pp. 777-786, 2002.
[8] W. N. Li and F. W. Meng, "Forced oscillation for certain systems of hyperbolic differential equations," Applied Mathematics and Computation, vol. 141, no. 2-3, pp. 313-320, 2003.
[9] W. N. Li and F. W. Meng, "Oscillation for systems of neutral partial differential equations with continuous distributed deviating arguments," Demonstratio Mathematica, vol. 34, no. 3, pp. 619-633, 2001.
[10] W. N. Li and B. T. Cui, "Oscillation of solutions of neutral partial functional-differential equations," Journal of Mathematical Analysis and Applications, vol. 234, no. 1, pp. 123-146, 1999.
[11] W. N. Li, "Oscillation properties for systems and hyperbolic differential equations of neutral type," Journal of Mathematical Analysis and Applications, vol. 248, no. 2, pp. 369-384, 2000.
[12] W. N. Li and B. T. Cui, "Oscillation of solutions of neutral partial functional-differential equations," Journal of Mathematical Analysis and Applications, vol. 234, no. 1, pp. 123-146, 1999.
[13] Y. V. Rogovchenko and F. Tuncay, "Oscillation criteria for second-order nonlinear differential equations with damping," Nonlinear Analysis, vol. 69, no. 1, pp. 208-221, 2008.


