# BERTRAND MATE OF SPACELIKE BIHARMONIC CURVES WITH TIMELIKE BINORMAL ACCORDING TO FLAT METRIC IN LORENTZIAN HEISENBERG GROUP 

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Abstract. In this paper, we study spacelike biharmonic curves with timelike binormal according to flat metric in the Lorentzian Heisenberg group Heis ${ }^{3}$. We characterize Bertrand mate of spacelike biharmonic curves with timelike binormal in terms of their curvature and torsion.

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## 1.Introduction

Bertrand curves discovered by J. Bertrand in 1850 are one of the important and interesting topic of classical special curve theory. A Bertrand curve is defined as a special curve which shares its principal normals with another special curve (called Bertrand mate). Note that Bertrand mates are particular examples of offset curves used in computer-aided design.

Harmonic maps $f:(M, g) \longrightarrow(N, h)$ between Riemannian manifolds are the critical points of the energy

$$
\begin{equation*}
E(f)=\frac{1}{2} \int_{M}|d f|^{2} v_{g} \tag{1.1}
\end{equation*}
$$

and they are therefore the solutions of the corresponding Euler-Lagrange equation. This equation is given by the vanishing of the tension field

$$
\begin{equation*}
\tau(f)=\operatorname{trace} \nabla d f . \tag{1.2}
\end{equation*}
$$

As suggested by Eells and Sampson in [5], we can define the bienergy of a map $f$ by

$$
\begin{equation*}
E_{2}(f)=\frac{1}{2} \int_{M}|\tau(f)|^{2} v_{g} \tag{1.3}
\end{equation*}
$$

and say that is biharmonic if it is a critical point of the bienergy.
Jiang derived the first and the second variation formula for the bienergy in [9], showing that the Euler-Lagrange equation associated to $E_{2}$ is

$$
\begin{align*}
\tau_{2}(f) & =-\mathcal{J}^{f}(\tau(f))=-\Delta \tau(f)-\operatorname{trace} R^{N}(d f, \tau(f)) d f  \tag{1.4}\\
& =0
\end{align*}
$$

where $\mathcal{J}^{f}$ is the Jacobi operator of $f$. The equation $\tau_{2}(f)=0$ is called the biharmonic equation. Since $\mathcal{J}^{f}$ is linear, any harmonic map is biharmonic. Therefore, we are interested in proper biharmonic maps, that is non-harmonic biharmonic maps.

In this paper, we study spacelike biharmonic curves with timelike binormal according to flat metric in the Lorentzian Heisenberg group Heis ${ }^{3}$. We characterize Bertrand mate of spacelike biharmonic curves with timelike binormal in terms of their curvature and torsion.

## 2.Preliminaries

The Heisenberg group Heis ${ }^{3}$ is a Lie group which is diffeomorphic to $\mathbb{R}^{3}$ and the group operation is defined as

$$
(x, y, z) *(\bar{x}, \bar{y}, \bar{z})=(x+\bar{x}, y+\bar{y}, z+\bar{z}-\bar{x} y+x \bar{y}) .
$$

The identity of the group is $(0,0,0)$ and the inverse of $(x, y, z)$ is given by $(-x,-y,-z)$. The left-invariant Lorentz metric on Heis ${ }^{3}$ is

$$
g=d x^{2}+(x d y+d z)^{2}-((1-x) d y-d z)^{2} .
$$

The following set of left-invariant vector fields forms an orthonormal basis for the corresponding Lie algebra:

$$
\begin{equation*}
\left\{\mathbf{e}_{1}=\frac{\partial}{\partial x}, \mathbf{e}_{2}=\frac{\partial}{\partial y}+(1-x) \frac{\partial}{\partial z}, \mathbf{e}_{3}=\frac{\partial}{\partial y}-x \frac{\partial}{\partial z}\right\} . \tag{2.1}
\end{equation*}
$$

The characterising properties of this algebra are the following commutation relations:

$$
\left[\mathbf{e}_{2}, \mathbf{e}_{3}\right]=0, \quad\left[\mathbf{e}_{3}, \mathbf{e}_{1}\right]=\mathbf{e}_{2}-\mathbf{e}_{3}, \quad\left[\mathbf{e}_{2}, \mathbf{e}_{1}\right]=\mathbf{e}_{2}-\mathbf{e}_{3},
$$

with

$$
\begin{equation*}
g\left(\mathbf{e}_{1}, \mathbf{e}_{1}\right)=g\left(\mathbf{e}_{2}, \mathbf{e}_{2}\right)=1, \quad g\left(\mathbf{e}_{3}, \mathbf{e}_{3}\right)=-1 . \tag{2.2}
\end{equation*}
$$

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Proposition 2.1. For the covariant derivatives of the Levi-Civita connection of the left-invariant metric $g$, defined above the following is true:

$$
\nabla=\left(\begin{array}{ccc}
0 & 0 & 0  \tag{2.3}\\
\mathbf{e}_{2}-\mathbf{e}_{3} & -\mathbf{e}_{1} & -\mathbf{e}_{1} \\
\mathbf{e}_{2}-\mathbf{e}_{3} & -\mathbf{e}_{1} & -\mathbf{e}_{1}
\end{array}\right)
$$

where the $(i, j)$-element in the table above equals $\nabla_{e_{i}} e_{j}$ for our basis

$$
\left\{\mathbf{e}_{k}, k=1,2,3\right\}=\left\{\mathbf{e}_{1}, \mathbf{e}_{2}, \mathbf{e}_{3}\right\}
$$

So we obtain that

$$
\begin{equation*}
R\left(\mathbf{e}_{1}, \mathbf{e}_{3}\right)=R\left(\mathbf{e}_{1}, \mathbf{e}_{2}\right)=R\left(\mathbf{e}_{2}, \mathbf{e}_{3}\right)=0 \tag{2.4}
\end{equation*}
$$

Then, the Lorentz metric g is flat.
3. Spacelike Biharmonic Curves with Timelike Binormal According to Flat Metric in the Lorentzian Heisenberg Group Heis ${ }^{3}$

An arbitrary curve $\gamma: I \longrightarrow$ Heis $^{3}$ is spacelike, timelike or null, if all of its velocity vectors $\gamma^{\prime}(s)$ are, respectively, spacelike, timelike or null, for each $s \in I \subset \mathbb{R}$. Let $\gamma: I \longrightarrow$ Heis $^{3}$ be a unit speed spacelike curve with timelike binormal and $\{\mathbf{t}, \mathbf{n}, \mathbf{b}\}$ are Frenet vector fields, then Frenet formulas are as follows

$$
\begin{align*}
\nabla_{\mathbf{t}} \mathbf{t} & =\kappa_{1} \mathbf{n} \\
\nabla_{\mathbf{t}} \mathbf{n} & =-\kappa_{1} \mathbf{t}+\kappa_{2} \mathbf{b}  \tag{3.1}\\
\nabla_{\mathbf{t}} \mathbf{b} & =\kappa_{2} \mathbf{n}
\end{align*}
$$

where $\kappa_{1}, \kappa_{2}$ are curvature function and torsion function, respectively and

$$
\begin{aligned}
g(\mathbf{t}, \mathbf{t}) & =1, g(\mathbf{n}, \mathbf{n})=1, g(\mathbf{b}, \mathbf{b})=-1 \\
g(\mathbf{t}, \mathbf{n}) & =g(\mathbf{t}, \mathbf{b})=g(\mathbf{n}, \mathbf{b})=0
\end{aligned}
$$

With respect to the orthonormal basis $\left\{\mathbf{e}_{1}, \mathbf{e}_{2}, \mathbf{e}_{3}\right\}$ we can write

$$
\begin{aligned}
\mathbf{t} & =t_{1} \mathbf{e}_{1}+t_{2} \mathbf{e}_{2}+t_{3} \mathbf{e}_{3} \\
\mathbf{n} & =n_{1} \mathbf{e}_{1}+n_{2} \mathbf{e}_{2}+n_{3} \mathbf{e}_{3} \\
\mathbf{b} & =b_{1} \mathbf{e}_{1}+b_{2} \mathbf{e}_{2}+b_{3} \mathbf{e}_{3} .
\end{aligned}
$$

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Theorem 3.1. ( see [11]) If $\gamma: I \longrightarrow H e i s^{3}$ is a unit speed spacelike biharmonic curve with timelike binormal according to flat metric, then

$$
\begin{align*}
\kappa_{1} & =\text { constant } \neq 0 \\
\kappa_{1}^{2}-\kappa_{2}^{2} & =0  \tag{3.2}\\
\kappa_{2} & =\text { constant }
\end{align*}
$$

Theorem 3.2. ( see [11]) Let $\gamma: I \longrightarrow H e i s{ }^{3}$ is a unit speed spacelike biharmonic curve with timelike binormal according to flat metric. Then the parametric equations of $\gamma$ are

$$
\begin{align*}
x(s)= & \cosh \varphi s+C_{1} \\
y(s)= & \frac{1}{\kappa_{1}} \sinh ^{2} \varphi\left[\cosh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]+\sinh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]\right]+C_{2}, \\
z(s)= & -\frac{\left(-1+C_{1}+\cosh \varphi s\right) \sinh \varphi}{\kappa_{1}} \cosh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]  \tag{3.3}\\
& +\frac{\sinh ^{2} \varphi \cosh \varphi}{\kappa_{1}^{2}}\left[\sinh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]+\cosh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]\right] \\
& \left.-\frac{\sinh \varphi\left(\cosh \varphi s+C_{1}\right)}{\kappa_{1}} \sinh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]\right]+C_{3}
\end{align*}
$$

where $C, C_{1}, C_{2}, C_{3}$ are constants of integration.
4. Bertrand Mate of Spacelike Biharmonic Curves According to Flat Metric in the Lorentzian Heisenberg Group Heis ${ }^{3}$

A curve $\gamma: I \longrightarrow$ Heis $^{3}$ with $\kappa_{1} \neq 0$ is called a Bertrand curve if there exist a curve $\gamma_{\mathcal{B}}: I \longrightarrow$ Heis $^{3}$ such that the principal normal lines of $\gamma$ and $\gamma_{\mathcal{B}}$ at $s \in I$ are equal. In this case $\gamma_{\mathcal{B}}$ is called a Bertrand mate of $\gamma$.

On the other hand, let $\gamma: I \longrightarrow H e i s^{3}$ be a Bertrand curve parametrized by arc length. A Bertrand mate of $\gamma$ is as follows:

$$
\begin{equation*}
\gamma_{\mathcal{B}}(s)=\gamma(s)+f \mathbf{n}(s), \quad \forall s \in I \tag{4.1}
\end{equation*}
$$

where $f$ is constant.
Lemma 4.1. Let $\gamma: I \longrightarrow H e i s^{3}$ be a unit speed spacelike biharmonic curve with timelike binormal according to flat metric. Then the position vector of $\gamma$ is

$$
\begin{align*}
\gamma(s)= & \left(\cosh \varphi s+C_{1}\right) \mathbf{e}_{1} \\
& +\left[\left(\cosh \varphi s+C_{1}\right)\left(\frac{1}{\kappa_{1}} \sinh ^{2} \varphi\left[\cosh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]+\sinh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]\right]+C_{2}\right)\right. \\
& -\frac{\left(-1+C_{1}+\cosh \varphi s\right) \sinh \varphi}{\kappa_{1}} \cosh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right] \\
& +\frac{\sinh ^{2} \varphi \cosh \varphi}{\kappa_{1}^{2}}\left[\sinh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]+\cosh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]\right]  \tag{4.2}\\
& \left.\left.-\frac{\sinh \varphi\left(\cosh \varphi s+C_{1}\right)}{\kappa_{1}} \sinh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]\right]+C_{3}\right] \mathbf{e}_{2} \\
& {\left[\left(\frac{1}{\kappa_{1}} \sinh ^{2} \varphi\left[\cosh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]+\sinh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]\right]+C_{2}\right)\right.} \\
& -\left[\left(\cosh \varphi s+C_{1}\right)\left(\frac{1}{\kappa_{1}} \sinh { }^{2} \varphi\left[\cosh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]+\sinh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]\right]+C_{2}\right)\right. \\
& -\frac{\left(-1+C_{1}+\cosh \varphi s\right) \sinh \varphi}{\kappa_{1}} \cosh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right] \\
& +\frac{\sinh ^{2} \varphi \cosh \varphi}{\kappa_{1}^{2}}\left[\sinh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]+\cosh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]\right] \\
& \left.\left.-\frac{\sinh ^{2}\left(\cosh \varphi s+C_{1}\right)}{\kappa_{1}} \sinh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]\right]+C_{3}\right] \mathbf{e}_{3},
\end{align*}
$$

where $C, C_{1}, C_{2}, C_{3}$ are constants of integration.
Proof. Using (2.1) and (3.3), we have above system.
Theorem 4.2. Let $\gamma: I \longrightarrow H e i s^{3}$ be a unit speed spacelike biharmonic curve with timelike binormal and $\gamma_{\mathcal{B}}$ its Bertrand mate on Heis ${ }^{3}$. Then,

$$
\begin{aligned}
\gamma_{\mathcal{B}}(s)= & {\left[\cosh \varphi s-\frac{f}{\kappa_{1}} \sinh ^{2} \varphi\left(\sinh ^{2}\left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right]\right.\right.} \\
& \left.\left.+\sinh \left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right] \cosh \left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right]\right)+C_{1}\right) \mathbf{e}_{1} \\
& +\left[\left(\cosh \varphi s+C_{1}\right)\left(\frac{1}{\kappa_{1}} \sinh ^{2} \varphi\left[\cosh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]+\sinh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]\right]+C_{2}\right)\right. \\
& -\frac{\left(-1+C_{1}+\cosh \varphi s\right) \sinh \varphi}{\kappa_{1}} \cosh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]
\end{aligned}
$$

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$$
\begin{align*}
& +\frac{\sinh ^{2} \varphi \cosh \varphi}{\kappa_{1}^{2}}\left[\sinh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]+\cosh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]\right]  \tag{4.3}\\
& \left.-\frac{\sinh \varphi\left(\cosh \varphi s+C_{1}\right)}{\kappa_{1}} \sinh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]\right] \\
& +\frac{f}{\kappa_{1}}\left(\kappa_{1} \cosh \left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right]+\sinh \varphi \cosh \varphi \sinh \left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right]\right. \\
& \left.\left.+\sinh \varphi \cosh \varphi \cosh \left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right]\right)+C_{3}\right] \mathbf{e}_{2} \\
& {\left[\left(\frac{1}{\kappa_{1}} \sinh ^{2} \varphi\left[\cosh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]+\sinh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]\right]+C_{2}\right)\right.} \\
& -\left[\left(\cosh \varphi s+C_{1}\right)\left(\frac{1}{\kappa_{1}} \sinh ^{2} \varphi\left[\cosh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]+\sinh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]\right]+C_{2}\right)\right. \\
& -\frac{\left(-1+C_{1}+\cosh \varphi s\right) \sinh \varphi}{\kappa_{1}} \cosh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right] \\
& +\frac{\sinh { }^{2} \varphi \cosh \varphi}{\kappa_{1}^{2}}\left[\sinh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]+\cosh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]\right] \\
& \left.-\frac{\sinh \varphi\left(\cosh \varphi s+C_{1}\right)}{\kappa_{1}} \sinh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]\right] \\
& +\frac{f}{\kappa_{1}}\left(\kappa_{1} \sinh \left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right]-\sinh \varphi \cosh \varphi \sinh \left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right]\right. \\
& \left.\left.-\sinh \varphi \cosh \varphi \cosh \left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right]\right)+C_{3}\right] \mathrm{e}_{3},
\end{align*}
$$

where $C, C_{1}, C_{2}, C_{3}$ are constants of integration.
Proof. We assume that $\gamma: I \longrightarrow$ Heis $^{3}$ be a unit speed spacelike biharmonic curve.

Using Lemma 4.1, we get

$$
\begin{equation*}
\mathbf{t}=\cosh \varphi \mathbf{e}_{1}+\sinh \varphi \sinh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right] \mathbf{e}_{2}+\sinh \varphi \cosh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right] \mathbf{e}_{3} . \tag{4.4}
\end{equation*}
$$

Therefore, (4.4) becomes

$$
\begin{align*}
\mathbf{t} & =\left(\cosh \varphi, \sinh \varphi \sinh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]+\sinh \varphi \cosh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right],\right.  \tag{4.5}\\
& \left.(1-x) \sinh \varphi \sinh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]-x \sinh \varphi \cosh \left[\frac{\kappa_{1} s}{\sinh \varphi}+C\right]\right) .
\end{align*}
$$

Using first equation of the system (3.2) and (2.3), we have

$$
\begin{aligned}
\nabla_{\mathbf{t}} \mathbf{t}=(\quad & \left.t_{1}^{\prime}-t_{2}^{2}-t_{2} t_{3}\right) \mathbf{e}_{1}+\left(t_{2}^{\prime}+t_{1} t_{2}+t_{1} t_{3}\right) \mathbf{e}_{2} \\
& +\left(t_{3}^{\prime}-t_{1} t_{2}-t_{1} t_{3}\right) \mathbf{e}_{3}
\end{aligned}
$$

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On the other hand, from above equation and (3.1), we obtain

$$
\begin{align*}
\nabla_{\mathbf{t}} \mathbf{t}= & -\sinh ^{2} \varphi\left(\sinh ^{2}\left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right]+\sinh \left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right] \cosh \left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right]\right) \mathbf{e}_{1} \\
& +\left(\kappa_{1} \cosh \left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right]+\sinh \varphi \cosh \varphi \sinh \left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right]\right. \\
& \left.+\sinh \varphi \cosh \varphi \cosh \left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right]\right) \mathbf{e}_{2}  \tag{4.6}\\
& +\left(\kappa_{1} \sinh \left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right]-\sinh \varphi \cosh \varphi \sinh \left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right]\right. \\
& \left.-\sinh \varphi \cosh \varphi \cosh \left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right]\right) \mathbf{e}_{3} .
\end{align*}
$$

By the use of Frenet formulas and above equation, we get

$$
\begin{align*}
\mathbf{n}= & -\frac{1}{\kappa_{1}} \sinh ^{2} \varphi\left(\sinh ^{2}\left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right]+\sinh \left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right] \cosh \left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right]\right) \mathbf{e}_{1} \\
& +\frac{1}{\kappa_{1}}\left(\kappa_{1} \cosh \left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right]+\sinh \varphi \cosh \varphi \sinh \left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right]\right. \\
& \left.+\sinh \varphi \cosh \varphi \cosh \left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right]\right) \mathbf{e}_{2}  \tag{4.7}\\
& +\frac{1}{\kappa_{1}}\left(\kappa_{1} \sinh \left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right]-\sinh \varphi \cosh \varphi \sinh \left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right]\right. \\
& \left.-\sinh \varphi \cosh \varphi \cosh \left[\frac{\kappa_{1} s}{\cosh \varphi}+C\right]\right) \mathbf{e}_{3} .
\end{align*}
$$

Combining (4.7) and (4.2), we obtain (4.3).

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